



Non-Hermitian matter-wave mixing in Bose-Einstein condensates: *Dissipation-induced amplification*

PRA **96** (2017) 013605

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MichiganTech

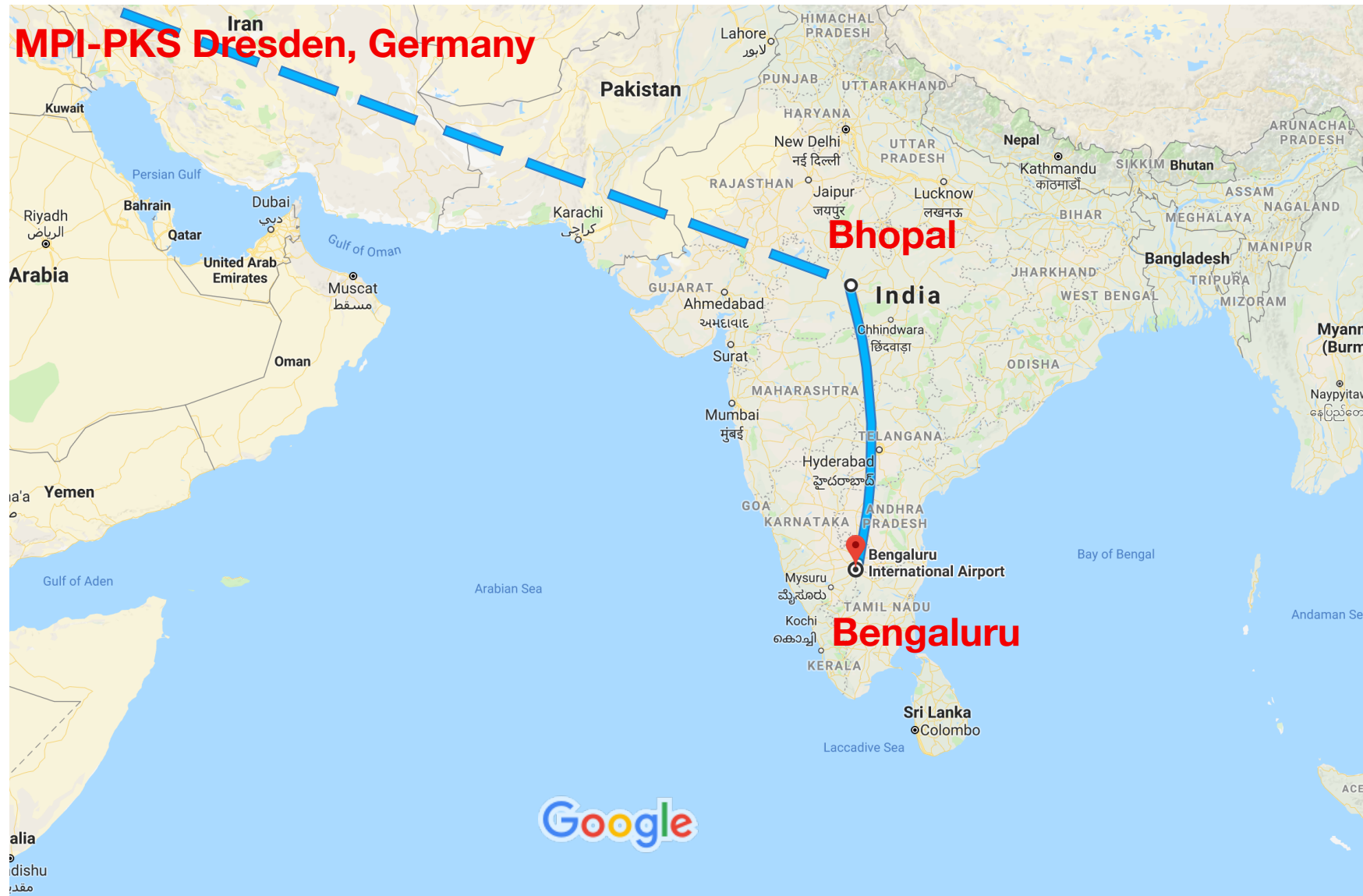
² Department of Physics, Michigan Technological University,
Houghton, Michigan 49931, USA



Geography

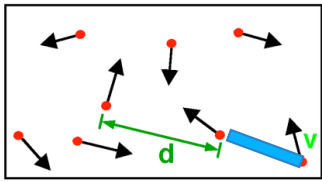
Ultra-cold and Rydberg atom theory in Bhopal: Max-Planck-IISER Partner Group

MAX-PLANCK-GESELLSCHAFT

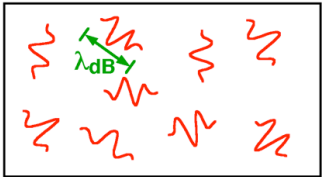


Motivation: BEC/ atomic collisions

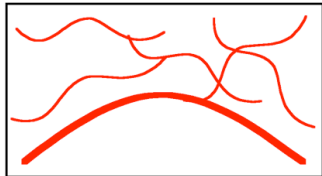
What is Bose-Einstein condensation (BEC)?



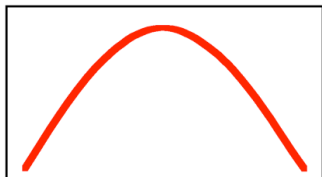
High Temperature T:
thermal velocity v
density d^{-3}
"Billiard balls"



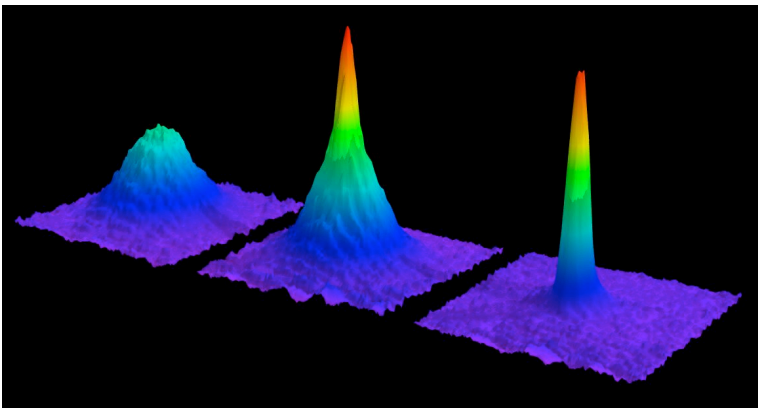
Low Temperature T:
De Broglie wavelength
 $\lambda_{dB} = h/mv \propto T^{-1/2}$
"Wave packets"



$T = T_{crit}$:
Bose-Einstein Condensation
 $\lambda_{dB} \approx d$
"Matter wave overlap"



$T=0$:
Pure Bose condensate
"Giant matter wave"



Images Ketterle group

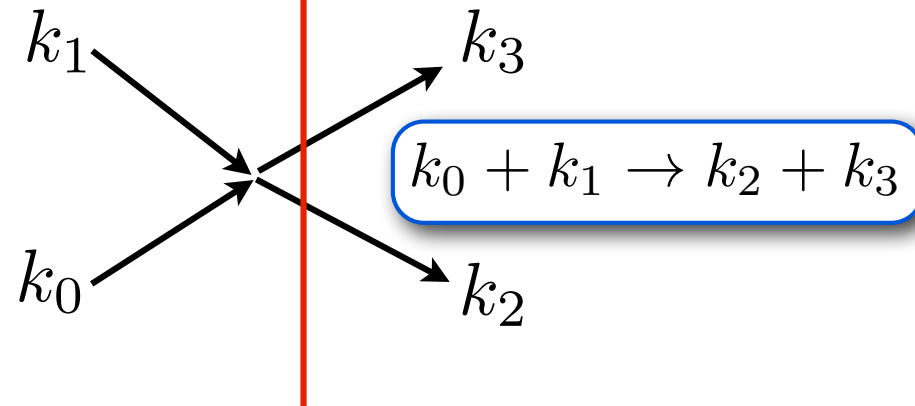
K.B. Davis et al., Phys. Rev. Lett. **75** (1995) 3969.

One wave function for all atoms

$$GPE: i\hbar \frac{\partial \Psi(x)}{\partial t} = \left[-\frac{\hbar^2}{2M} \frac{\partial^2}{\partial x^2} + U|\Psi(x)|^2 \right] \Psi(x)$$

Gross-Pitaevskii equation (GPE)

s-wave collisions



Non-interacting atom,
dispersion relation

$$E(k) = \frac{\hbar^2 k^2}{2m}$$

$$E(k_0) + E(k_1) = E(k_2) + E(k_3)$$

Motivation: wave mixing

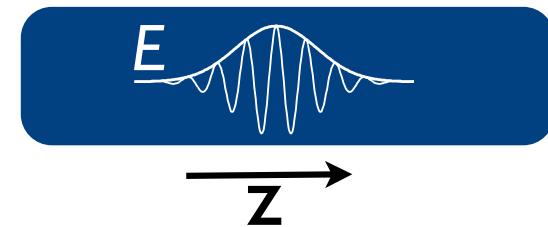
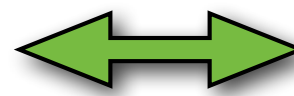
Analogy to non-linear optics with cubic (Kerr) nonlinearity

$$i\hbar \frac{\partial \Psi(x)}{\partial t} = \left[-\frac{\hbar^2}{2M} \frac{\partial^2}{\partial x^2} + U|\Psi(x)|^2 \right] \Psi(x)$$

GPE

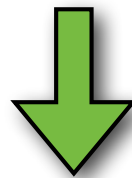
$$i \frac{\partial E}{\partial z} = \left[-a \frac{\partial^2}{\partial \tau^2} + b|E|^2 \right] E$$

NLSE



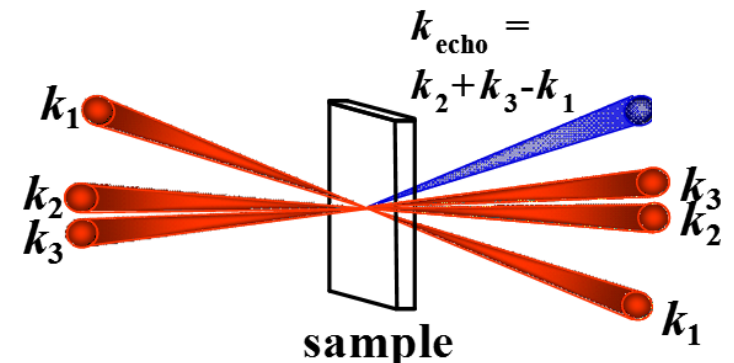
Nonlinear optics, wave mixing

$$E(t) = \frac{1}{2} E_1 e^{-i\omega_1 t} + \frac{1}{2} E_2 e^{-i\omega_2 t} + \text{c.c.}$$



$$P(t) = \epsilon_0 \chi^{(2)} E^2(t) = \frac{\epsilon_0 \chi^{(2)}}{4} \left[E_1^2 e^{-i2\omega_1 t} + E_2^2 e^{-i2\omega_2 t} + 2E_1 E_2 e^{-i(\omega_1 + \omega_2)t} + 2E_1 E_2 e^{-i(\omega_1 - \omega_2)t} + \dots \right]$$

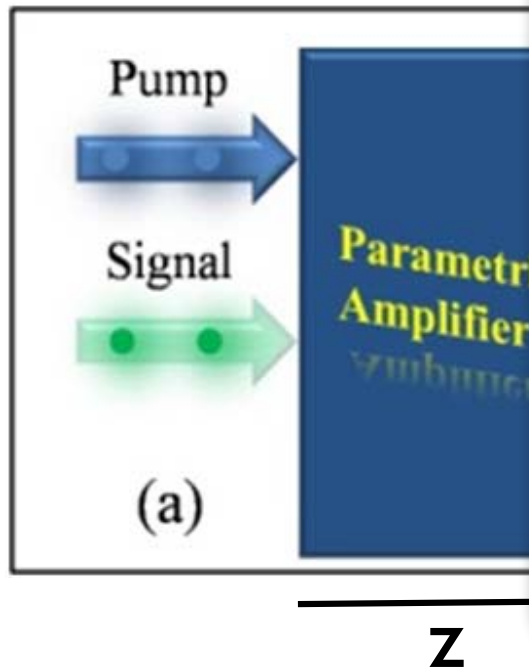
Phase matching
 $\exp [i(\omega_n t - \underline{k_n x})]$



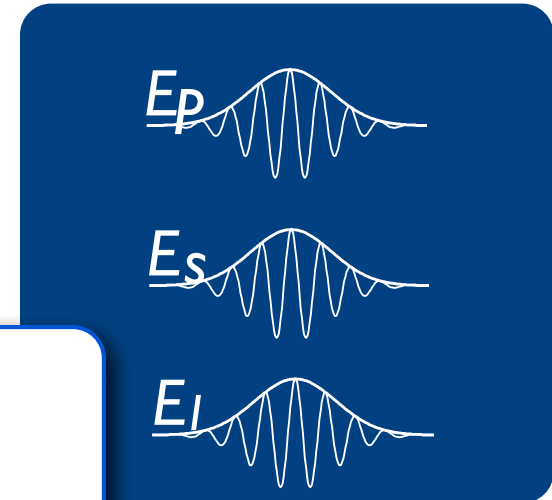
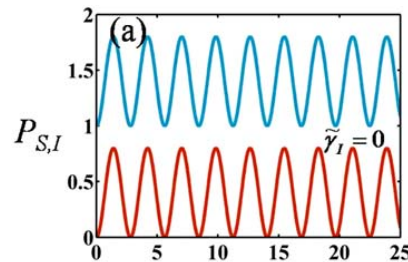
(img: Fayer lab webpage, Stanford)

Motivation: non-Hermitian optical amplification

Optical parametric amplification (OPA) / three-wave mixing



No amplification



$$\exp(i\Delta\beta z) \cdot \text{phase-matching}$$

$$\exp(i\Delta\beta z)$$

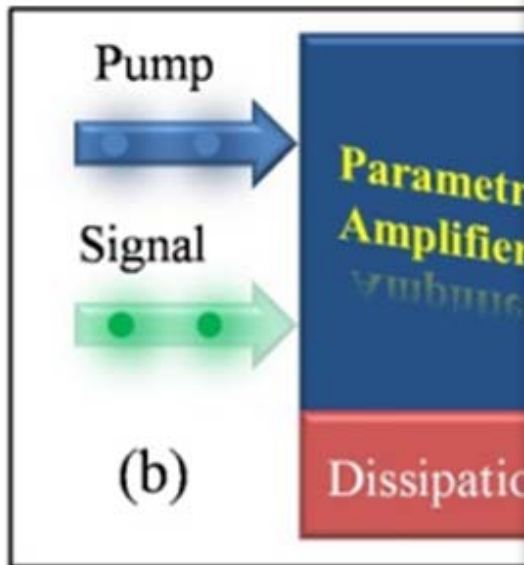
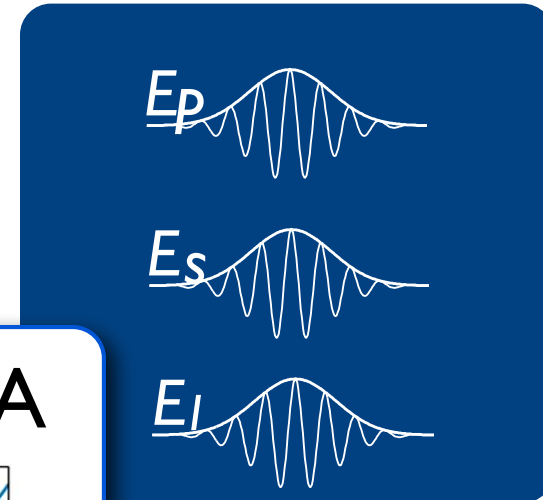
$$\exp(-i\Delta\beta z)$$

$$\Delta\beta = \beta_P - (\beta_S + \beta_I)$$

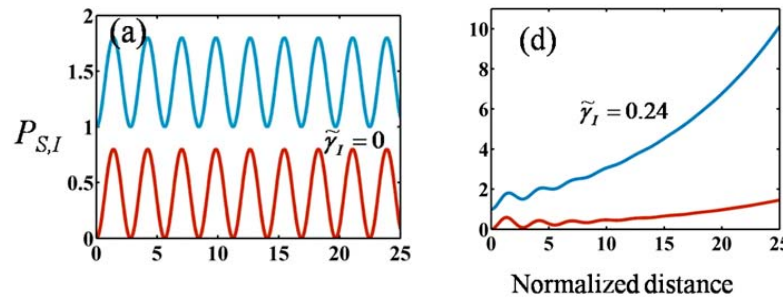
β dimless wavenumbers

Motivation: non-Hermitian optical amplification

Optical parametric amplification (OPA) / three-wave mixing



Non-Hermitian OPA



R. El-Ganainy et al., Opt. Lett. **40** (2015) 5086.
 see also:
 T. Wasak et al., Opt. Lett. **40** (2015) 5291.
 D.A. Antonosyan et al., Opt. Lett. **40** (2015) 4575.

$\exp(i\Delta\beta z)$
 phase-matching

$\exp(i\Delta\beta z) - \gamma_I E_I$

$\exp(-i\Delta\beta z)$,

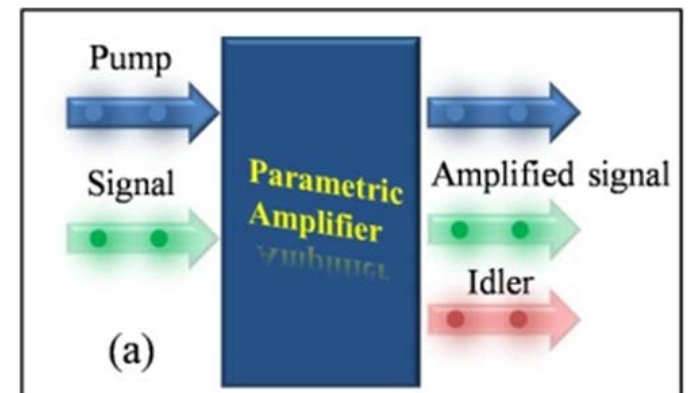
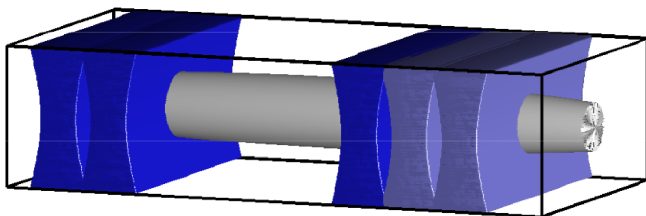
$-(\beta_S + \beta_I)$
 wavenumbers

Ramifications for BEC?

R. El-Ganainy et al.

Outline

1. Motivation
2. Introduction, BEC four-wave mixing
3. Matter-wave amplification in homogeneous 1D condensate
4. Two-dimensional wave mixing
5. Engineering atomic losses
6. Summary and Outlook



BEC four-wave mixing

GPE again, position space

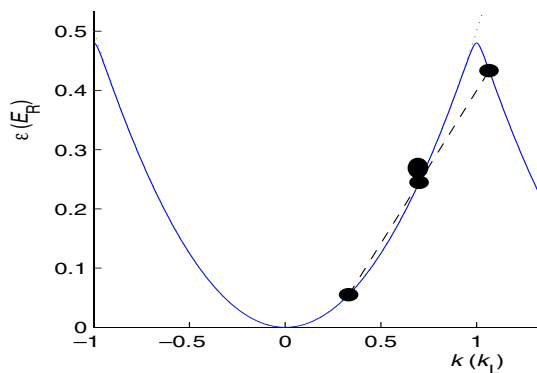
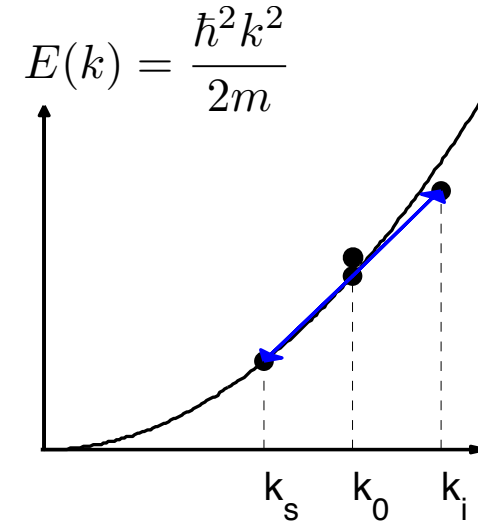
$$i\hbar \frac{\partial \Psi(x)}{\partial t} = \left[-\frac{\hbar^2}{2M} \frac{\partial^2}{\partial x^2} + U |\Psi(x)|^2 \right] \Psi(x)$$

BEC four-wave mixing

GPE in momentum-space and rotating frame

$$i\hbar \frac{\partial \phi_p}{\partial t} = U \sum_{n,m,l} \delta_{k_p+k_n-k_m-k_l} \times \phi_n^* \phi_m \phi_l e^{i(E_p+E_n-E_m-E_l)t/\hbar}$$

- for large Δk = four-wave mixing
- 1D: no energy-momentum conservation / phase matching

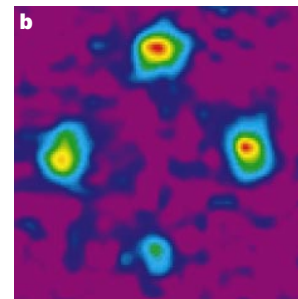
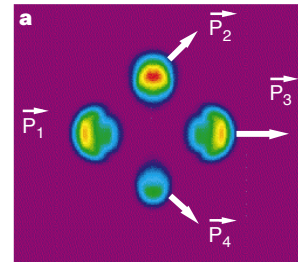


K. Mølmer, NJP **8** (2006) 170.

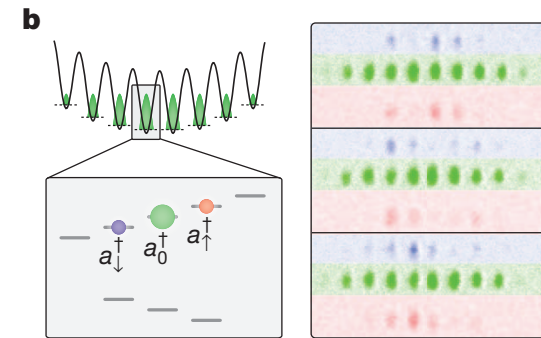
G. K. Campbell *et al.*
PRL **96** (2006) 020406.

change dispersion

L. Deng *et al.*
Nature **398**
(1999) 218.



2D



C. Gross *et al.*
Nature **480** (2011) 219.

R. Bückner *et al.*
Nat. Phys. **7** (2011) 608.

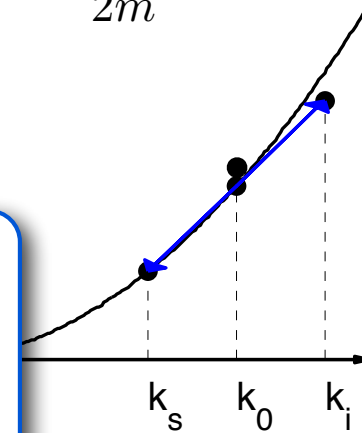
internal /trap energy

BEC four-wave mixing

GPE in momentum-space and rotating frame

$$i\hbar \frac{\partial \phi_p}{\partial t} = U \sum_{n,m,l} \delta_{k_p+k_n-k_m-k_l} \times \phi_n^* \phi_m \phi_l e^{i(E_p+E_n-E_m-E_l)t/\hbar}$$

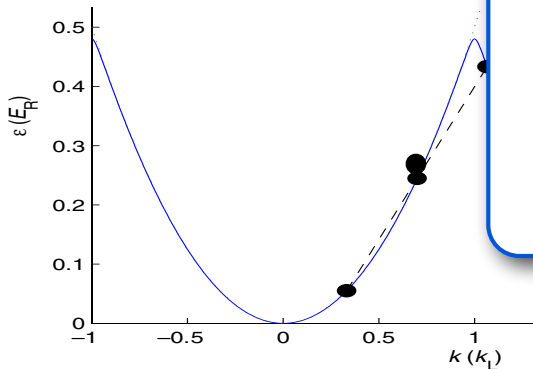
$$E(k) = \frac{\hbar^2 k^2}{2m}$$



- for large $\Delta k =$ four-wave mixing
- 1D: no energy- n phase matching

Four-wave mixing applications

- fundamental studies
- EPR entangled atom pairs
- Atom lasers and interferometry

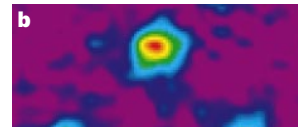


K. Mølmer, NJP **8** (2006) 170

G. K. Campbell et al. PRL **96** (2006) 020406

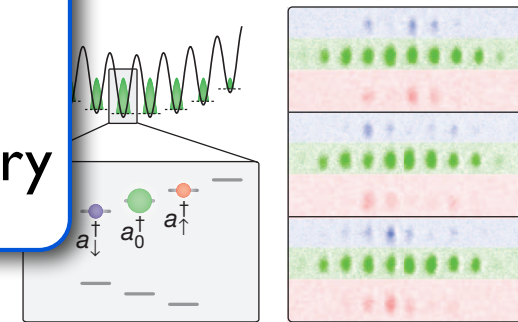
change dispersion

Nature **398** (1999) 218.



2D

0.57 mm



C. Gross et al. Nature **480** (2011) 219.

ner et al. Phys. **7** (2011) 608.

internal /trap energy

New channels in 1D?

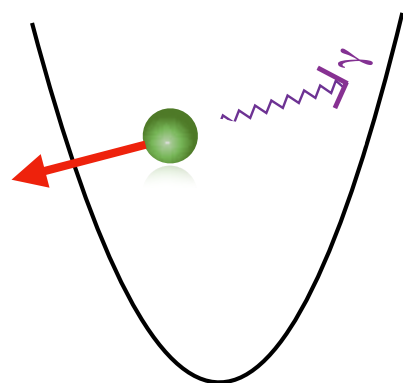
BEC - Non-Hermiticity

$$\text{GPE: } i\hbar \frac{\partial \Psi(x)}{\partial t} = \left[-\frac{\hbar^2}{2M} \frac{\partial^2}{\partial x^2} + U|\Psi(x)|^2 \right] \Psi(x)$$

Many-body H
Hermitian

Unavoidable atom losses:

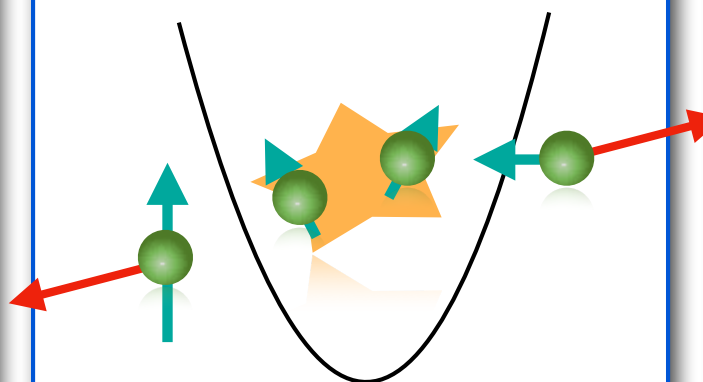
One-body loss



- vacuum imperfections
- BB-photons
- trap photon recoil

$$\dot{\Psi} = -\gamma \Psi$$

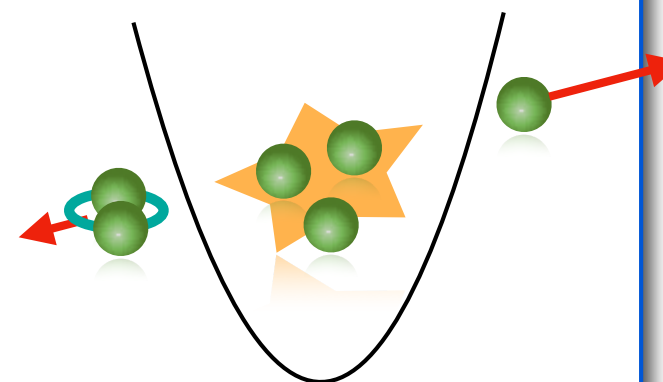
Two-body loss



- spin-flip collisions

$$\dot{\Psi} = -K_2 |\Psi|^2 \Psi$$

Three-body loss



- molecule formation
- main density & life-time limitation

$$\dot{\Psi} = -K_3 \frac{|\Psi|^4}{2} \Psi$$

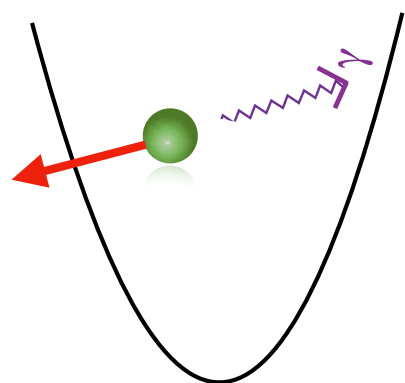
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Many-body H
Hermitian

Unavoidable atom losses:

One-body loss



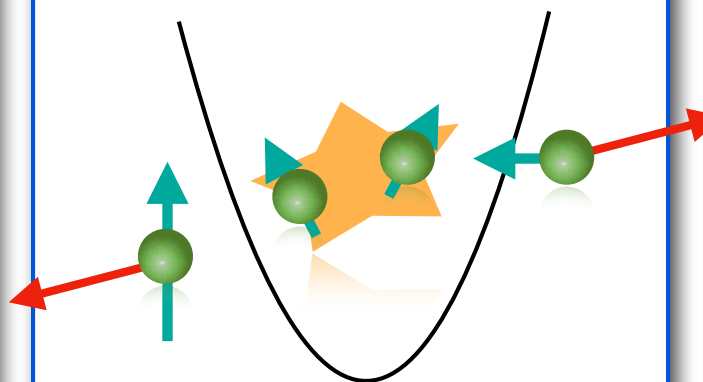
- vacuum

Here & Laser Cooling

- trap photon recoil

$$\dot{\Psi} = -\gamma \Psi$$

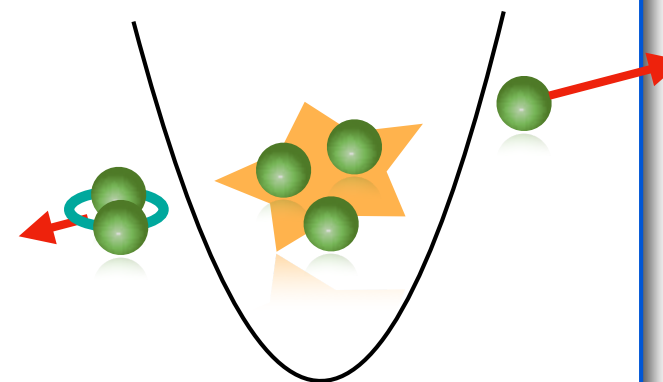
Two-body loss



- spin-flip collisions

$$\dot{\Psi} = -K_2 |\Psi|^2 \Psi$$

Three-body loss

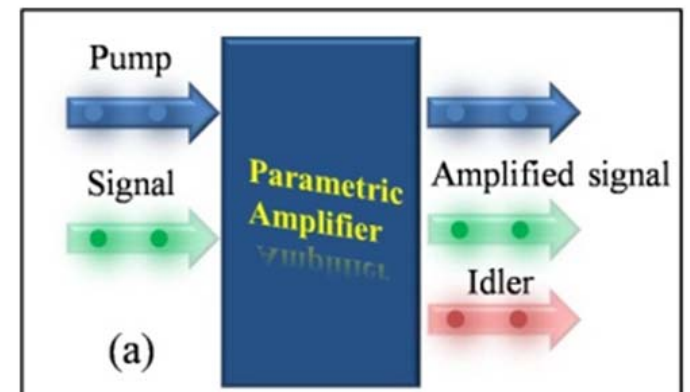
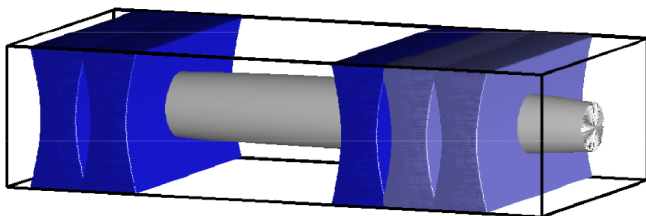


- molecule formation
- main density & life-time limitation

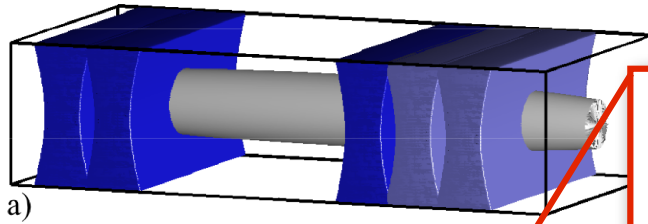
$$\dot{\Psi} = -K_3 \frac{|\Psi|^4}{2} \Psi$$

Outline

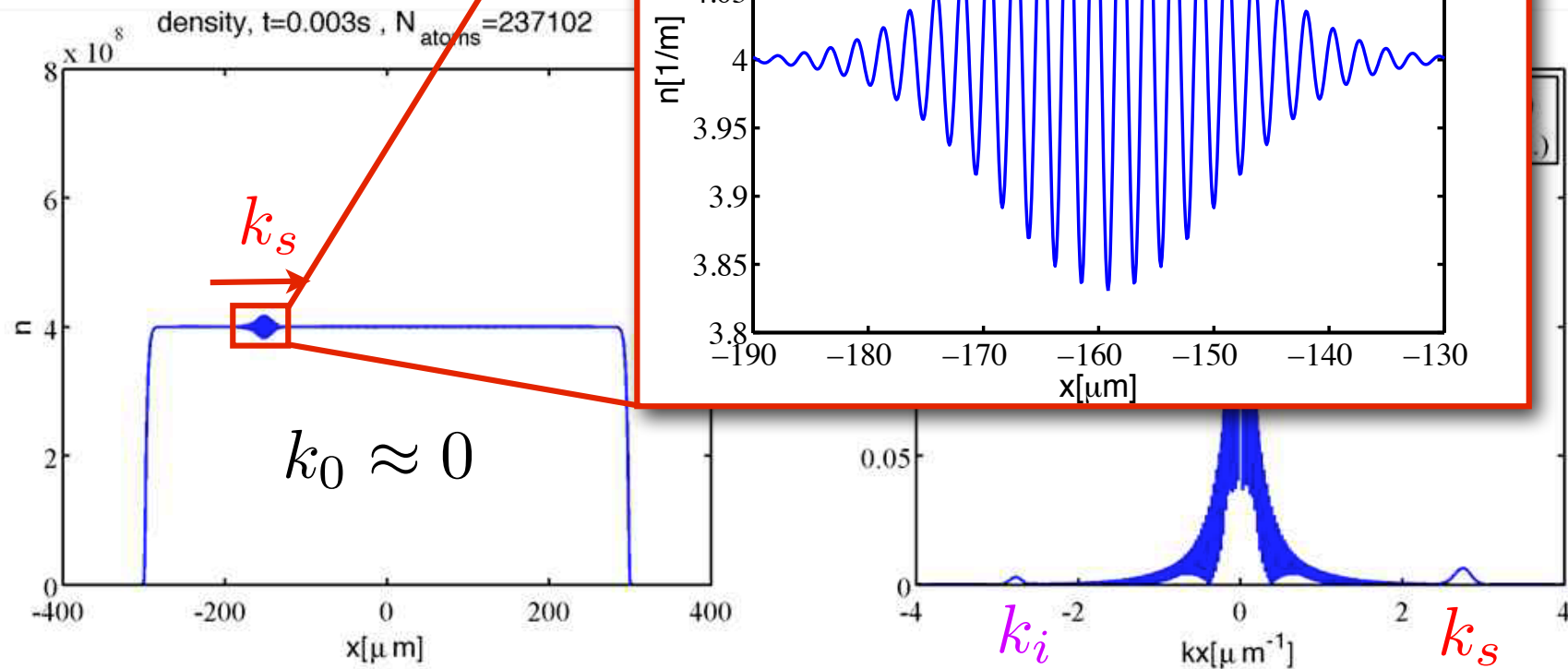
1. Motivation
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No wave mixing in homogeneous 1D condensate



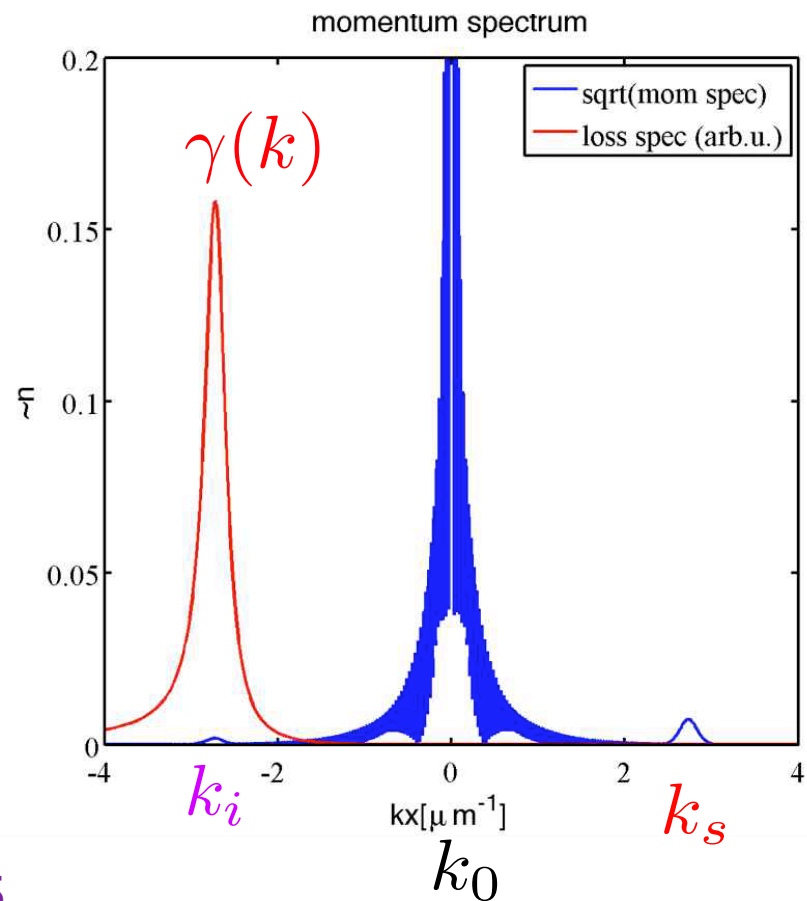
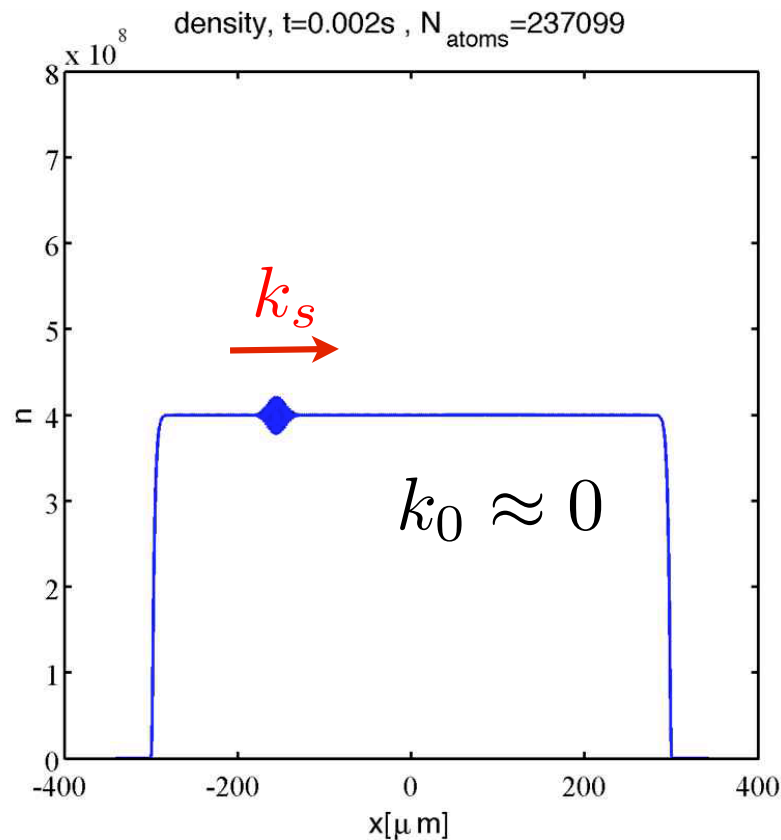
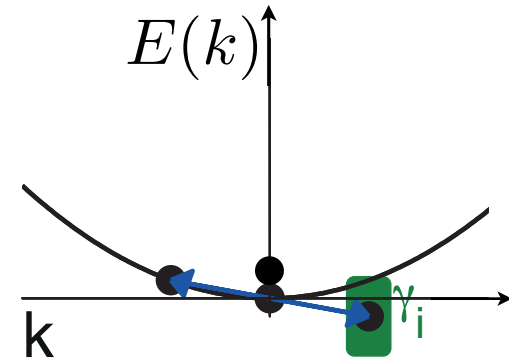
e.g. Meyerath et al. PRA 71 (2005) 0



Matter-wave amplification in homogeneous 1D condensate

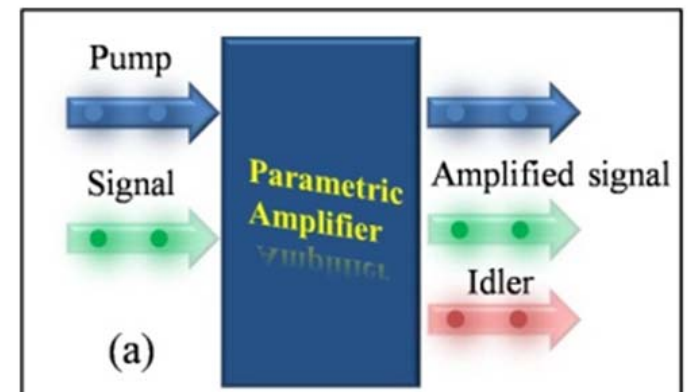
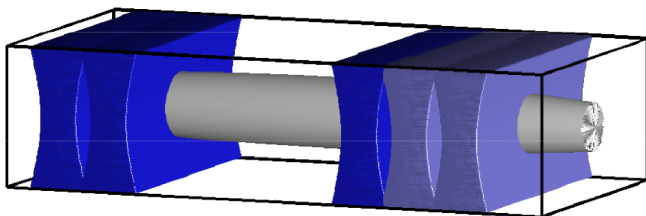
$$i\hbar \frac{\partial \phi(k)}{\partial t} = \frac{\hbar^2 k^2}{2m} \phi(k) + U \int_{-\infty}^{\infty} dk_1 \int_{-\infty}^{\infty} dk_2 \int_{-\infty}^{\infty} dk_3$$

$$\times \delta(k_2 + k_3 - k_1 - k) \phi^*(k_1) \phi(k_2) \phi(k_3) - i\hbar \gamma(k) \phi(k)$$



Outline

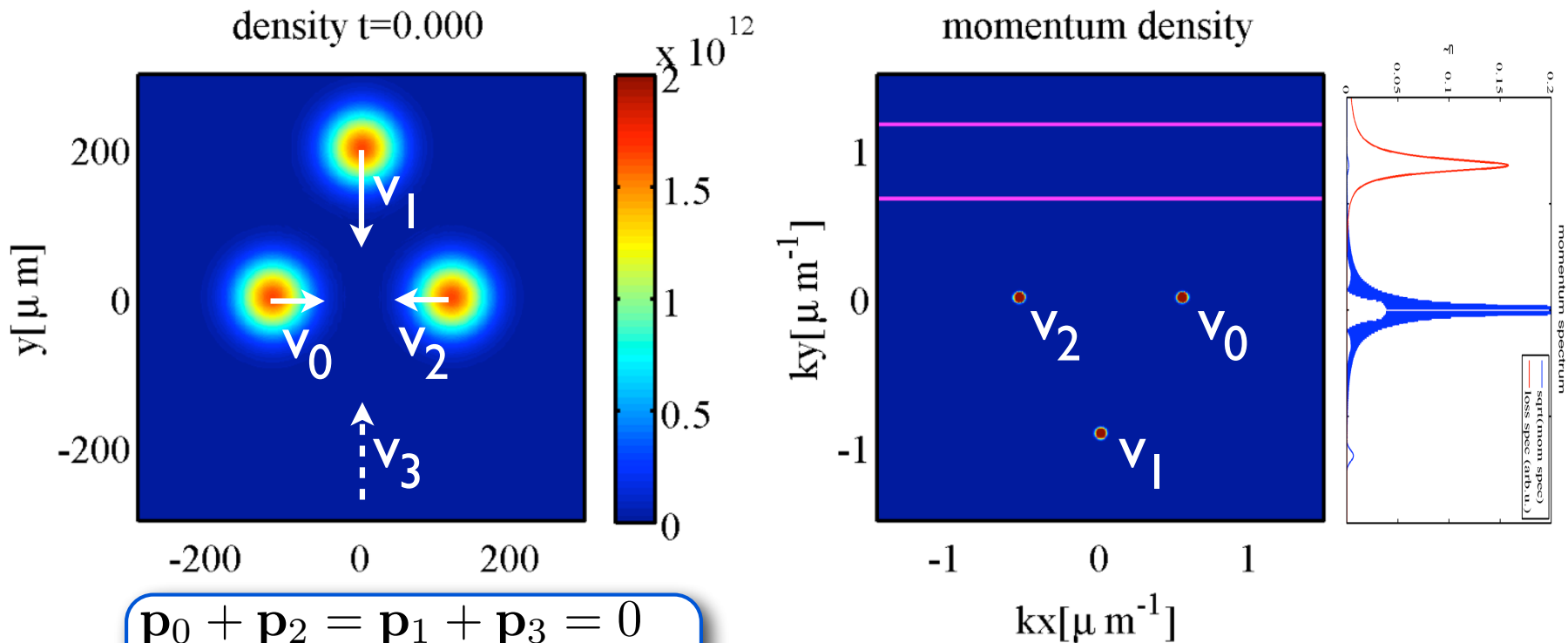
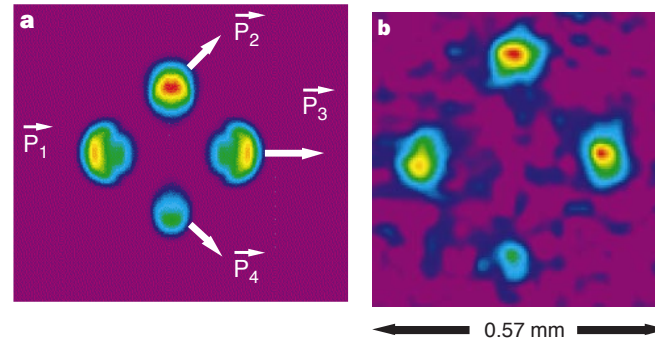
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Two dimensional wave mixing

L. Deng *et al.*
Nature **398**
(1999) 218.

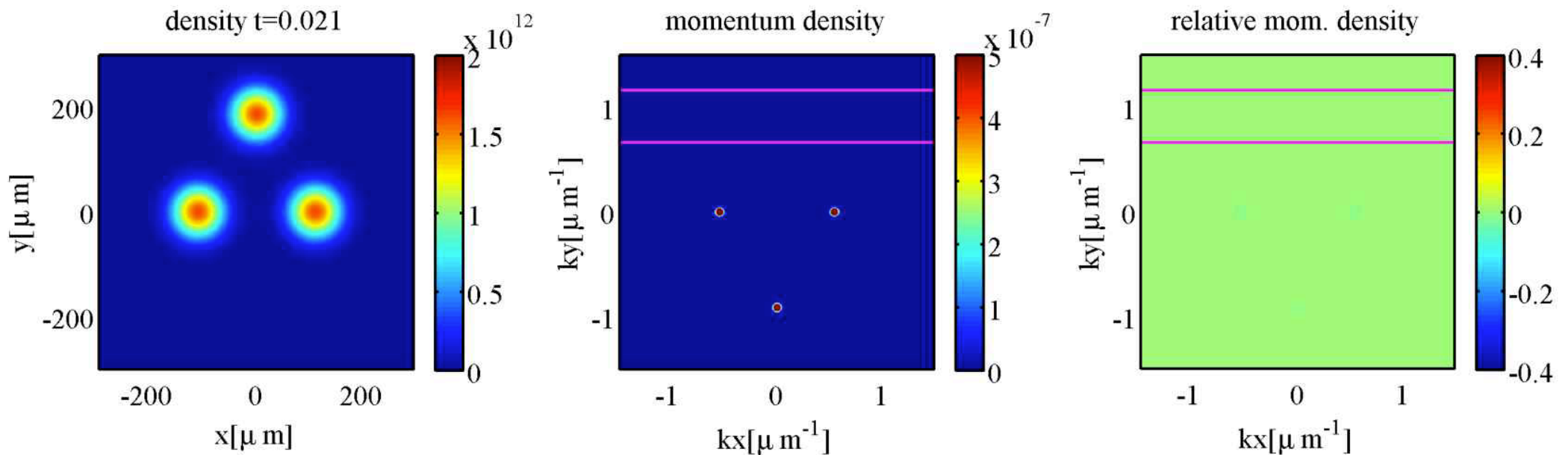
2D



$$\mathbf{p}_0 + \mathbf{p}_2 = \mathbf{p}_1 + \mathbf{p}_3 = 0$$

$$E_0 + E_2 \neq E_1 + E_3$$

Two dimensional wave mixing

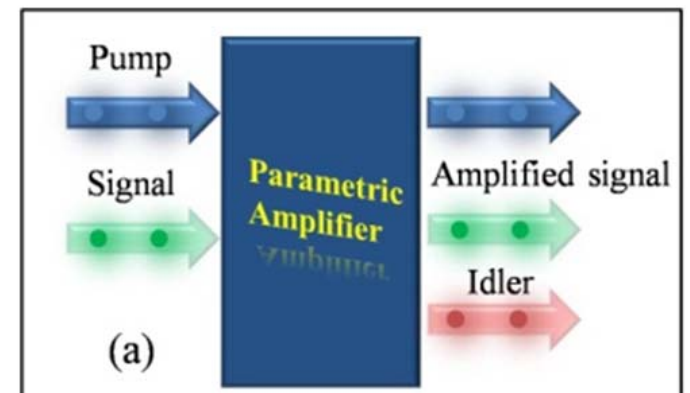
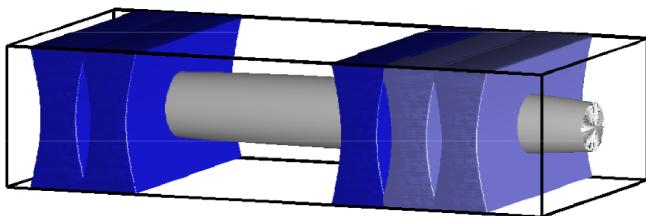


S. Wüster and R. El-Ganainy, PRA **96** (2017) 013605.

New channels in 1D and 2D, phonon amplifier!

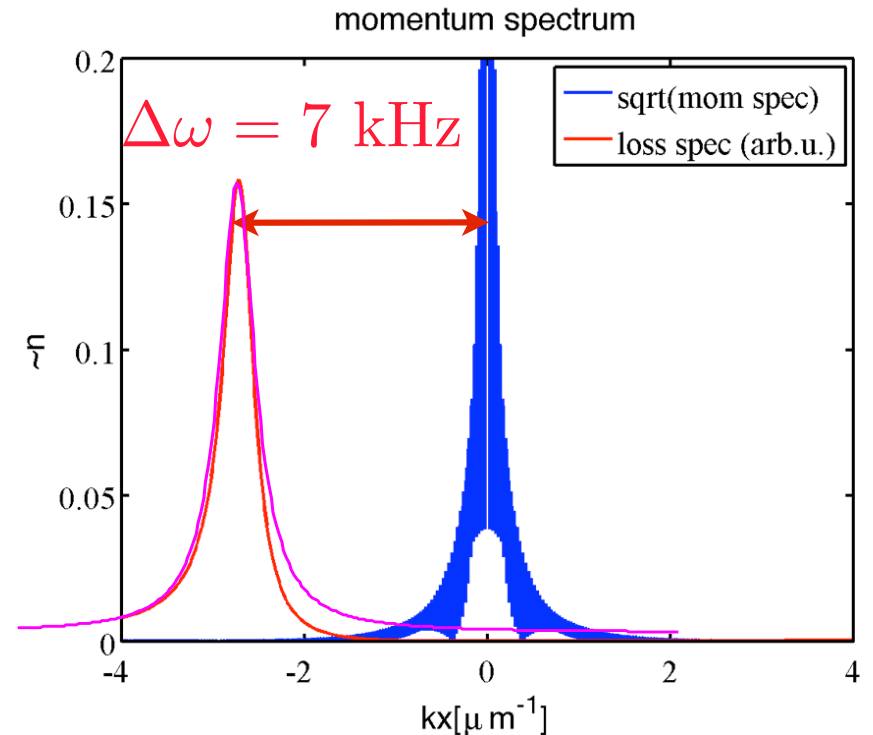
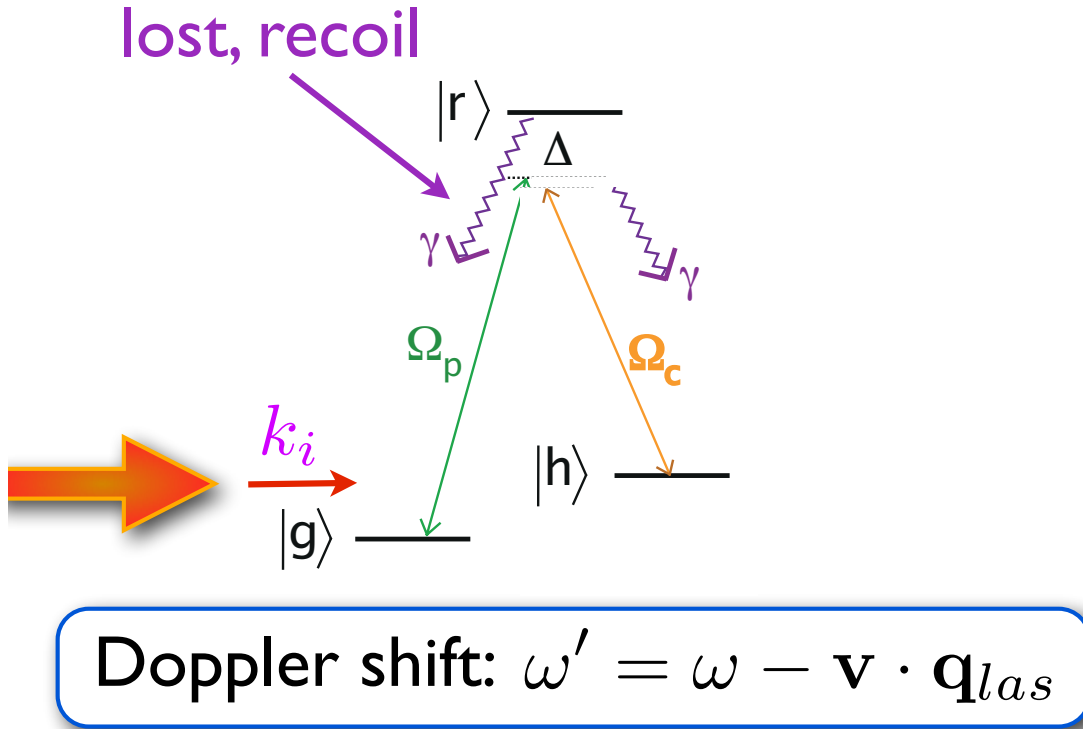
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Engineering atomic losses

Challenge: $\gamma(k)$ should act on *small range of k* only



(1) BEC velocities small, need relatively long-lived state $|r\rangle$
(metastable triplett or Rydberg $n=70, p$ $\gamma = 5 \text{ kHz}$)

(2) Tackle “tails” of Lorentzian: quantum interference, EIT

adaptation of [theory] G. Morigi *et al.* PRL **85** (2000) 4458. [expt] C. F. Roos *et al.* *ibid.* 5547.

Engineering atomic losses

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adaptation of [theory] G. Morigi *et al.* PRL **85** (2000) 4458. [expt] C. F. Roos *et al.* *ibid.* 5547.

Here:

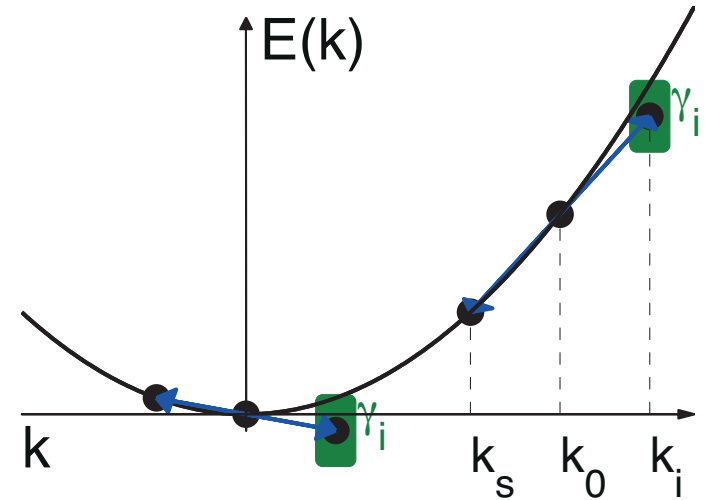
- Metastable excited states
- Decaying atoms lost from system
 - Weak trap
 - Hierarchy of time-scales: $T_{\text{BEC}} \gg T_{\text{loss}} \gg T_{\gamma}$

S. Wüster and R. El-Ganainy, PRA **96** (2017) 013605.

Few mode model

Analytical understanding of amplification mechanism

$$i\hbar \frac{\partial \phi_p}{\partial t} = U \sum_{n,m,l} \delta_{k_p+k_n-k_m-k_l} \times \phi_n^* \phi_m \phi_l e^{i(E_p+E_n-E_m-E_l)t/\hbar} - i\hbar \gamma_p \phi_p$$

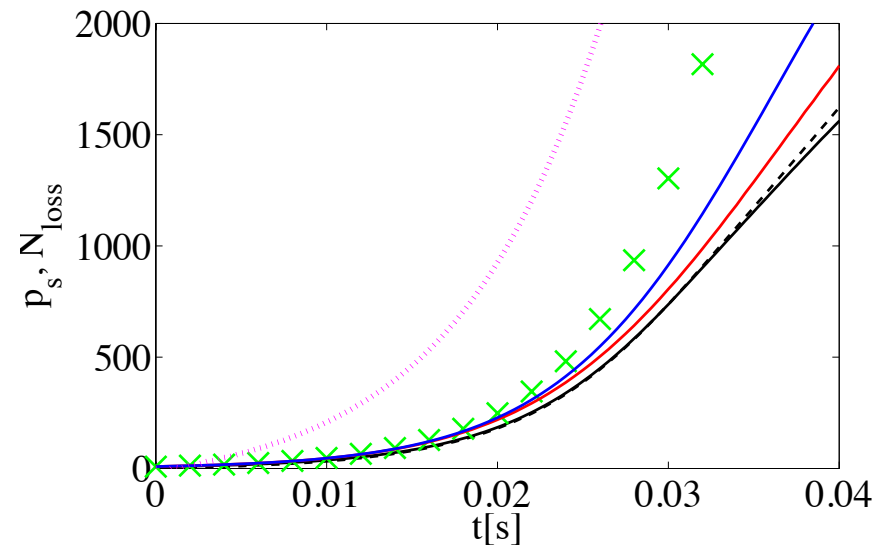


- Consider only three modes (0,s,i).
- Growth rates of solutions

$$\lambda'' = \frac{(U\rho/\hbar)^2 \gamma}{\gamma^2 + (\Delta E/\hbar)^2}$$

$$\Delta E = 2E_0 - E_s - E_i$$

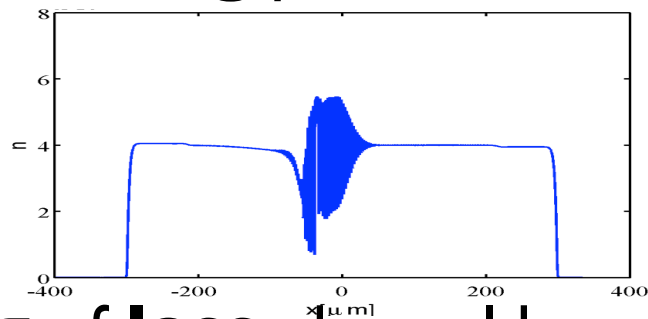
- Max for $\frac{\Delta E}{\hbar} = \gamma$, but “frequency” $\lambda'' \ll (U\rho/\hbar)$



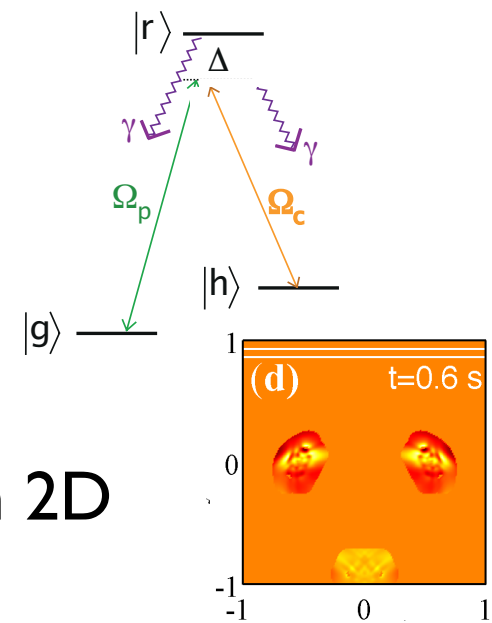
Summary

S. Wüster and R. El-Ganainy, PRA **96** (2017) 013605.

Spectrally selective **atom loss** in BEC **enables** otherwise forbidden **scattering** processes in **1D**



Engineering of **loss-channel** by adaptation of laser cooling techniques with **Rydberg** states

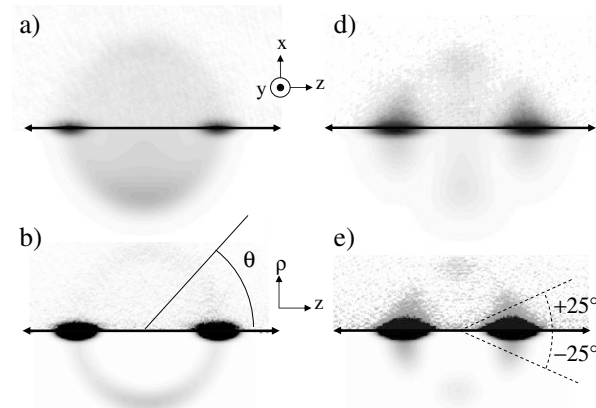


Simple studies of Non-Hermitian wave mixing in 2D

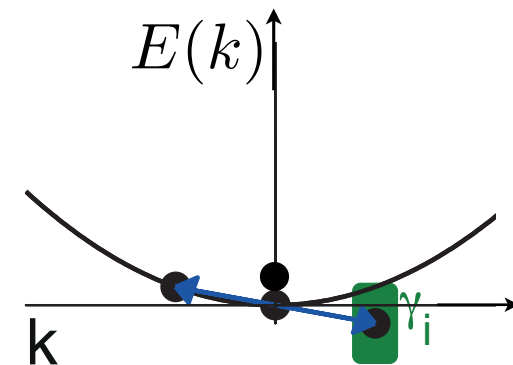
Outlook

Amplify scattering seeded by
vacuum fluctuations (TWA)?

Ch. Buggle et al. PRL **93** (2004) 173202.



Quantum correlations /
entanglement between pump
and signal (TWA)?



BEC naturally exhibit **non-linear**
loss (e.g. three-body recombination)

$$-i\hbar K_n |\Psi|^{2(n-1)} \Psi$$

Thanks for your attention