

# Quantum Many-Body Physics with Multimode Cavity QED

Jonathan Keeling



University of  
St Andrews

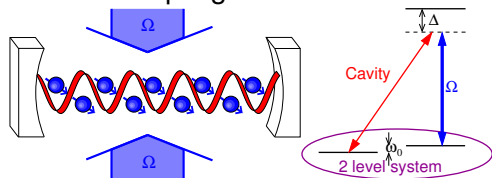
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1413



Open Quantum Systems, ICTS-TIFR, July 2017

# Synthetic cavity QED: Raman driving

- Tunable coupling via Raman



$$H_{\text{eff}} = \dots \frac{\Omega g}{\Delta} (\sigma_n^+ a + \text{H.c.})$$

• Real systems: loss  $\partial_t \rho = -i[H, \rho] + \kappa C[a, \rho] + \dots$

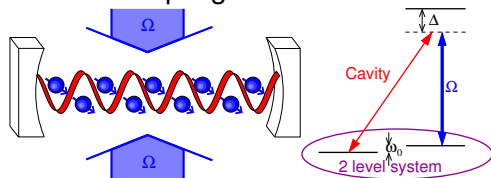
• To balance loss, counter-rotating:

$$H_{\text{eff}} = \dots \frac{\Omega g}{\Delta} \sigma_n^x (a + a^\dagger)$$

[Dimer *et al.* PRA '07]

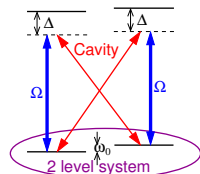
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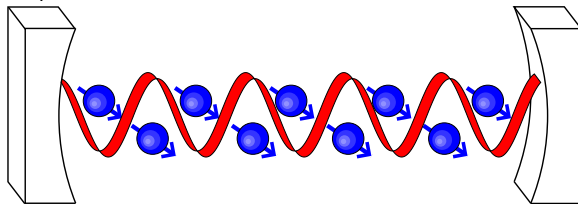
[Dimer *et al.* PRA '07]

# (Multimode) cavity QED

$$H = \sum_k \omega_k a_k^\dagger a_k + \sum_n \omega_0 \sigma_n^+ \sigma_n^- + \sum_{n,k} g_{k,n} (a_k^\dagger + a_{-k}) (\sigma_n^+ + \sigma_n^-)$$

$$\dot{\rho} = -i[H, \rho] + \kappa \sum_k \mathcal{L}[a_k, \rho] + \gamma \sum_i \mathcal{L}[\sigma_i^-, \rho]$$

• Compare  $g$  (or  $g\sqrt{N}$ ) vs:  
 $\kappa, \gamma$





# (Multimode) cavity QED

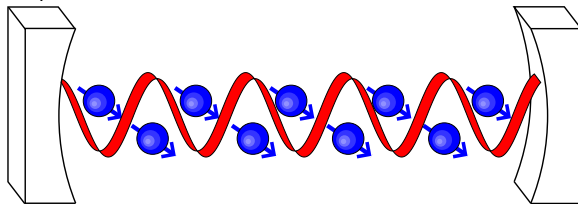
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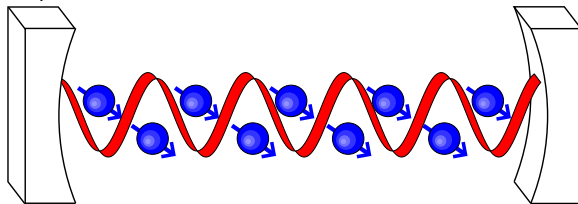
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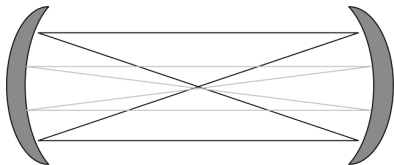
- Compare  $g$  (or  $g\sqrt{N}$ ) vs:

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- ▶ **bandwidth**
- ▶  $\omega_k, \omega_0$

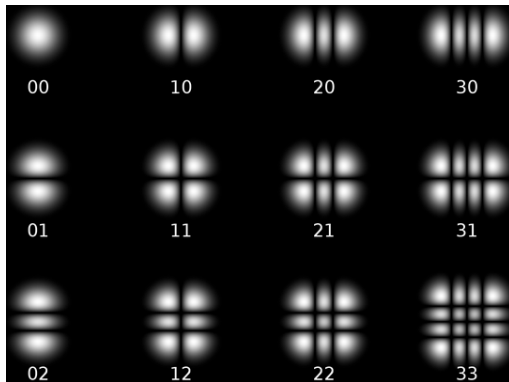
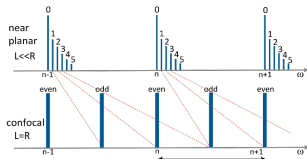


# Multimode cavities

Confocal cavity  $L = R$ :



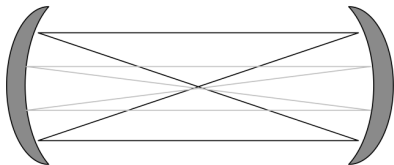
- Modes  $\Xi_{l,m}(\mathbf{r}) = H_l(x)H_m(y)$ ,  $l + m$  fixed parity



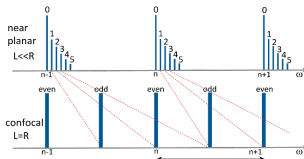
• Tune between  
degenerate vs non-degenerate

# Multimode cavities

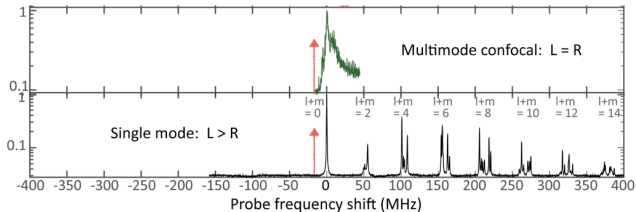
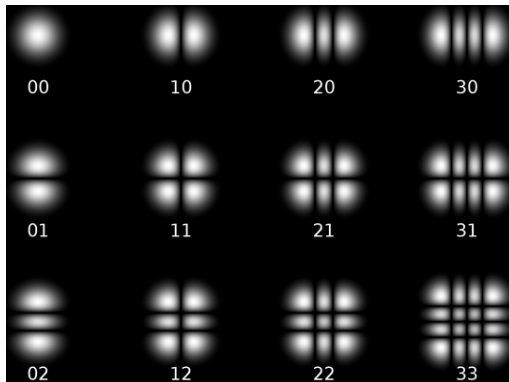
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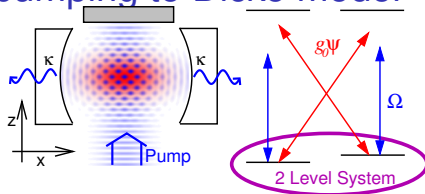


- Tune between degenerate vs non-degenerate



- 1 Introduction: Tunable multimode Cavity QED
- 2 Density wave polaritons
  - Superradiance transition
  - Supermode density wave polariton condensation
  - Degenerate limit
- 3 Spin wave polaritons
  - Spin glass
  - Effects of loss
- 4 Meissner-like effect

# Mapping transverse pumping to Dicke model



- Atomic states:  $\psi(\mathbf{r}) = \psi_{\downarrow} + \psi_{\uparrow} \cos(qx) \cos(qz)$

$$H_{\text{eff}} = \underbrace{(\omega_c - \omega_p)}_{-\Delta_c} a^\dagger a + \sum_n \frac{\omega_0}{2} \sigma_n^z + \underbrace{\frac{\Omega g_0}{\Delta}}_{g_{\text{eff}}} \sigma_n^x (a + a^\dagger)$$

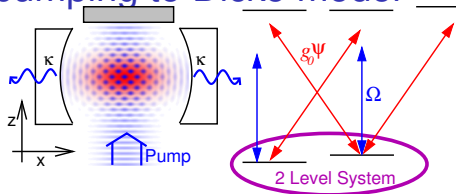
- Extra "feedback" term  $U$ , cavity loss  $\kappa$

$Z_2$  symmetry

- Phase of light

Atom checkerboard

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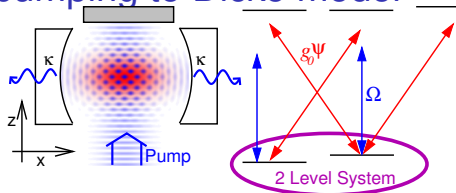
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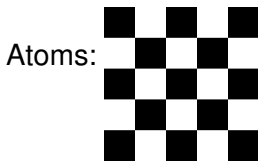
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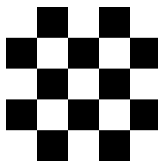
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- 1 Phase of light.
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vs





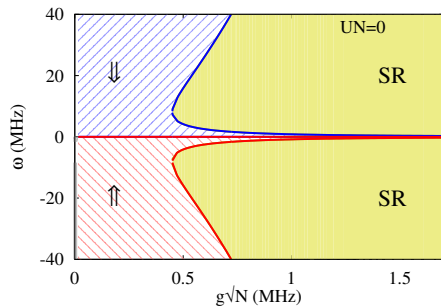
# Classical dynamics

Changing  $U$ :

$$U = 0$$

$$U < 0$$

$$U > 0$$



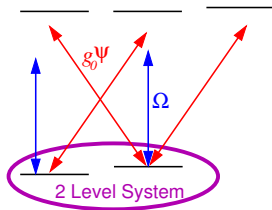
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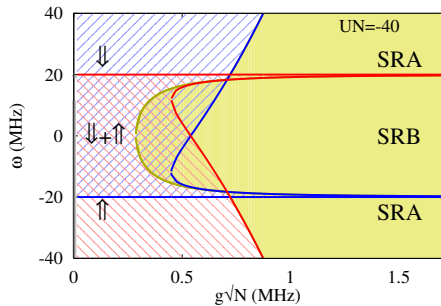
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$$U \propto \frac{g_0^2}{\omega_c - \omega_a}$$



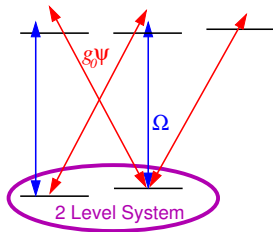
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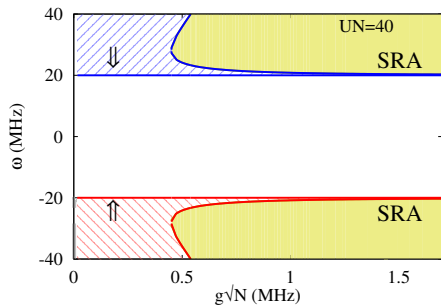
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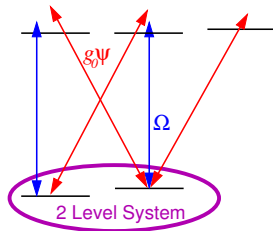
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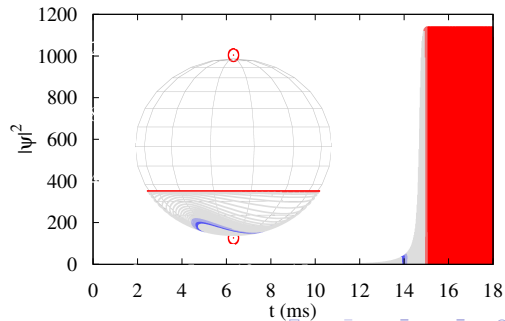
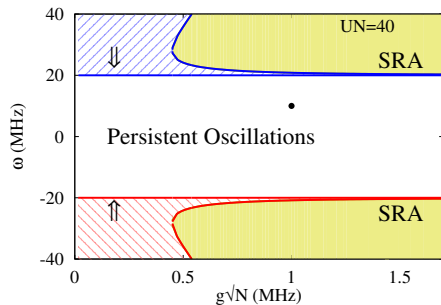
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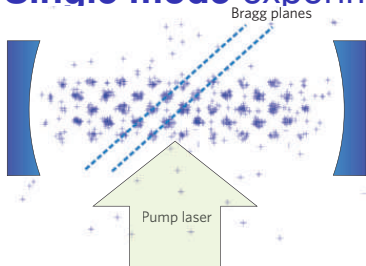
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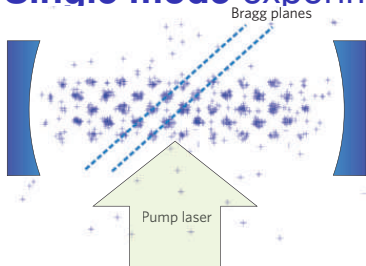


# Single mode experiments



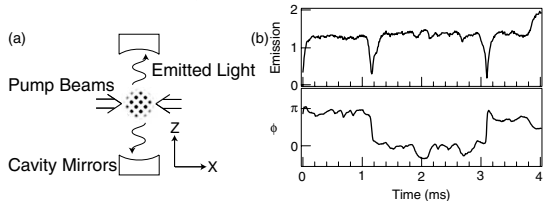
Ritsch *et al.* PRL '02

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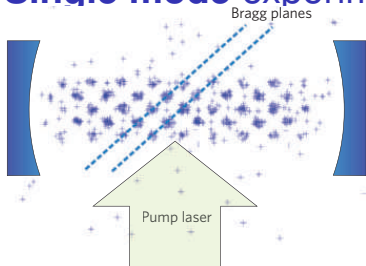
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## Thermal atoms, momentum state



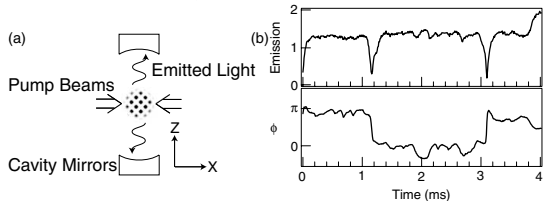
Vuletic *et al.* PRL '03 (MIT)

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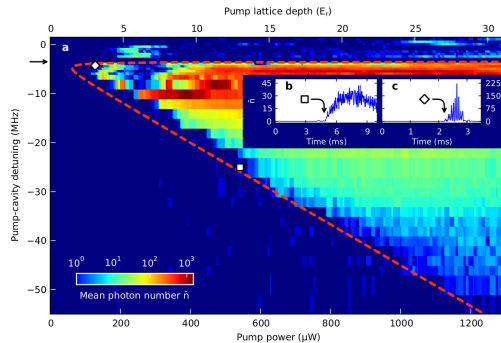
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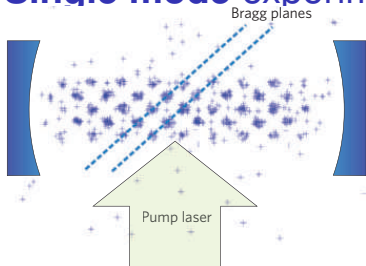
## BEC, momentum state



Baumann *et al.* Nature '10 (ETH)

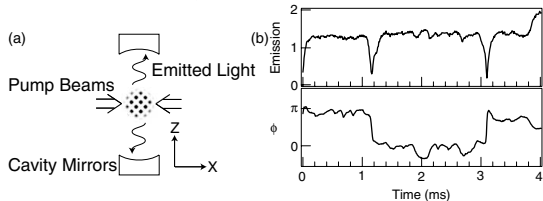
Kinder *et al.* PRL '15 (Hamburg)

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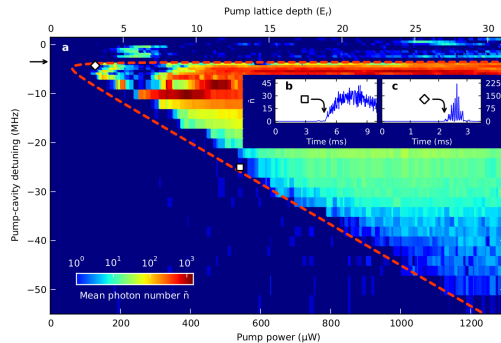
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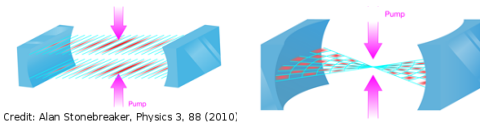
## BEC, hyperfine states

Baden *et al.* PRL '14 (Singapore)



# Synthetic cQED Possibilities

- Single mode vs **multimode**



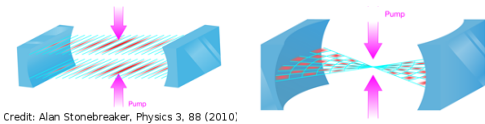
● Momentum state vs hyperfine state

● XY vs Ising

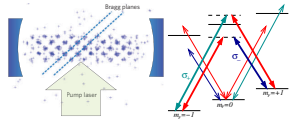
● Thermal gas vs BEC vs disorder localised

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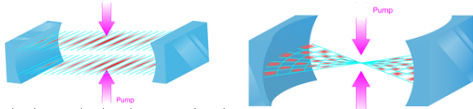


● XY vs Ising

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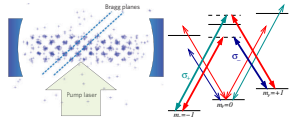
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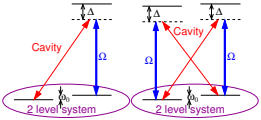


Credit: Alan Stonebreaker, Physics 3, 88 (2010)

- Momentum state vs hyperfine state



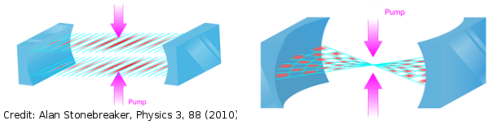
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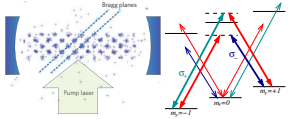
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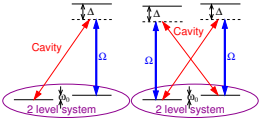
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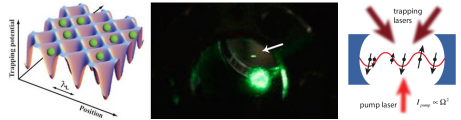
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# Density wave polaritons

1 Introduction: Tunable multimode Cavity QED

2 **Density wave polaritons**

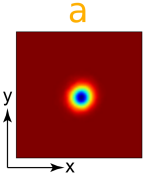
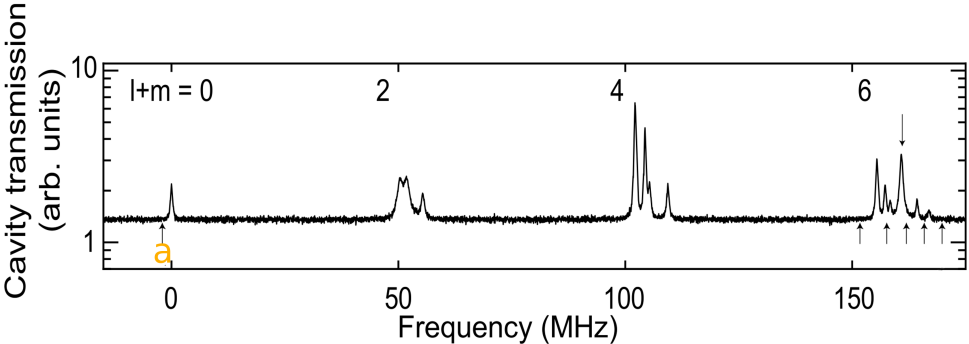
- Superradiance transition
- **Supermode density wave polariton condensation**
- Degenerate limit

3 Spin wave polaritons

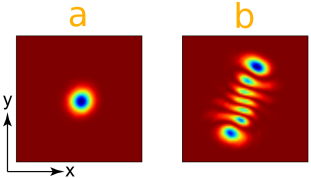
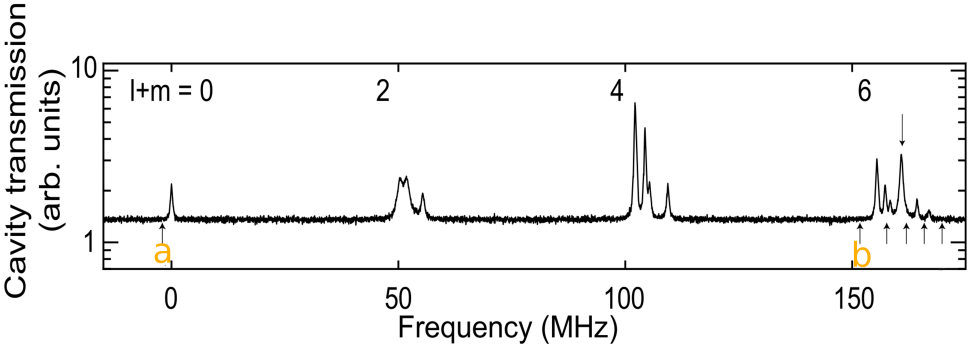
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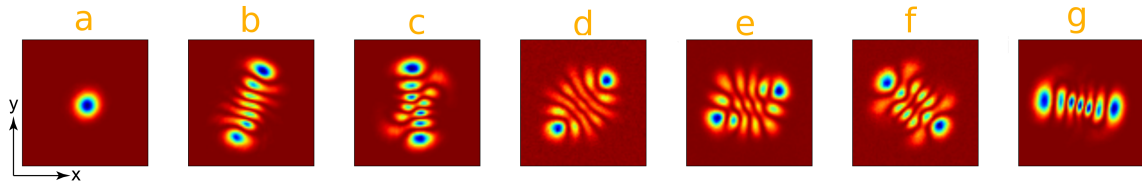
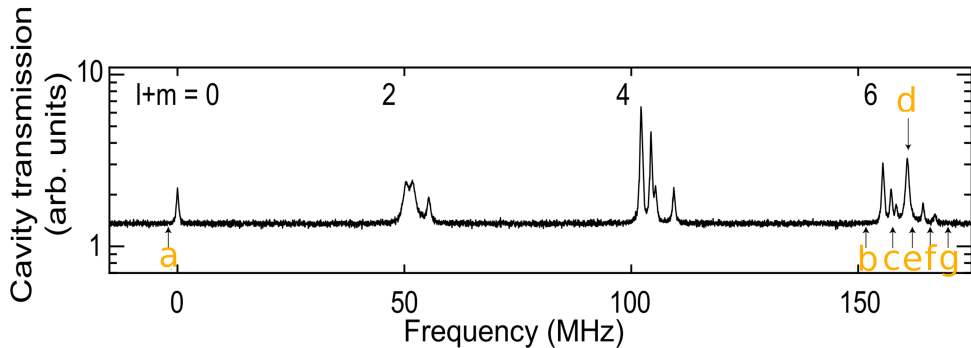
# Superradiance in multimode cavity: Even family



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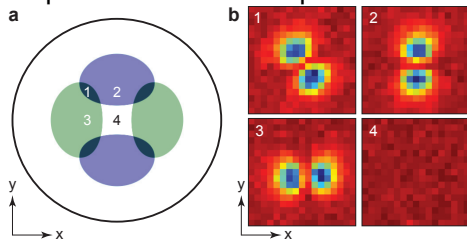




# Superradiance in multimode cavity: Odd family

• Atomic time-of-flight  $\rightarrow$  structure factor  
 $\psi(r) = \psi_0(r) + \psi_1(r) \cos(qy) \cos(qz)$

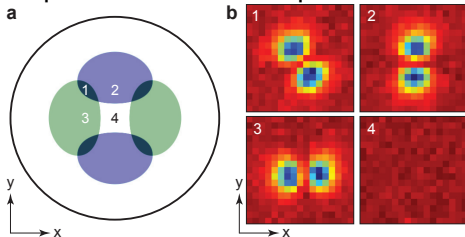
- Dependence on cloud position



- Near-degeneracy of (1, 0), (0, 1) modes broken by matter-light coupling.

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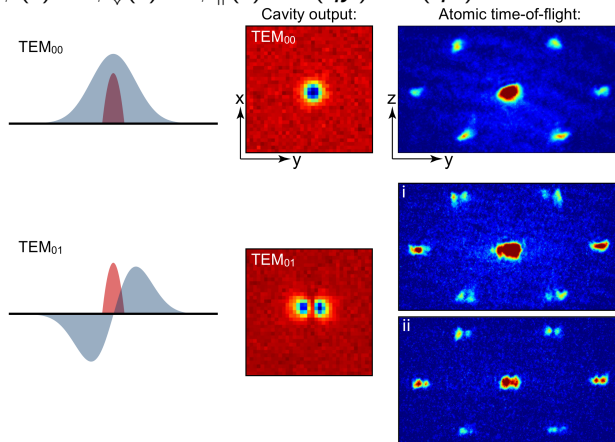
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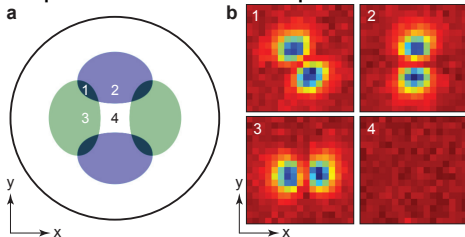
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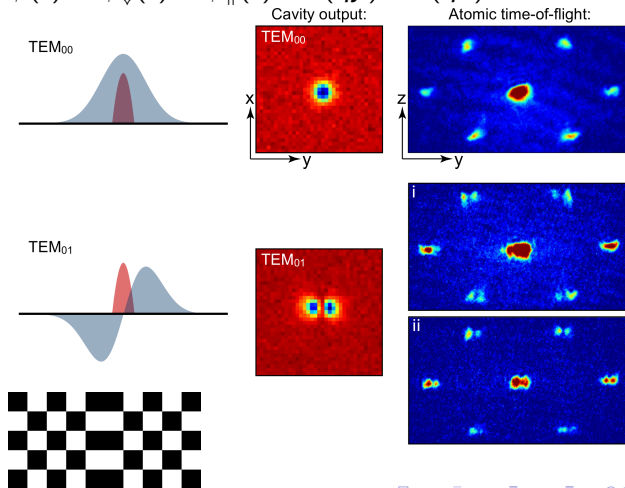
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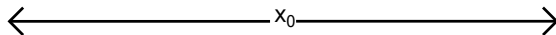
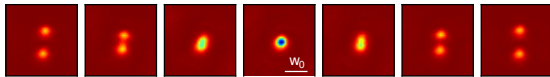
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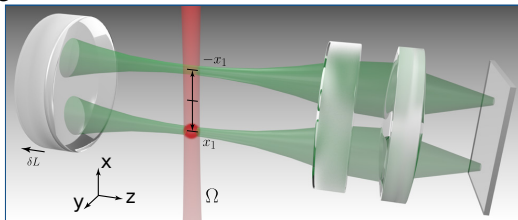
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# Degenerate cavity limit

- Multimode cavity — light follows atom cloud



- Confocal cavity: mirror image

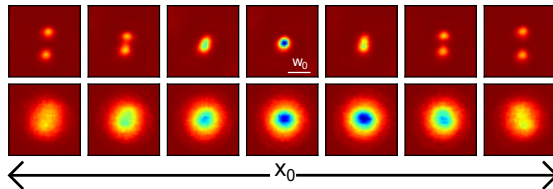


• Single mode cavity, Gaussian.

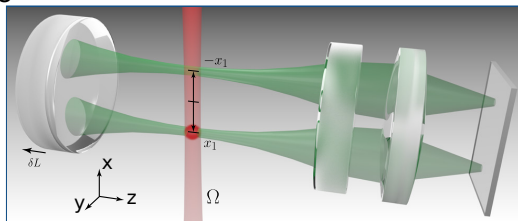
[Vaidya, Guo, Kroeze, Ballantine, Kollar, JK, Lev, in preparation]

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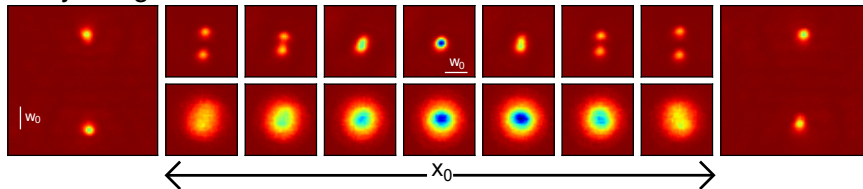


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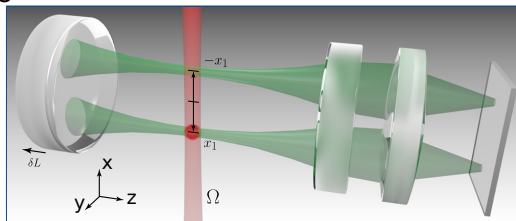
[Vaidya, Guo, Kroeze, Ballantine, Kollar, JK, Lev, in preparation]

# Degenerate cavity limit

- Multimode cavity — light follows atom cloud



- Confocal cavity: mirror image

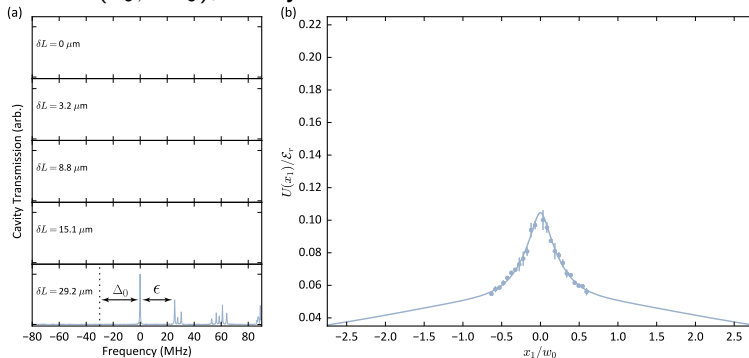


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# Measuring atom-image interaction

- Threshold measures  $U(x_0, -x_0)$ , cavity Green's function



- Small detuning  $\epsilon$ , short-range part  $U(x, -x) = u(x, -x) + u(x, -x)$ :

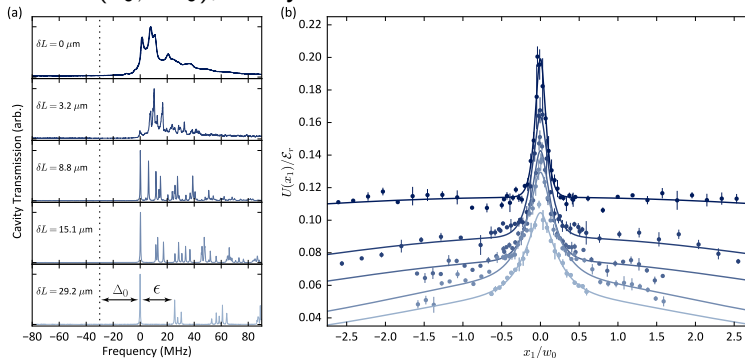
$$u(x, x') \propto \frac{M^* \exp(-2M^*|x-x'|)}{\sqrt{1 + ((x+x')/4M^*)^2}}, \quad M^* = \sqrt{\frac{\Delta_0}{2\epsilon}}$$

- Controllable range cavity-mediated interaction:



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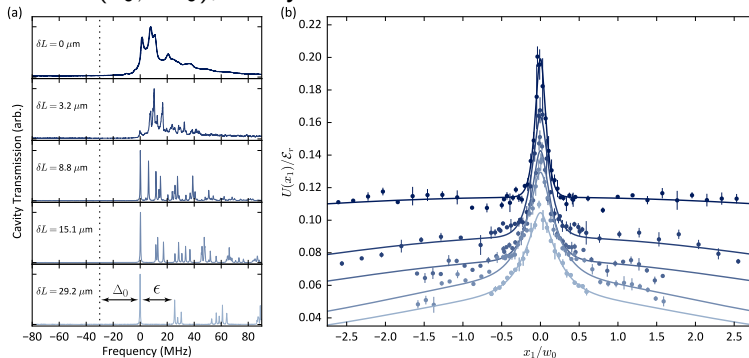
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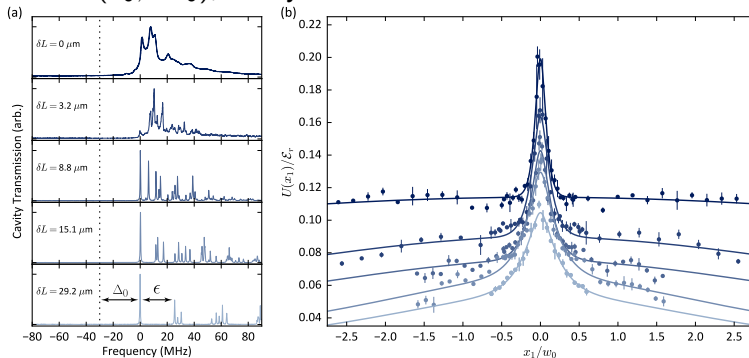


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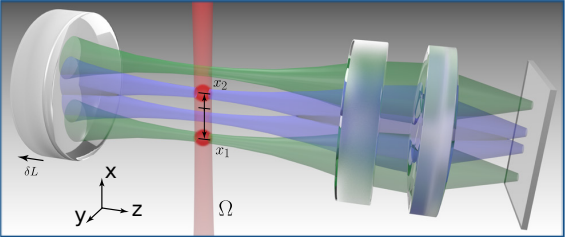
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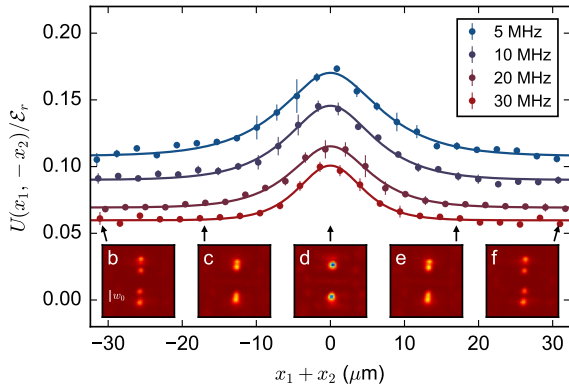
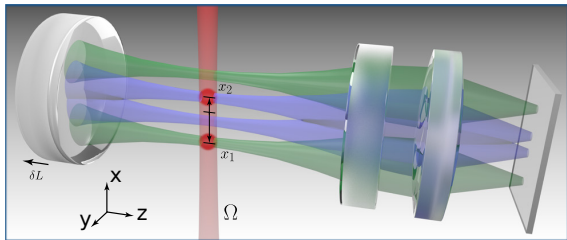
# Measuring atom-atom interaction

- Two BECs,



# Measuring atom-atom interaction

- Two BECs, vary  $x_1, x_2$



# Long-range part of interaction

- Atom-atom Green's function at  $\epsilon \rightarrow 0$ :

$$U(x, x') = \frac{1}{4} \left( \delta(x - x') + \delta(x + x') + \frac{1}{\pi} \cos(xx') \right)$$

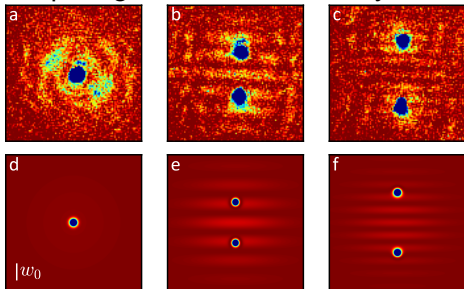
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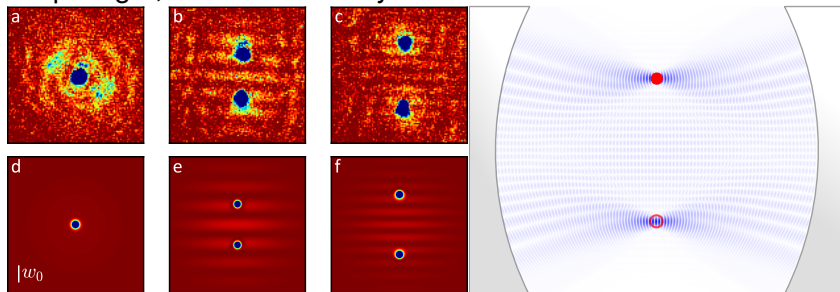


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# Spin wave polaritons

1 Introduction: Tunable multimode Cavity QED

2 Density wave polaritons

- Superradiance transition
- Supermode density wave polariton condensation
- Degenerate limit

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# Disordered atoms

- Multimode cavity, Hyperfine states,

$$H_{\text{eff}} = - \sum_{\mu} \Delta_{\mu} a_{\mu}^{\dagger} a_{\mu} + \sum_n \frac{\omega_0}{2} \sigma_n^z + \frac{\Omega g_0}{\Delta} \sum_{\mu} \Xi_{\mu}(\mathbf{r}_n) \sigma_n^x (a_{\mu} + a_{\mu}^{\dagger})$$

• Random atom positions – quenched disorder

• Effective XY/Ising spin glass

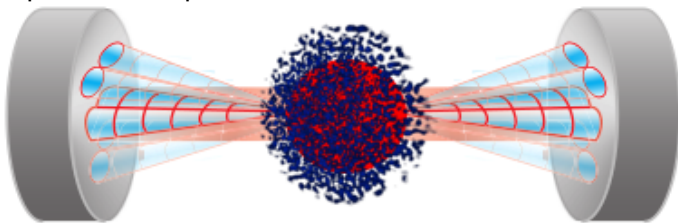
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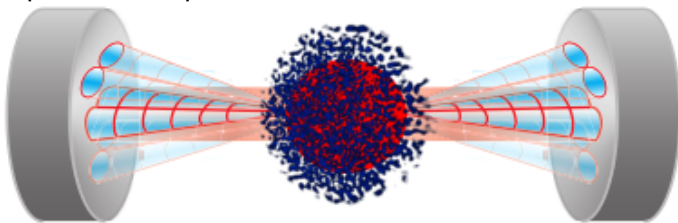
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## Internal states: Effect of particle losses

- Dicke Hamiltonian:  $H = \omega a^\dagger a + \sum_i \omega_0 \sigma_i^+ \sigma_i^- + g \sigma_i^x (a^\dagger + a)$
- Adding other loss terms

$$\partial_t \rho = -i[H, \rho] + \kappa \mathcal{L}[\hat{a}] + \sum_i \Gamma_\downarrow \mathcal{L}[\sigma_i^-] + \Gamma_\phi \mathcal{L}[\sigma_i^z]$$

$$\mathcal{L}[X] = X\rho X^\dagger - (X^\dagger X\rho + \rho X^\dagger X)/2$$

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[Dalla Torre *et al.*, PRA (Rapid) 2016, Kirton & JK, PRL 2017]

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# Effect of particle losses

- Wigner function  $W(\hat{a} = \hat{x} + i\hat{p})$

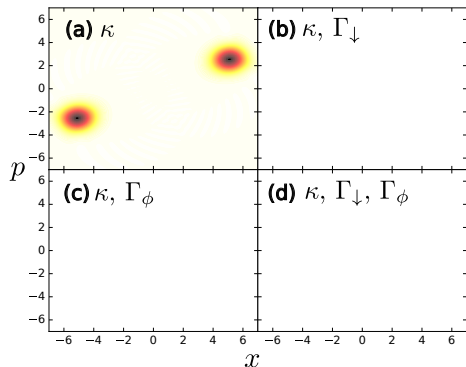
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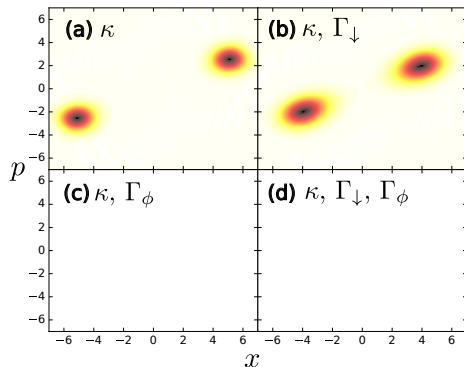
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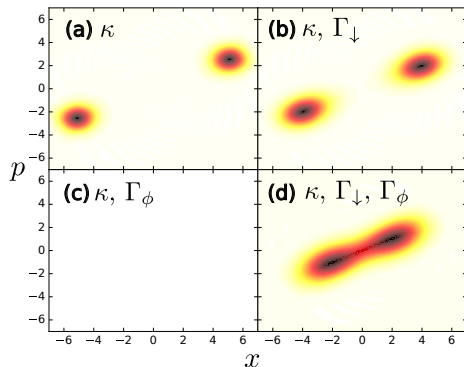
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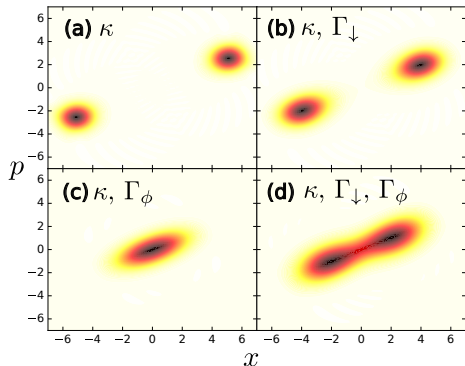
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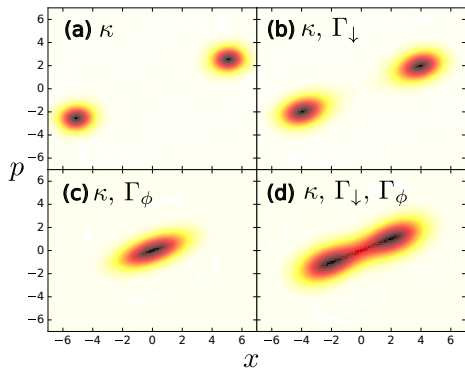
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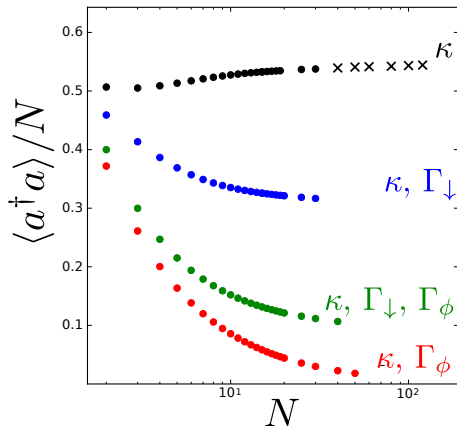
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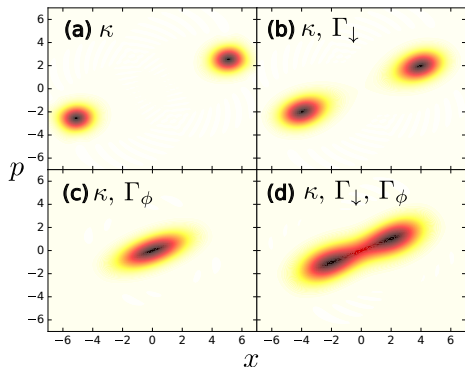
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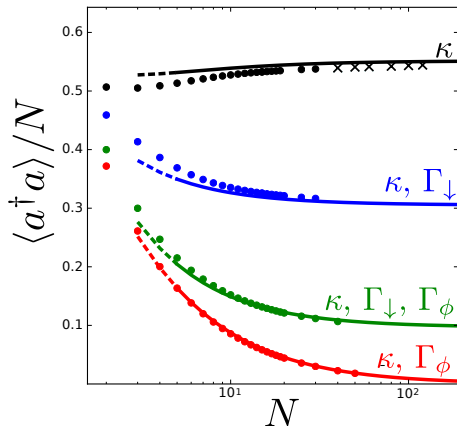
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# Meissner-like effect

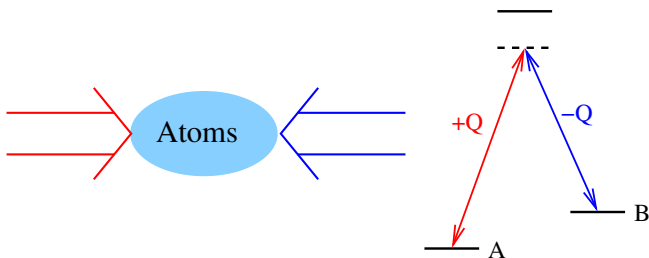
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# Cavity QED and synthetic gauge fields

- [Spielman, PRA '09] scheme, hyperfine states  $A, B$

$$H = (\psi_A \quad \psi_B) \begin{pmatrix} E_A + (\nabla - Q\hat{x})^2 & \\ & \Omega/2 \end{pmatrix} \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix} \begin{pmatrix} \\ E_B + (\nabla + Q\hat{x})^2 \end{pmatrix}$$



• Feedback

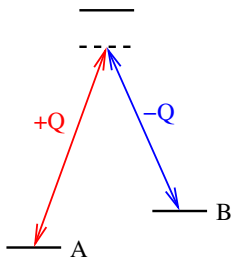
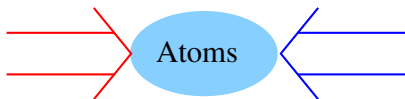
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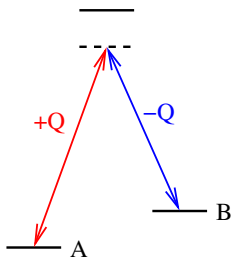
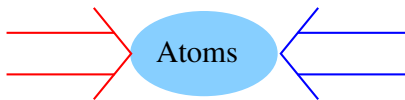
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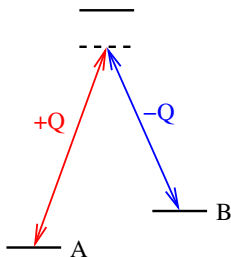
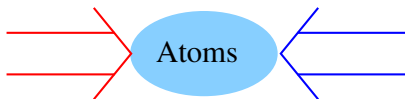
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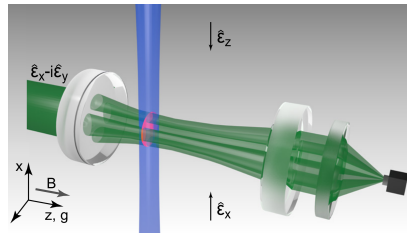
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- Follow Spielman scheme

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- $E_A, E_B \propto |v|^2$  from cavity Stark shift
- Ground state  $E_-(\mathbf{k}) \propto (\mathbf{k} - Q\hat{x}|v|^2)^2$

- Multimode cQED  $\rightarrow$  local matter-light coupling
- Variable profile synthetic gauge field?
- Reciprocity: matter affects field



[Ballantine *et al.* PRL 2017]

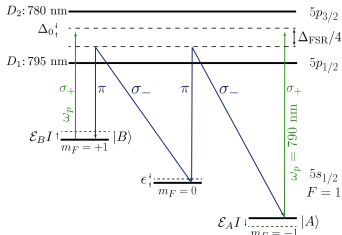
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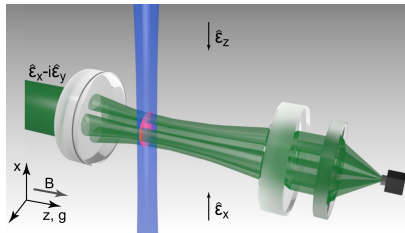
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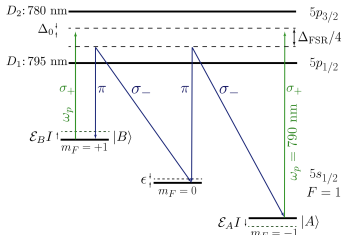
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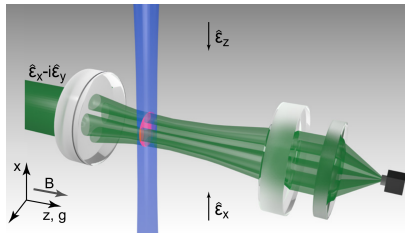
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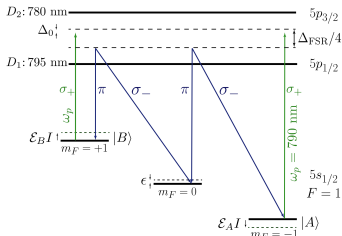
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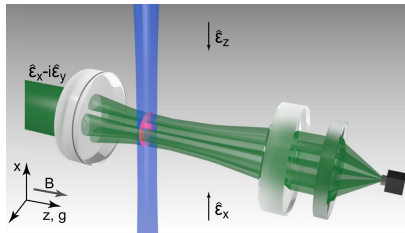
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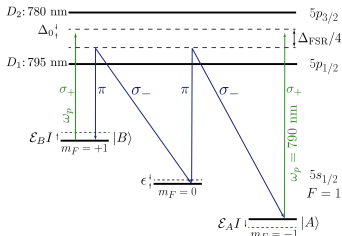
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- Follow Spielman scheme

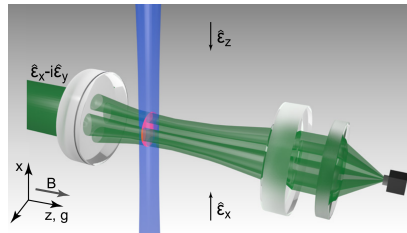
$$\begin{pmatrix} E_A + (\nabla - Q\hat{x})^2 & \Omega/2 \\ \Omega/2 & E_B + (\nabla + Q\hat{x})^2 \end{pmatrix}$$

- $E_A, E_B \propto |\varphi|^2$  from **cavity** Stark shift

- Ground state  $E_-(\mathbf{k}) \propto (\mathbf{k} - Q\hat{x}|\varphi|^2)^2$

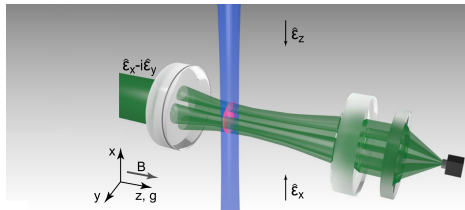


[Ballantine *et al.* PRL 2017]



- ▶ Multimode cQED  $\rightarrow$  local matter-light coupling
- ▶ Variable profile synthetic gauge field?
- ▶ Reciprocity: matter affects field

# Meissner-like physics: setup



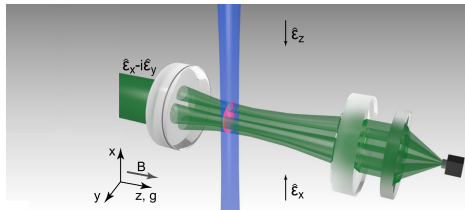
• Atoms: 
$$i\partial_t \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix} = \left[ -\frac{\nabla^2}{2m} + \begin{pmatrix} -\mathcal{E}_\Delta |\varphi|^2 + i\frac{q}{m}\partial_x & \Omega/2 \\ \Omega/2 & \mathcal{E}_\Delta |\varphi|^2 - i\frac{q}{m}\partial_x \end{pmatrix} + \dots \right] \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix}.$$

• Light: 
$$i\partial_t \varphi = \left[ \frac{\delta}{2} \left( -\nabla^2 + \frac{1}{R^2} \right) - \Delta_0 - i\kappa - N\mathcal{E}_\Delta (|\psi_A|^2 - |\psi_B|^2) \right] \varphi$$

• Low energy: 
$$|\psi_A|^2 - |\psi_B|^2 = \frac{Q}{im\Omega} (\psi_-^* \partial_x \psi_- - \psi_- \partial_x \psi_-^*) + 2\frac{\mathcal{E}_\Delta}{\Omega} |\psi_-|^2 |\varphi|^2$$

[Ballantine *et al.* PRL 2017]

# Meissner-like physics: setup

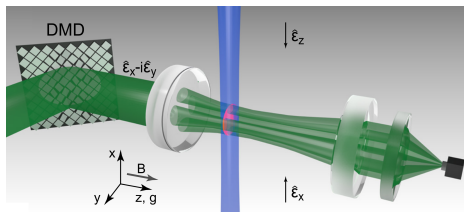


- Atoms: 
$$i\partial_t \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix} = \left[ -\frac{\nabla^2}{2m} + \begin{pmatrix} -\mathcal{E}_\Delta |\varphi|^2 + i\frac{q}{m}\partial_x & \Omega/2 \\ \Omega/2 & \mathcal{E}_\Delta |\varphi|^2 - i\frac{q}{m}\partial_x \end{pmatrix} + \dots \right] \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix} .$$
- Light: 
$$i\partial_t \varphi = \left[ \frac{\delta}{2} \left( -f^2 \nabla^2 + \frac{r^2}{f^2} \right) - \Delta_0 - i\kappa - N\mathcal{E}_\Delta (|\psi_A|^2 - |\psi_B|^2) \right] \varphi .$$

• Low energy:  $|\psi_A|^2 - |\psi_B|^2 = \frac{q}{im\Omega} (\psi_-^* \partial_x \psi_- - \psi_- \partial_x \psi_-^*) + 2\frac{\mathcal{E}_\Delta}{\Omega} |\psi_-|^2 |\varphi|^2$

[Ballantine *et al.* PRL 2017]

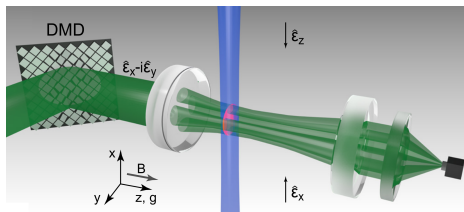
# Meissner-like physics: setup



- Atoms: 
$$i\partial_t \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix} = \left[ -\frac{\nabla^2}{2m} + \begin{pmatrix} -\mathcal{E}_\Delta |\varphi|^2 + i\frac{q}{m}\partial_x & \Omega/2 \\ \Omega/2 & \mathcal{E}_\Delta |\varphi|^2 - i\frac{q}{m}\partial_x \end{pmatrix} + \dots \right] \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix}.$$
- Light: 
$$i\partial_t \varphi = \left[ \frac{\delta}{2} \left( -l^2 \nabla^2 + \frac{r^2}{l^2} \right) - \Delta_0 - i\kappa - N\mathcal{E}_\Delta (|\psi_A|^2 - |\psi_B|^2) \right] \varphi.$$
- Low energy: 
$$|\psi_A|^2 - |\psi_B|^2 = \frac{Q}{im\Omega} (\psi_-^* \partial_x \psi_- - \psi_- \partial_x \psi_-^*) + 2\frac{\mathcal{E}_\Delta}{\Omega} |\psi_-|^2 |\varphi|^2$$

[Ballantine *et al.* PRL 2017]

# Meissner-like physics: setup

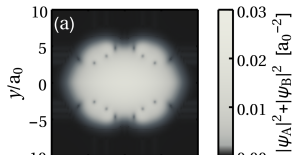


- Atoms: 
$$i\partial_t \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix} = \left[ -\frac{\nabla^2}{2m} + \begin{pmatrix} -\mathcal{E}_\Delta |\varphi|^2 + i\frac{q}{m}\partial_x & \Omega/2 \\ \Omega/2 & \mathcal{E}_\Delta |\varphi|^2 - i\frac{q}{m}\partial_x \end{pmatrix} + \dots \right] \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix}.$$
- Light: 
$$i\partial_t \varphi = \left[ \frac{\delta}{2} \left( -\ell^2 \nabla^2 + \frac{r^2}{\ell^2} \right) - \Delta_0 - i\kappa - N\mathcal{E}_\Delta (|\psi_A|^2 - |\psi_B|^2) \right] \varphi + f(\mathbf{r}).$$
- Low energy: 
$$|\psi_A|^2 - |\psi_B|^2 = \frac{Q}{im\Omega} (\psi_-^* \partial_x \psi_- - \psi_- \partial_x \psi_-^*) + 2\frac{\mathcal{E}_\Delta}{\Omega} |\psi_-|^2 |\varphi|^2$$

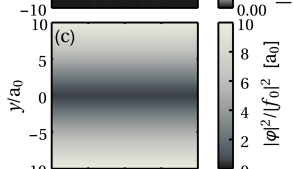
[Ballantine *et al.* PRL 2017]

# Meissner-like physics: numerical simulations

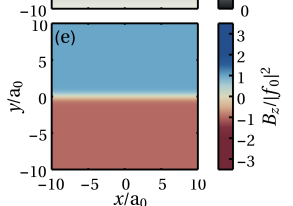
Atoms



Cavity light



Synthetic field

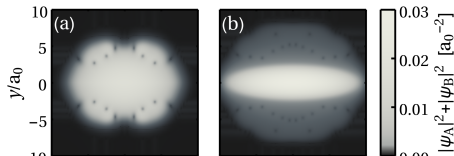


- Consider  $f(\mathbf{r})$  such that  $|\varphi|^2 \propto y$ .
- Without feedback ( $\mathcal{E}_\Delta = 0$ ) for field
- With feedback

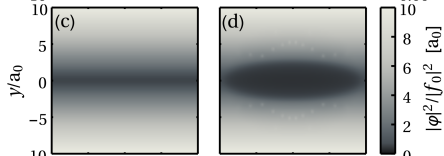
[Ballantine *et al.* PRL 2017]

# Meissner-like physics: numerical simulations

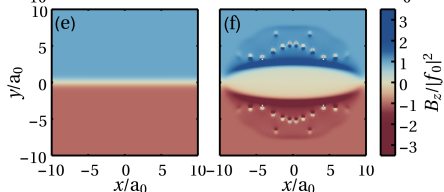
Atoms



Cavity light



Synthetic field



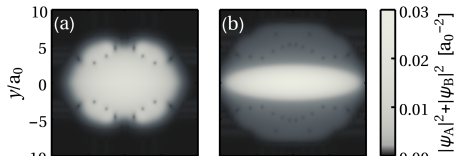
- Consider  $f(\mathbf{r})$  such that  $|\phi|^2 \propto y$ .
- Without feedback ( $\mathcal{E}_\Delta = 0$ ) for field
- With feedback

• Field expelled  
• Cloud shrinks

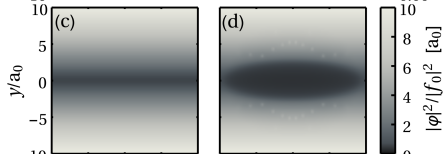
[Ballantine *et al.* PRL 2017]

# Meissner-like physics: numerical simulations

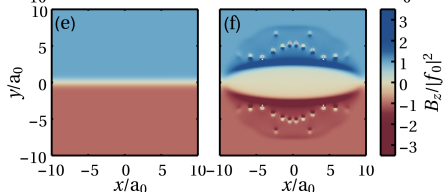
Atoms



Cavity light



Synthetic field



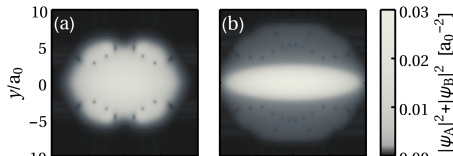
- Consider  $f(\mathbf{r})$  such that  $|\phi|^2 \propto y$ .
- Without feedback ( $\mathcal{E}_\Delta = 0$ ) for field
- With feedback
  - ▶ Field expelled
  - ▶ Cloud shrinks

[Ballantine *et al.* PRL 2017]

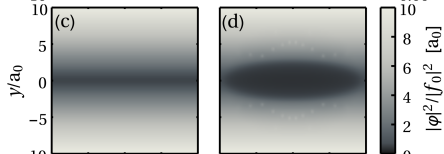


# Meissner-like physics: numerical simulations

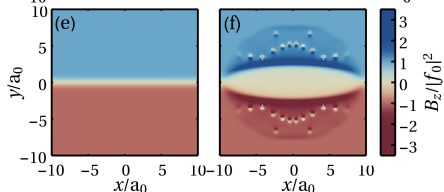
Atoms



Cavity light



Synthetic field

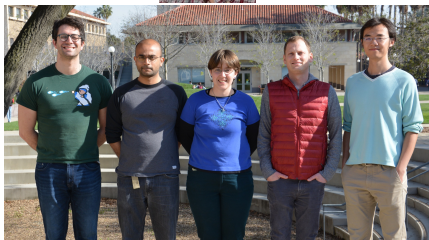


- Consider  $f(\mathbf{r})$  such that  $|\phi|^2 \propto y$ .
- Without feedback ( $\mathcal{E}_\Delta = 0$ ) for field
- With feedback
  - ▶ Field expelled
  - ▶ Cloud shrinks

[Ballantine *et al.* PRL 2017]

# Acknowledgments

Experiment (Stanford):  
Benjamin Lev



Theory:



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**Peter Kirton, Kyle Ballantine, Laura Staffini** (St Andrews)



The Leverhulme Trust

EPSRC

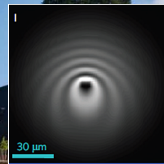
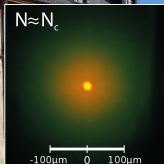
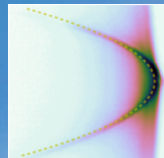
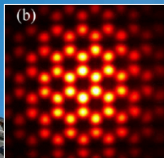
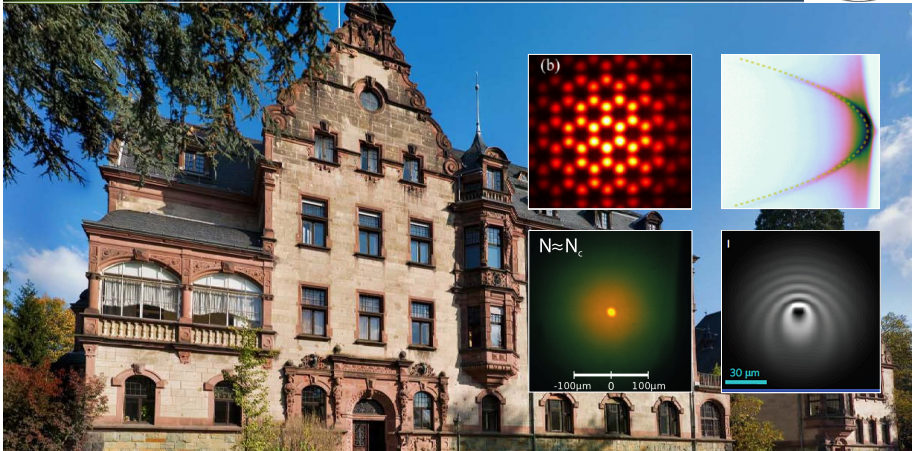
Engineering and Physical Sciences  
Research Council



Topological Protection and  
Non-Equilibrium States in  
Strongly Correlated Electron  
Systems

# WE-Heraeus-Seminar: Condensates of Light

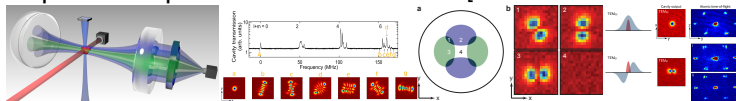
Physikzentrum Bad Honnef, Germany



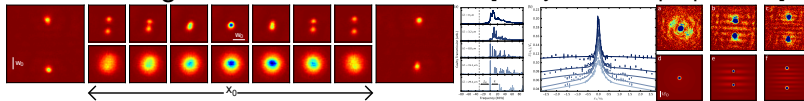
**14th – 17th JANUARY 2018**

# Summary

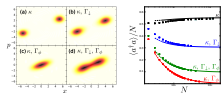
- Supermode polariton condensation [Kollár *et al.* Nat. Comms. 2017]



- Variable range atom-atom interaction [Vaidya *et al.* in preparation]



- Open Dicke model,  $\kappa$ ,  $\Gamma_\phi$ ,  $\Gamma_\downarrow$  [Kirton & JK, PRL 2017]



- Meissner like effect [Ballantine *et al.* PRL 2017]

