

Soliton and Some of its Geophysical/Oceanographic Applications

M. Lakshmanan

Centre for Nonlinear Dynamics
Bharathidasan University
Tiruchirappalli – 620024
India

Discussion Meeting, ICTS-TIFR

21 July 2016

Theme

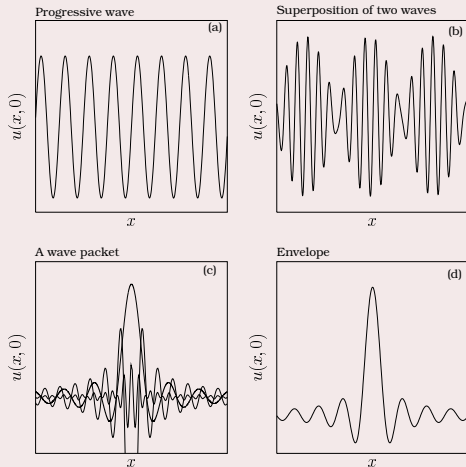
Soliton is a counter-intuitive wave entity arising because of a delicate balance between dispersion and nonlinearity. Its major attraction is its localized nature, finite energy and preservation of shape under collision \implies a remarkably stable structure. Consequently, it has ramifications in as wider areas as hydrodynamics, tsunami dynamics, condensed matter, magnetism, particle physics, nonlinear optics, cosmology and so on. A brief overview will be presented with emphasis on oceanographic applications and potential problems will be indicated.

Plan

- ➊ Introduction
- ➋ Historical: (i) Scott-Russel (ii) Zabusky/Kruskal
- ➌ Korteweg-de Vries equation and solitary wave/soliton
- ➍ Other ubiquitous equations
- ➎ Mathematical theory of solitons
- ➏ Geophysical/Oceanographic applications
- ➐ Solitons in higher dimensions
- ➑ Problems and potentialities

Solitons

- Soliton as a counter-intuitive entity
- Linear dispersive waves

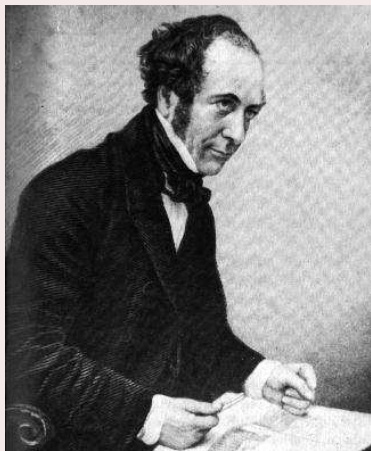


Solitons

- Nonlinear dispersive waves: Solitary waves/solitons
- Many facets of solitons:
 - a) Tsunami as soliton?
 - b) Internal soliton
 - c) Rossby soliton
 - d) Bore soliton
 - e) Capillary soliton
 - f) Optical soliton
 - g) Magnetic soliton
 - h) Topological soliton
 - i) Plane soliton
 - j) Lump soliton
 - k) Dromion
 - l) Light bullet

John Scott-Russel

John Scott-Russel (1808-82)



- Russel's observations -
Phenomenological relation:

$$c = \sqrt{g(h + k)}$$

J. S. Russel, British Association Reports (1844)

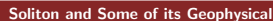
“I was observing the motion of a boat which was rapidly drawn along a narrow channel by a pair of horses, when the boat suddenly stopped - not so the mass of water in the channel which it had put in motion; it accumulated round the prow of the vessel in a state of violent agitation, then suddenly leaving it behind, rolled forward with great velocity, assuming the form of a large solitary elevation, a rounded, smooth and well-defined heap of water, which continued its course along the channel apparently without change of form or diminution of speed. I followed it on horseback, and overtook it still rolling on at a rate of some eight or nine miles an hour, preserving its original figure some thirty feet long and a foot to a foot and a half in height. Its height gradually diminished, and after a chase of one or two miles I lost it in the windings of the channel. Such, in the month of August 1834, was my first chance interview with that singular and beautiful phenomenon! which I have called the Wave of Translation”.

Solitons

Soliton creation



- Scott Russell Aqueduct – Soliton creation – Heriot-Watt University, 1995.
- Korteweg-de Vries Formulation (1895)
- Unidirectional shallow water wave propagation in one dimension
- Incompressible, inviscible fluid



Solitons

Fluid motion:

$$\vec{V}(x, y, t) = u(x, y, t)\vec{i} + v(x, y, t)\vec{j}$$

$$\Rightarrow \text{Irrotational: } \vec{\nabla} \times \vec{V} = 0 \quad \Rightarrow \quad \vec{V} = \vec{\nabla}\phi$$

$$\phi(x, y, t) : \text{ velocity potential}$$

(i) Conservation of density:

$$\frac{d\rho}{dt} = \rho_t + \vec{\nabla} \cdot (\rho \vec{V}) = 0, \quad \rho = \rho_0 = \text{constant}$$

$$\Rightarrow \nabla^2 \phi(x, y, t) = 0 - \text{Laplace equation.}$$

Solitons

(ii) Euler's Equation:

$$\begin{aligned} \frac{d\vec{V}}{dt} &= \frac{\partial \vec{V}}{\partial t} + (\vec{V} \cdot \vec{\nabla}) \vec{V} \\ &= -\frac{1}{\rho_0} \vec{\nabla} p - g\vec{j} \end{aligned} \quad \Rightarrow \quad \phi_t + \frac{1}{2} (\vec{\nabla} \phi)^2 + \frac{p}{\rho_0} + gy = 0$$

(iii) Boundary conditions:

$$\phi_y(x, 0, t) = 0$$

$$y = h + \eta(x, t); \quad \phi_{1y} = \eta_t + \eta_x \phi_{1x}$$

$$p_1 = 0 \quad \Rightarrow \quad u_{1t} + u_1 u_{1x} + v_1 v_{1x} + g\eta_x = 0$$

Solitons

(iv) Taylor Expansion:

$$\phi(x, y, t) = \sum_{n=0}^{\infty} y^n \phi_n(x, t)$$

$$u_1 = \phi_{1x} = f - \frac{1}{2} y_1^2 f_{xx} + \dots,$$

$$v_1 = \phi_{1y} = -y_1 f_x + \frac{1}{6} y_1^3 f_{xxx} + \dots, \quad f = \frac{\partial \phi_0}{\partial x}$$

(v) Introducing ϵ and δ :

$$u_1 = \epsilon c_0 u'_1, \quad v_1 = \epsilon \delta c_0 v'_1, \quad f = \epsilon c_0 f'$$

$$y_1 = h + \eta(x, t) = h (1 + \epsilon \eta' (x', t'))$$

Solitons

$$\eta_t + f_x + \epsilon \eta f_x + \epsilon f \eta_x - \frac{1}{6} \delta^2 f_{xxx} = 0$$

$$f_t + \epsilon f f_x + \frac{ga}{\epsilon c_0^2} \eta_x - \frac{1}{2} \delta^2 f_{xxt} = 0$$

(vi) Perturbation analysis:

$$f = f^{(0)} + \epsilon f^{(1)} + \delta^2 f^{(2)} + \text{higher order}$$

$$f^{(0)} = \eta, \quad f^{(1)} = -\frac{1}{4} \eta^2, \quad f^{(2)} = \frac{1}{3} \eta_{xx}$$

$$\eta_t + \eta_x + \frac{3}{2} \epsilon \eta \eta_x + \frac{\delta^2}{6} \eta_{xxx} = 0$$

Solitons

(vii) The standard contemporary form:

$$u_t \pm 6uu_x + u_{xxx} = 0$$

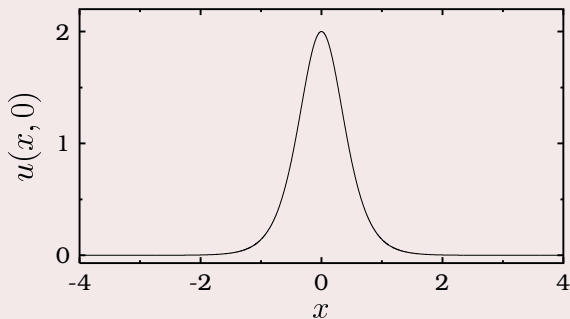
(ix) Cnoidal and solitary wave solution:

$$u(x, t) = 2f(\zeta), \quad \zeta = (x - ct) \implies -c \frac{\partial f}{\partial \zeta} + 12f \frac{\partial f}{\partial \zeta} + \frac{\partial^3 f}{\partial \zeta^3} = 0$$

$$\implies f(\zeta) = \alpha_3 - (\alpha_3 - \alpha_2) \operatorname{sn}^2(\sqrt{\alpha_3 - \alpha_1}(x - ct), m)$$

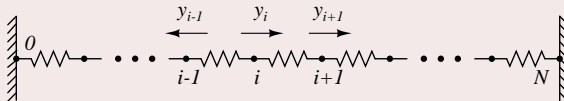
$$\implies \frac{c}{2} \operatorname{sech}^2 \frac{\sqrt{c}}{2}(x - ct + \delta), \quad \delta = \text{constant} \quad (m = 1)$$

Solitons



III. FPU Problem (1955) and Zabusky-Kruskal Analysis (1965)

Wave propagation in anharmonic lattices



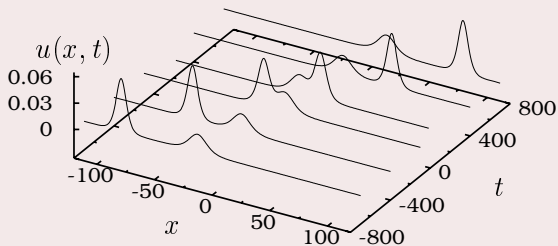
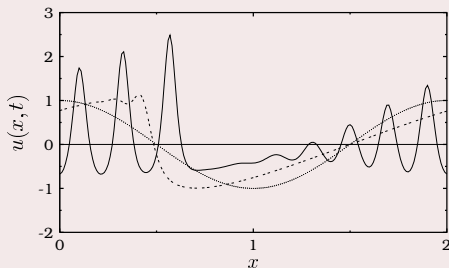
Equation of motion:

$$m \frac{\partial^2 y_i}{\partial t^2} = f(y_{i+1} - y_i) - f(y_i - y_{i-1}), \quad i = 1, 2, \dots, N-1,$$

$$y_i \longrightarrow y(x, t), \quad y_{i+1} = y(x, t) \pm \frac{\partial y}{\partial x} \pm \frac{a^2}{2} \frac{\partial^2 y}{\partial x^2} \pm \dots$$

$$\implies \text{KdV equation:} \quad u_t + 6uu_x + u_{xxx} = 0$$

Zabusky-Kruskal Numerical Analysis: Birth of soliton



Soliton

Solitary wave which retains its shape and speed under collision except for a phase shift – Elastic collision.

Other ubiquitous equations

Boussinesq equation

$$u_t + uu_x + g\eta_x - \frac{1}{3}h^2 u_{txx} = 0$$
$$\eta_t + [u(h + \eta)]_x = 0$$

Benjamin-Bona-Mahoney (BBM) equation

$$u_t + u_x + uu_x - u_{xxt} = 0$$

Camassa-Holm equation

$$u_t + 2\kappa u_x + 3uu_x - u_{xxt} = 3u_x u_{xx} + uu_{xxx}$$

Other ubiquitous equations

Kadomtsev-Petviashvile (KP) equation

$$(u_t + 6uu_x + u_{xxx})_x + 3\sigma^2 u_{yy} = 0$$

$(\sigma^2 = -1: \text{KP-I}, \sigma^2 = +1: \text{KP-II})$

Modified KdV equation

$$u_t \pm 6u^2 u_x + u_{xxx} = 0$$

Sine-Gordon (sG) equation (in light-cone coordinates)

$$u_{xt} + m^2 \sin u = 0, \quad m^2 = \text{constant}$$

Other ubiquitous equations

Nonlinear Schrödinger equation

$$iq_t + q_{xx} \pm 2|q|^2 q = 0, \quad q \in \mathbb{C}$$

Continuous Heisenberg ferromagnetic spin equation

$$\vec{S}_t = \vec{S} \times \vec{S}_{xx}, \quad \vec{S} = (S_1, S_2, S_3), \quad \vec{S}^2 = 1$$

Toda lattice equation

$$\ddot{u}_n = \exp[-(u_n - u_{n-1})] - \exp[-(u_{n+1} - u_n)], \quad i = 1, 2, \dots, N$$

Mathematical theory of solitons: Cauchy IVP

KdV equation \implies Linearization

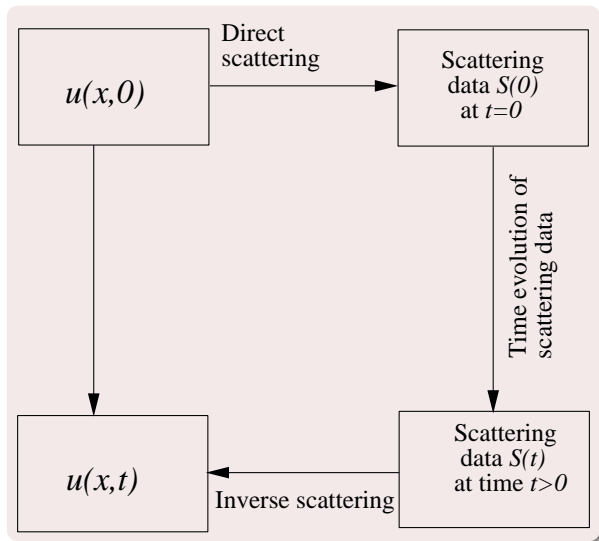
Given: $u(x, 0)$: Initial value

(i) $\psi_{xx} + [\lambda - u(x, t)]\psi = 0, \quad u(x, t) \longrightarrow 0, \text{ at } |x| \rightarrow \infty$

(ii) $\psi_t = -4\psi_{xxx} + 3u\psi_x + 6u_x\psi$

$\implies (\psi_{xx})_t \equiv (\psi_t)_{xx} \implies \text{KdV}$

Inverse Scattering Transform method



Inverse Scattering Transform method

1-soliton solution:

$$u = -2\kappa^2 \operatorname{sech}^2 [\kappa (x - 4\kappa^2 t) + \delta]$$

2-soliton solution:

$$u(x, t) = -2 (\kappa_2^2 - \kappa_1^2) \frac{\kappa_2^2 \operatorname{cosech}^2 \gamma_2 + \kappa_1^2 \operatorname{sech}^2 \gamma_1}{(\kappa_2 \coth \gamma_2 - \kappa_1 \tanh \gamma_1)^2}$$

N-soliton solution

- Bilinearization: $u = 2 \frac{\partial^2}{\partial x^2} \log F$



$$F_{xt}F - F_xF_t + F_{xxxx}F - 4F_{xxx}F_x + 3F_{xx}^2 = 0$$

- Bäcklund transformation
- Infinite number of conservation laws/symmetries
- Geometric structures/Group theoretic
- Hamiltonian/Bi-Hamiltonian structure
- Completely integrable infinite dimensional nonlinear dynamical systems

VI. Geophysical applications

A. Tsunami as solitons?



Geophysical applications

Tsunami as solitons?



Geophysical applications

Tsunami as solitons?

- 2004 Indian ocean / 1960 Chilean earth quake
- Typical $\lambda \simeq 200$ kms, $h \simeq 1$ to 4 kms, $a \simeq 60$ cms
- Shallowness condition: $\epsilon = \frac{a}{h} \simeq 10^{-4} \ll 1$, $\delta^2 = \frac{h^2}{l^2} \simeq 10^{-4}$
- Propagation of tsunami waves as KdV like soliton

Recent generalization

- Phase modulated solitary waves controlled by boundary conditions at the bottom.
(A. Mukherjee and M.S. Janaki, Phys. Rev E(2014)).
- Lump solitons & Rogue waves controlled by ocean currents.
(A. Kundu, A. Mukherjee, T. Naskar, Proc. R. Soc. Lond. (2014))

Geophysical applications

B. Internal solitons?

- Seafarers through strait of Malacca often note that in the Andaman sea bands of strongly increasing surface roughness occur.
⇒ Internal solitons (SAR images of ERS-1/2 satellites) (Solitary wave like distortion of the boundary layer between warm upper sea and cold lower depths)
- Peculiar striations of 100 km long, separated by 6-15 kms and grouped in packets of 4-8 of the surface of Andaman and Sulu seas.
⇒ Secondary phenomena of internal solitons.
- Internal solitons are traveling edges of warm water with enormous energy affecting oil rigs

Geophysical applications

B. Internal solitons

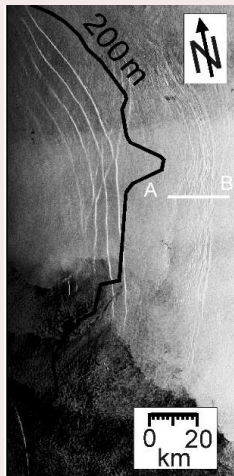
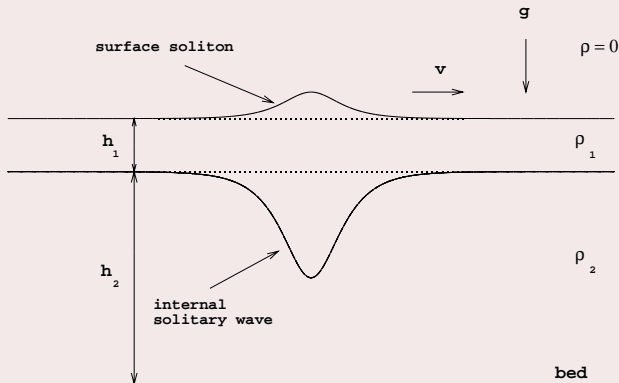


Figure: SAR image of the Andaman Sea acquired by the ERS-2 satellite showing two internal solitary wave packets

Geophysical applications

B. Internal solitons



Geophysical applications

B. Internal solitons

- Study of Osborne and Birch (1980) of underwater currents by an oil rig drilling at 36000 ft depth in the Andaman sea \sim KdV solitons
- Other observations
- Theory: Consider two incompressible, immiscible fluid, densities ρ_1 , ρ_2 and depths h_1 , h_2
- KdV
- Intermediate long wavelength (ILW)
- Benjamin-Ono

Geophysical applications

C. Rossby solitons

- In the atmosphere of a rotating planet, fluid particle \implies rotation rate (Coriolis freq)/latitude
- Motion in N/S direction inhibited by constant angular momentum

Rossby Waves:

- Large scale atmospheric waves caused by variation of Coriolis frequency with latitude \implies KdV equation can model Rossby waves (E-W zonal flow) (long wave/incompressible fluid/ β plane approximation etc)
Experimental confirmation: Eight years of Topex/Poseidon altimeter observations. (Pacific dynamics) – 1998

Geophysical applications

D. Bore solitons

- When a tidal wave/tsunami/storm surge hits an estuary \implies hydraulic jump (step-wise perturbations like a shock wave) in the water height and speed which propagates upstream.
- Far less dangerous but very similar is the bore (mascaret in French), a tidal wave which can propagate in river for considerable distances.

Examples

- Seine river / Hooghly river
- 1983 Japan sea tsunami \implies bores ascended many rivers
- IHOP of Kansas

Geophysical applications



Figure: Tidal bore at the mouth of the Aragnau River (Brazil)

- E. Capillary solitons
- F. Resonant three and four wave solitons
- G. Optical solitons

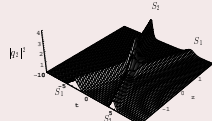
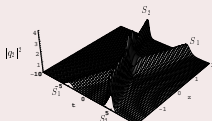
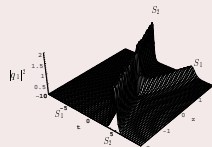
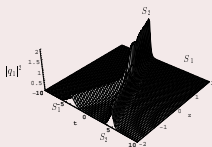
VII. Solitons in other situations

A. Multicomponent solitons:

- Manakov model

$$iq_{1,z} + q_{1,tt} + 2(|q_1|^2 + |q_2|^2)q_1 = 0,$$

$$iq_{2,z} + q_{2,tt} + 2(|q_1|^2 + |q_2|^2)q_2 = 0.$$

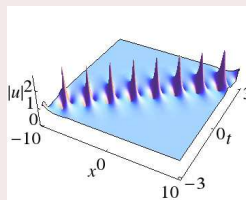


Breathers and Rogue Waves

- Nonlinear Schrödinger Equation

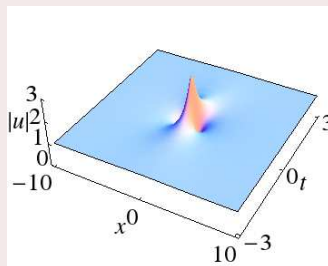
$$iu_t + u_{xx} + q|u|^2 u = 0, \quad u \in \mathbb{C}$$

- Coupled NLS Manakov Equation



Breather solution of NLS equation

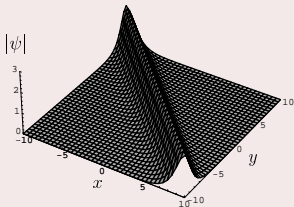
Breathers and Rogue Waves



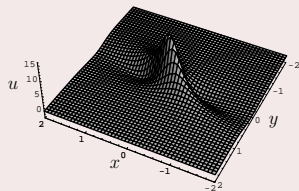
Rogue wave solution of NLS equation

- Rogue waves in BECs/Optics

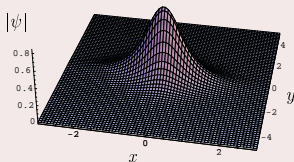
B. (2+1)- dimensions:



- Plane solitons



- Lump solitons



● Dromions

C. (3+1)- dimensions

● Localized solutions?

VIII. Conclusion

- Solitons are very stable nonlinear entities
- Occur in $(1+1)$ and $(2+1)$ dimensions
- Have considerable physical relevance
- What happens in $(2+1)$ and $(3+1)$ dimensions?

References



M. Lakshmanan and S. Rajasekar (2003) *Nonlinear Dynamics: Integrability and Chaos*, Springer, Berlin



M. Lakshmanan (2009) *Tsunamis and other Oceanographical Applications of Solitons* To appear in *Encyclopedia on Complex Systems*, Springer, Berlin



M. J. Ablowitz and P. A. Clarkson (1991) *Solitons, Nonlinear Evolution Equations and Inverse Scattering*. Cambridge University Press, Cambridge (UK)



T. Dauxois and M. Peyrard (2006) *Physics of Solitons*, Cambridge University Press, Cambridge (UK)