

# Dephasing Enabled Fast Charging of Quantum Batteries

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# Where do I come from?



# Acknowledgements



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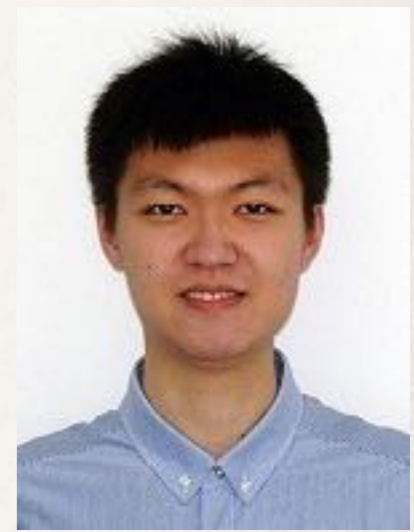
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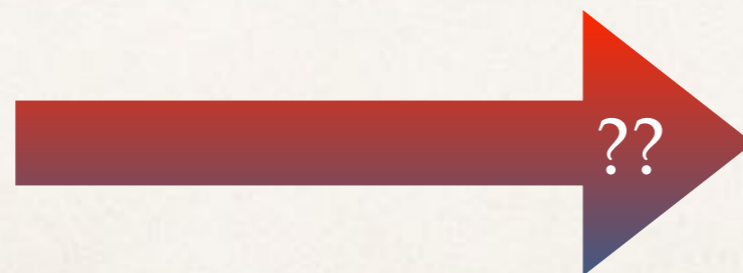
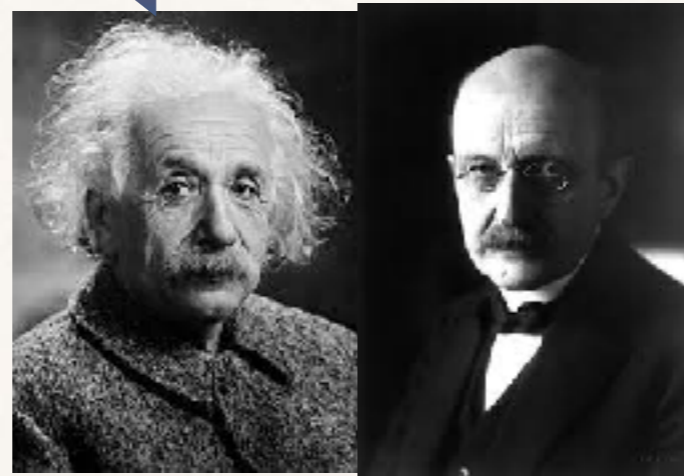
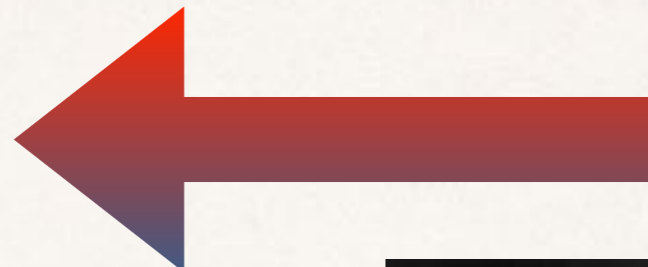
# Talk Outline

- Introduction
- Battery-Charger Set-up & Figures of Merit
- Results

# Talk Outline

- **Introduction**
- Battery-Charger Set-up & Figures of Merit
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# Quantum



# Thermodynamics

<u>École Polytechnique</u>	<u>Glasgow school</u>
	
<u>Sadi Carnot</u> (1796-1832)	<u>William Thomson</u> (1824-1907)
<u>Vienna school</u>	<u>Gibbsian school</u>
	
<u>Ludwig Boltzmann</u> (1844-1906)	<u>Willard Gibbs</u> (1839-1903)

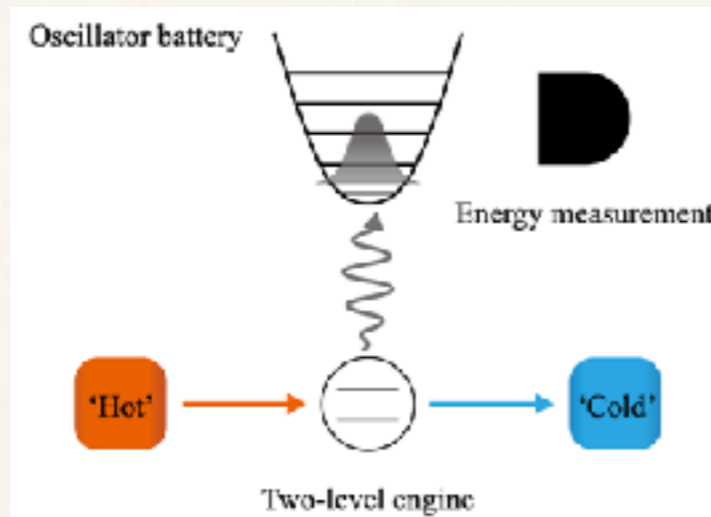
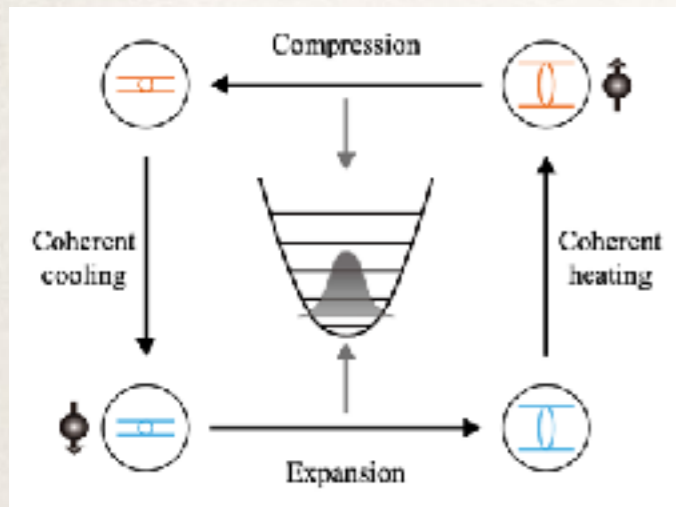
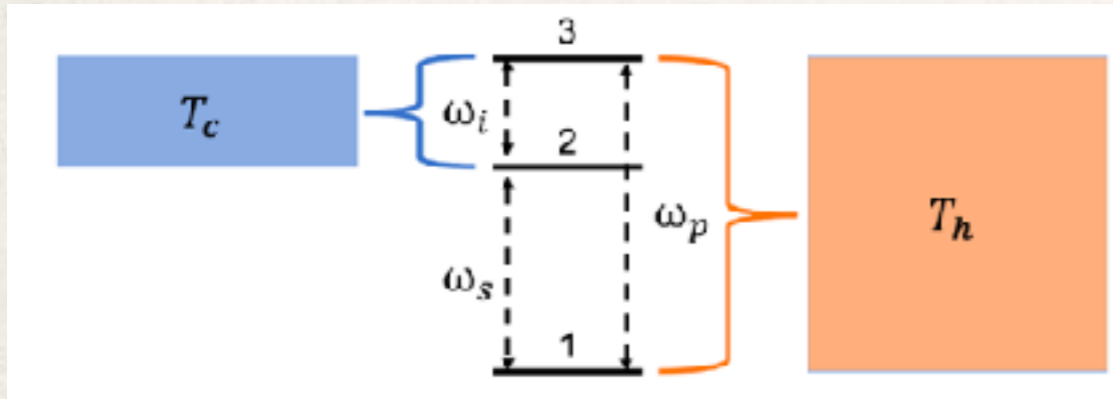
## Non-equilibrium Quantum Thermodynamics

S. Vinjanampathy and J. Anders, *Contemporary Physics*, 57, 545 (2016)

R. Kosloff, arXiv:1801.08314

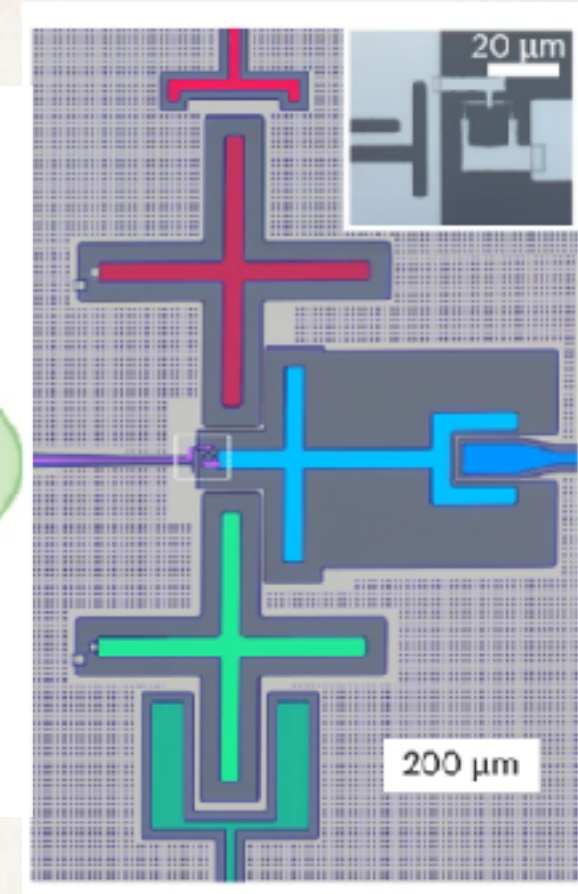
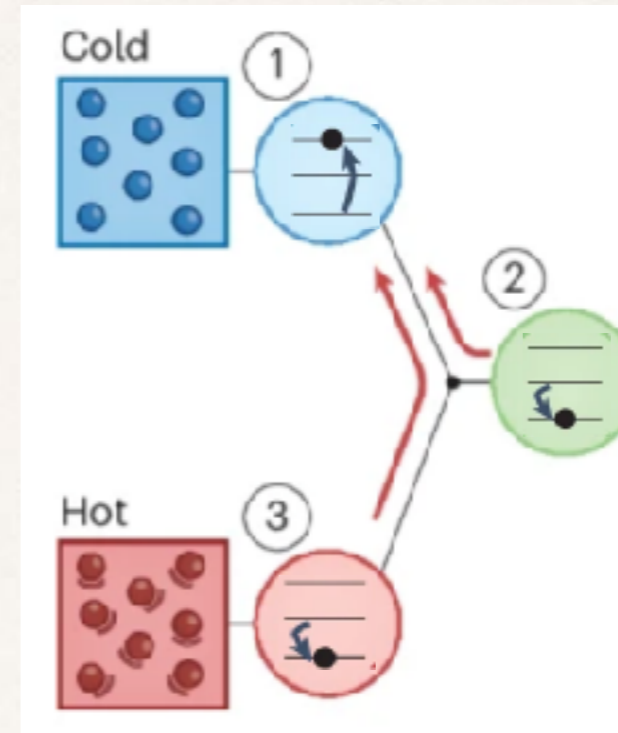
Slide Credit: Quantum Steampunk, Nicole Yunger-Halpern

# Quantum Thermal Machines



W. Hou et.al, arXiv:2404.15075

## Heat Engines



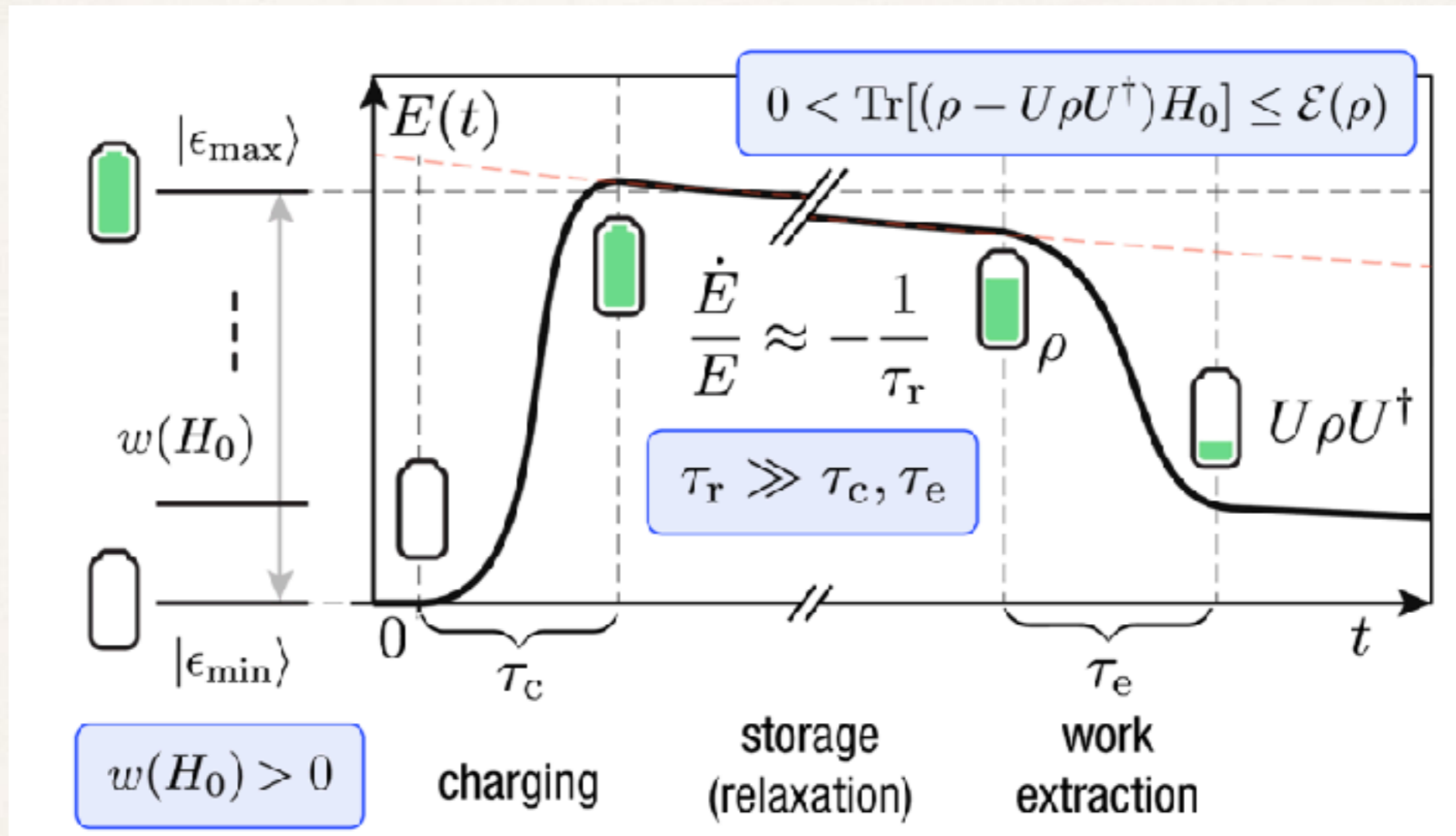
Aamir MA *et.al.*, Nat. Phys. (2025)  
Blok & Landi, Nat. Phys. (2025)

## Refrigerators

Quantum Effects for Enhanced Performance - Coherence, Entanglement, Non-Standard Baths, Correlated Many-Body States, Quantum Statistics...

S. Bhattacharjee and A. Dutta, EPJ B 94, 231 (2021)  
Myers *et.al.*, AVS Quantum Sci. 4, 027101(2022)

# Quantum Batteries



## Quantum Systems as Energy Storage Devices

Reviews:

S. Bhattacharjee and A. Dutta, EPJ B 94, 231 (2021)

J. Quach et.al., Joule 7, 2195 (2023)

F. Campaioli et.al., Rev. Mod. Phys. 96, 031001 (2024)



# Quantum Batteries

High Capacity, Fast Charging / Discharging, Stable Charging

## Entanglement

Alicki & Fannes, PRE **87**, 042123 (2013)

Hovhannisyanyan *et al.*, PRL **111**, 240401 (2013)

## Collective Effects Enhancement

F. Binder *et al.*, NJP **17**, 075015 (2015).

Campaioli *et al.*, PRL **118**, 150601 (2017)

S. Julià-Farré *et al.*, PRR **2**, 023113 (2020)

## Coherence

L.P. García-Pintos *et al.*, PRL **125**, 040601 (2020)

## Collective Effects Implementations

D. Ferraro *et al.*, PRL **120**, 117702 (2018)

D. Rossini *et al.*, PRL **125**, 236402 (2020)

D. Rosa *et al.*, JHEP **2020**, 67 (2020)

## Open Quantum Batteries

F. Barra, PRL **122**, 210601 (2019).

D. Farina *et al.*, PRB **99**, 035421 (2019).

J. Q. Quach, PR Applied **14**, 024092 (2020)

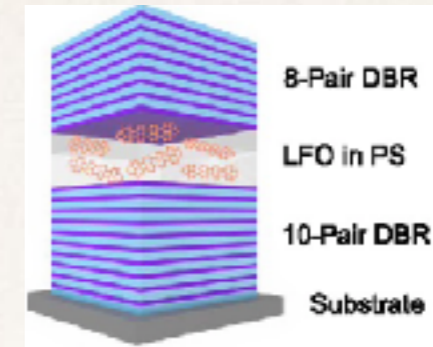
V. Shaghghi *et al.*, *Entropy*, **25**, 430 (2023)

..... many other works [pls see Rev. Mod. Phys. **96**, 031001 (2024)]

# Experimental Platforms

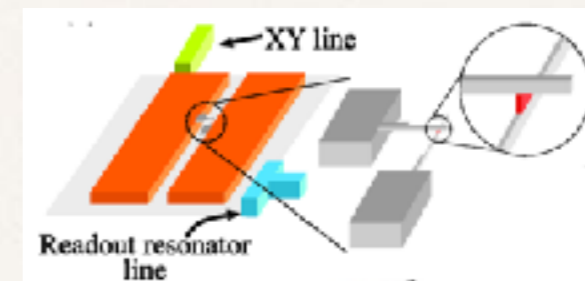
- Organic semiconductor microcavities

Quach *et al.*, *Sci. Adv.* 8 3160 (2022)



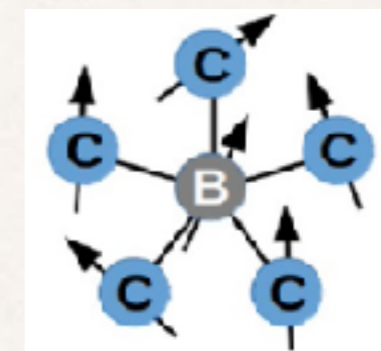
- Superconducting Circuits

Hu *et al.*, *Q. Sci. Tech.* 7, 045018 (2022)



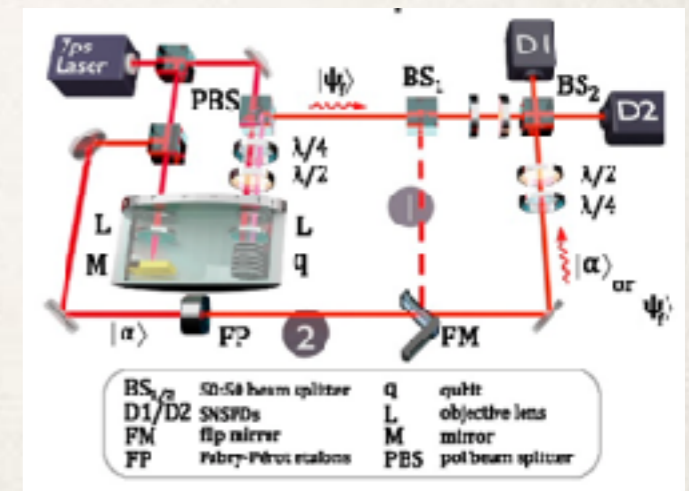
- NMR Spins

Joshi & Mahesh, *PRA* 106, 042601 (2022)



- Quantum Dots

Wenniger *et al.*, *PRL* 131, 260401 (2023)

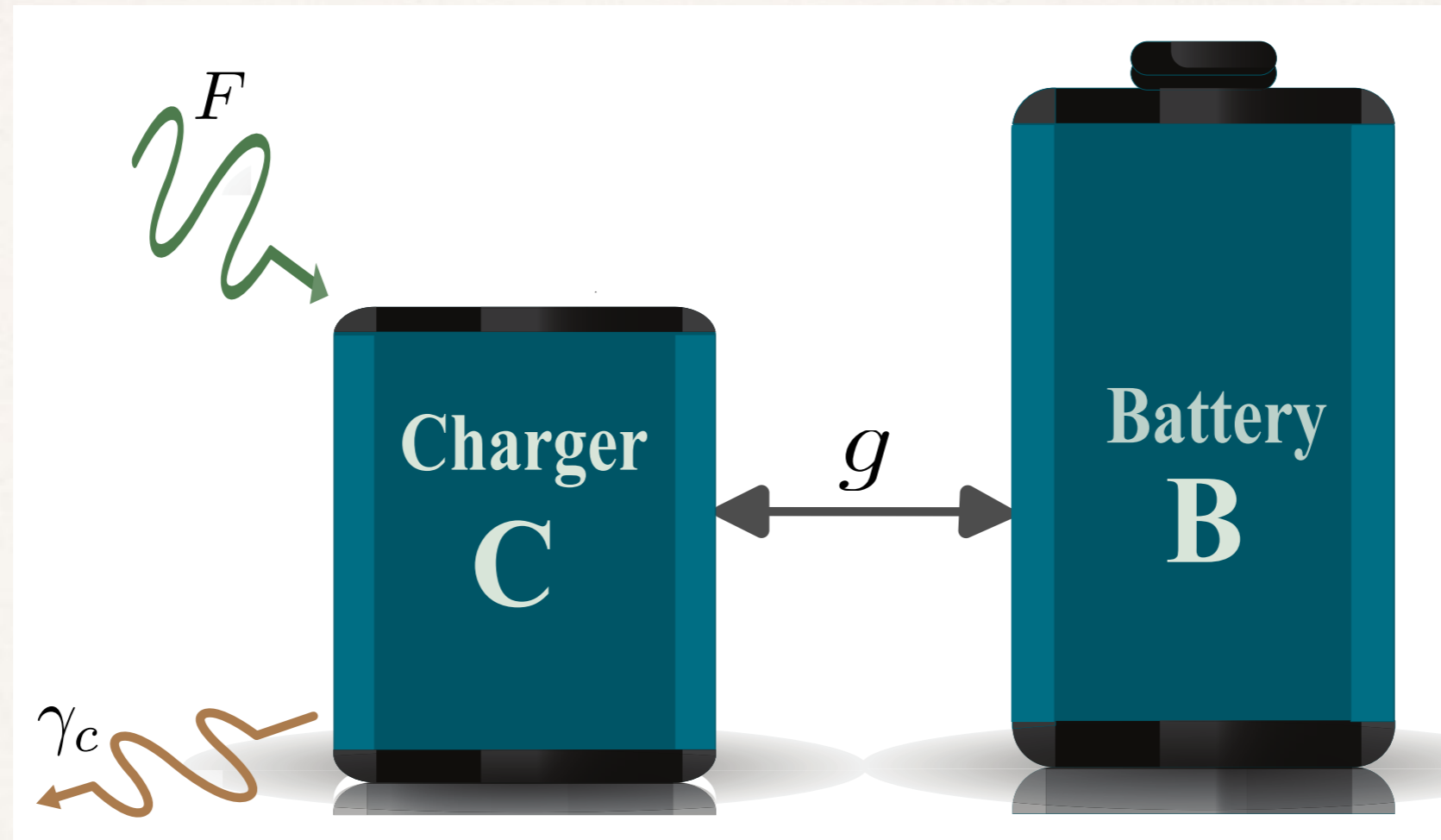


- Quantum Computers

Gemme *et al.*, *Batteries* 8, 43 (2022)



# This Talk



Open Model of Quantum Battery - Dissipation (Dephasing) as a resource

# Talk Outline

- Introduction
- **Battery-Charger Set-up & Figures of Merit**
- Results

# Dephasing

$$\begin{pmatrix} \rho_{ee} & \rho_{eg} \\ \rho_{ge} & \rho_{gg} \end{pmatrix} \rightarrow \begin{pmatrix} \rho_{ee} & \rho_{eg}e^{-\gamma t} \\ \rho_{ge}e^{-\gamma t} & \rho_{gg} \end{pmatrix}$$

Decay of coherences in energy basis

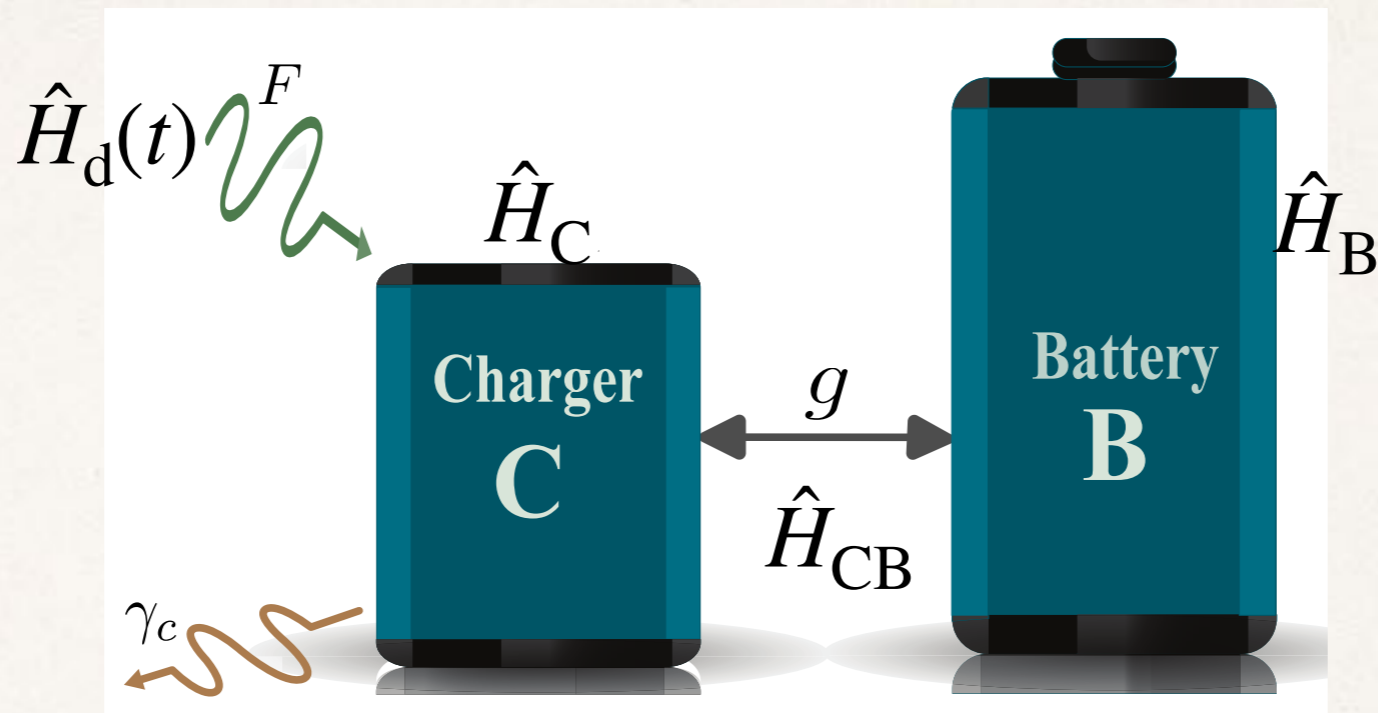
$$\frac{d\hat{\rho}(t)}{dt} = \gamma \left( \hat{L}\hat{\rho}(t)\hat{L} - \frac{\{\hat{L}^2, \hat{\rho}(t)\}}{2} \right) \quad \hat{L} \propto \hat{H}$$

Physical Realization:

- Continuous energy Measurement
- Noisy external classical field for TLS

$$\hat{H}_{\text{stoch}}(t) = B(t)\hat{\sigma}_z, \quad B(t) = \sqrt{\frac{\gamma}{2}}\xi(t)$$

# Battery + Dephased Charger

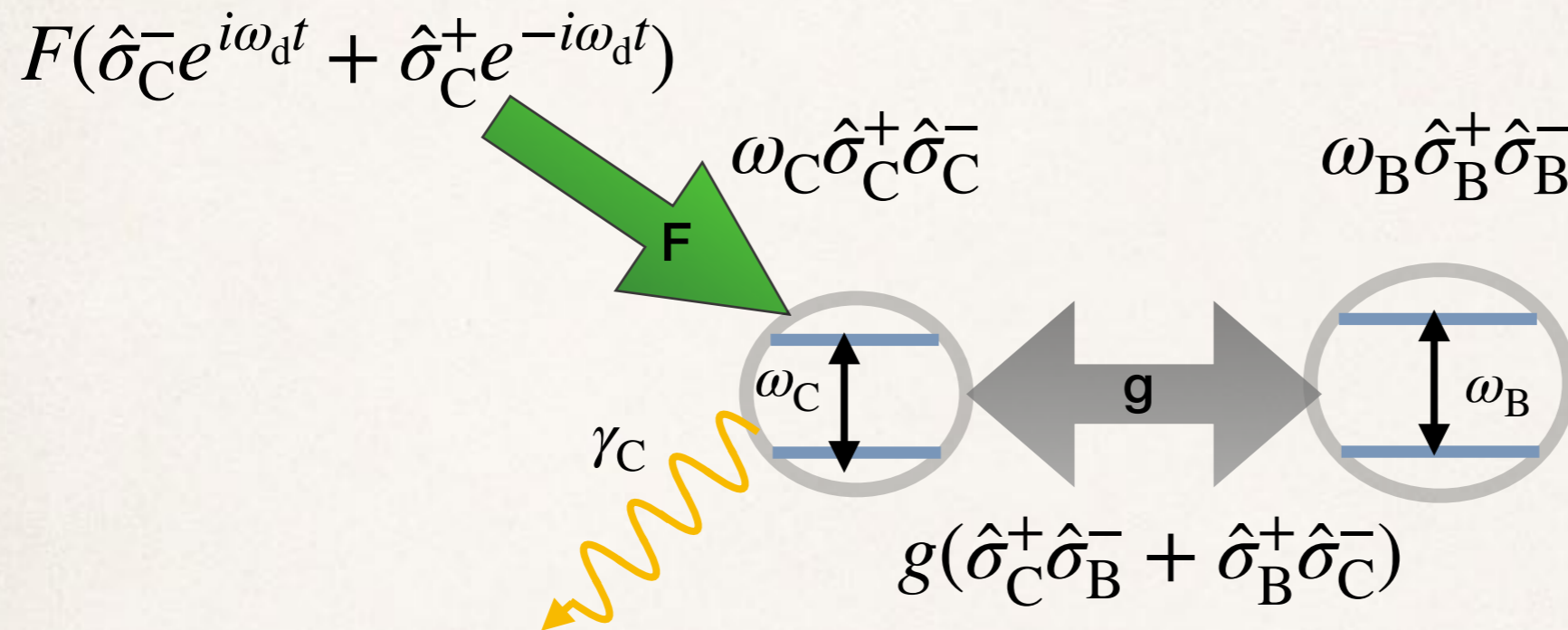


$$\gamma_C \left( \hat{L}_C \hat{\rho}(t) \hat{L}_C - \frac{\{\hat{L}_C^2, \hat{\rho}(t)\}}{2} \right), \quad [\hat{L}_C, \hat{H}_C] = 0$$

$$\hat{H}(t) = \hat{H}_C + \hat{H}_d(t) + \hat{H}_B + \hat{H}_{CB}$$

$$\frac{d\hat{\rho}(t)}{dt} = -i [\hat{H}, \hat{\rho}(t)] + \gamma_C \left( \hat{L}_C \hat{\rho}(t) \hat{L}_C - \frac{\{\hat{L}_C^2, \hat{\rho}(t)\}}{2} \right)$$

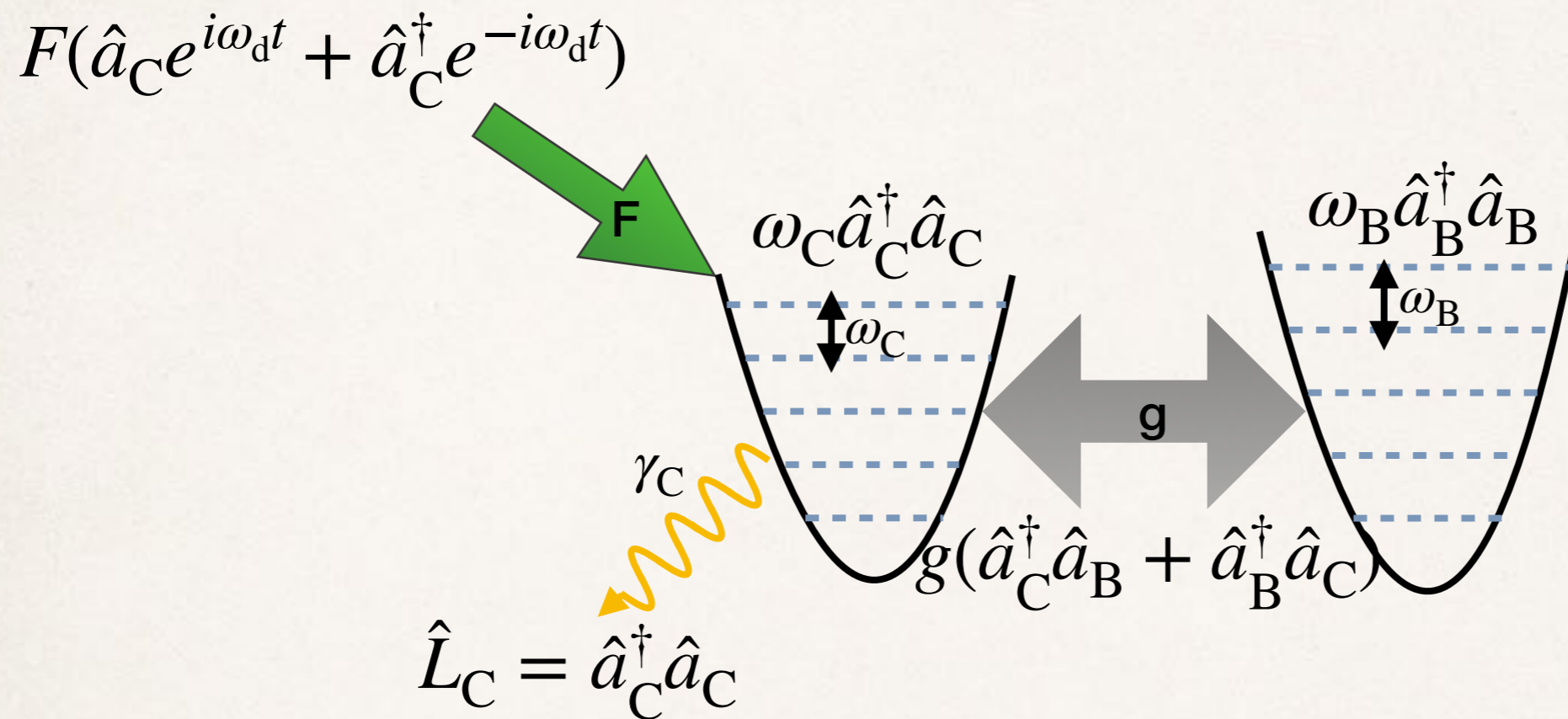
# Two Level Systems (TLSs)



$$\hat{L}_C = \hat{\sigma}_C^+ \hat{\sigma}_C^- = (\mathbb{1} + \hat{\sigma}_C^z)/2$$

$$\frac{d\hat{\rho}(t)}{dt} = -i \left[ \hat{H}, \hat{\rho}(t) \right] + \gamma_C \left( \hat{L}_C \hat{\rho}(t) \hat{L}_C - \frac{\{\hat{L}_C^2, \hat{\rho}(t)\}}{2} \right)$$

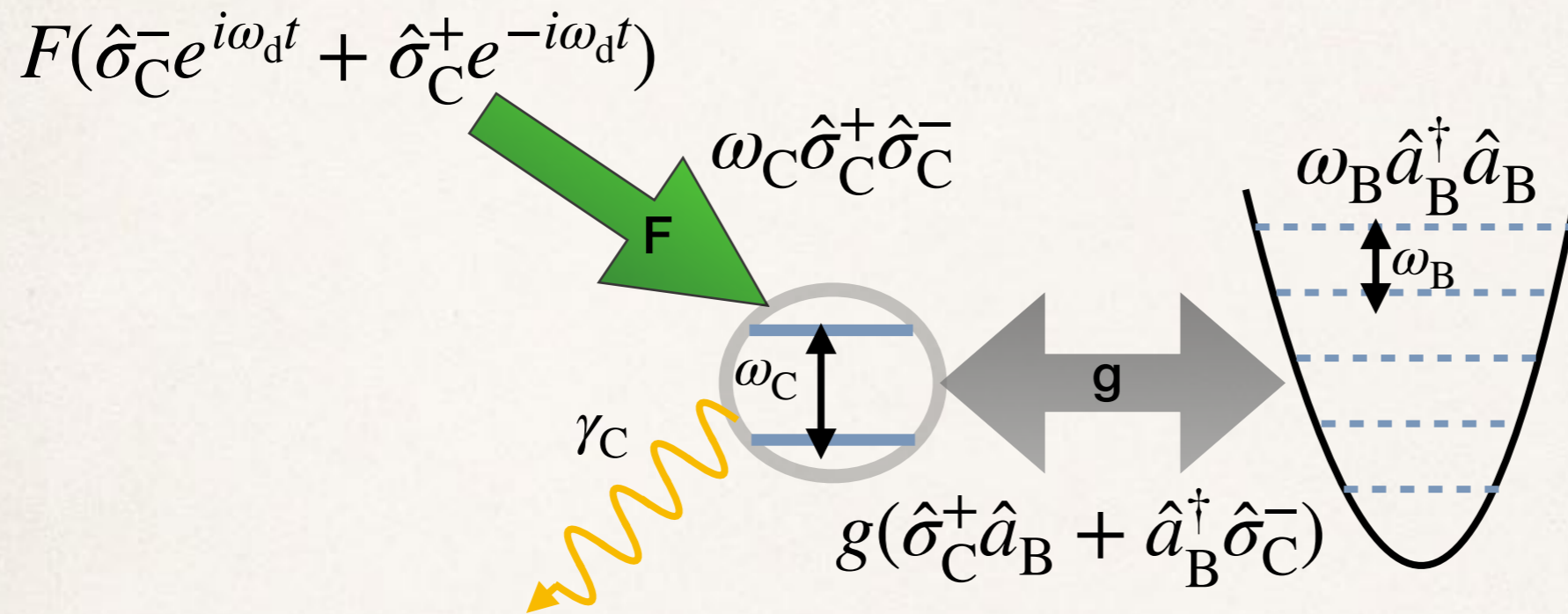
# Harmonic Oscillators (HOs)



$$\frac{d\hat{\rho}(t)}{dt} = -i \left[ \hat{H}, \hat{\rho}(t) \right] + \gamma_C \left( \hat{L}_C \hat{\rho}(t) \hat{L}_C - \frac{\{\hat{L}_C^2, \hat{\rho}(t)\}}{2} \right)$$



# Hybrid TLS-HO



$$\hat{L}_C = \hat{\sigma}_C^+ \hat{\sigma}_C^- = (\mathbb{1} + \hat{\sigma}_C^z)/2$$

$$\frac{d\hat{\rho}(t)}{dt} = -i \left[ \hat{H}, \hat{\rho}(t) \right] + \gamma_C \left( \hat{L}_C \hat{\rho}(t) \hat{L}_C - \frac{\{\hat{L}_C^2, \hat{\rho}(t)\}}{2} \right)$$

# Figures of Merit

## Battery Energy

$$E_B = \text{Tr}_B \left[ \hat{\rho}_B \hat{H}_B \right]$$

Maximum extractable energy

## Ergotropy

“quality of energy”

Allahverdyan *et al.* EPL 67, 565 (2004)

$$\mathcal{E}_B = E_B - \min_{\hat{U}_B} \text{Tr}_B \left[ \hat{U}_B \hat{\rho}_B \hat{U}_B^\dagger \hat{H}_B \right] = \text{Tr}[\hat{\rho}_B \hat{H}_B] - \text{Tr}[\hat{\rho}_P \hat{H}_B]$$

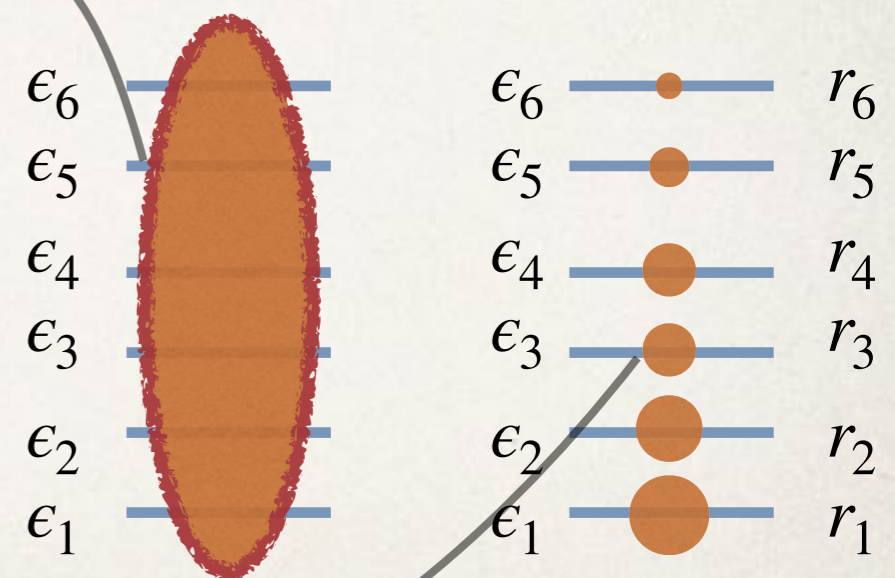
$$\hat{H}_B = \sum_{j=1}^N \epsilon_j |\epsilon_j\rangle \langle \epsilon_j|$$

$$\hat{\rho}_B = \sum_{j=1}^N r_j |r_j\rangle \langle r_j|$$

$$\epsilon_1 \leq \epsilon_2 \leq \dots \leq \epsilon_N$$

$$r_1 \geq r_2 \geq \dots \geq r_N$$

$$\hat{\rho}_P = \sum_{j=1}^N r_j |\epsilon_j\rangle \langle \epsilon_j|$$



# Figures of Merit

## Battery Energy

$$E_B = \text{Tr}_B \left[ \hat{\rho}_B \hat{H}_B \right]$$

## Ergotropy

$$\mathcal{E}_B = E_B - \min_{\hat{U}_B} \text{Tr}_B \left[ \hat{U}_B \hat{\rho}_B \hat{U}_B^\dagger \hat{H}_B \right]$$

## Charging Time

$$\left| \frac{E_B(\tau) - E_B(\infty)}{E_B(0) - E_B(\infty)} \right| = e^{-n}$$

# Figures of Merit - TLS

## Battery Energy

$$E_B = \frac{\omega_B}{2} (\langle \hat{\sigma}_B^z \rangle + 1)$$

## Ergotropy

$$\mathcal{E}_B = \frac{\omega_B}{2} \left( \sqrt{\langle \hat{\sigma}_B^z \rangle^2 + 4\langle \hat{\sigma}_B^+ \rangle \langle \hat{\sigma}_B^- \rangle} + \langle \hat{\sigma}_B^z \rangle \right)$$

## Charging Time

$$\left| \frac{E_B(\tau) - E_B(\infty)}{E_B(0) - E_B(\infty)} \right| = e^{-n}$$

# Figures of Merit - HO

## Battery Energy

$$E_B = \omega_B \langle \hat{a}_B^\dagger \hat{a}_B \rangle$$

## Ergotropy

$$\mathcal{E}_B = E_B - \min_{\hat{U}_B} \text{Tr}_B \left[ \hat{U}_B \hat{\rho}_B \hat{U}_B^\dagger \hat{H}_B \right]$$

## Charging Time

$$\left| \frac{E_B(\tau) - E_B(\infty)}{E_B(0) - E_B(\infty)} \right| = e^{-n}$$

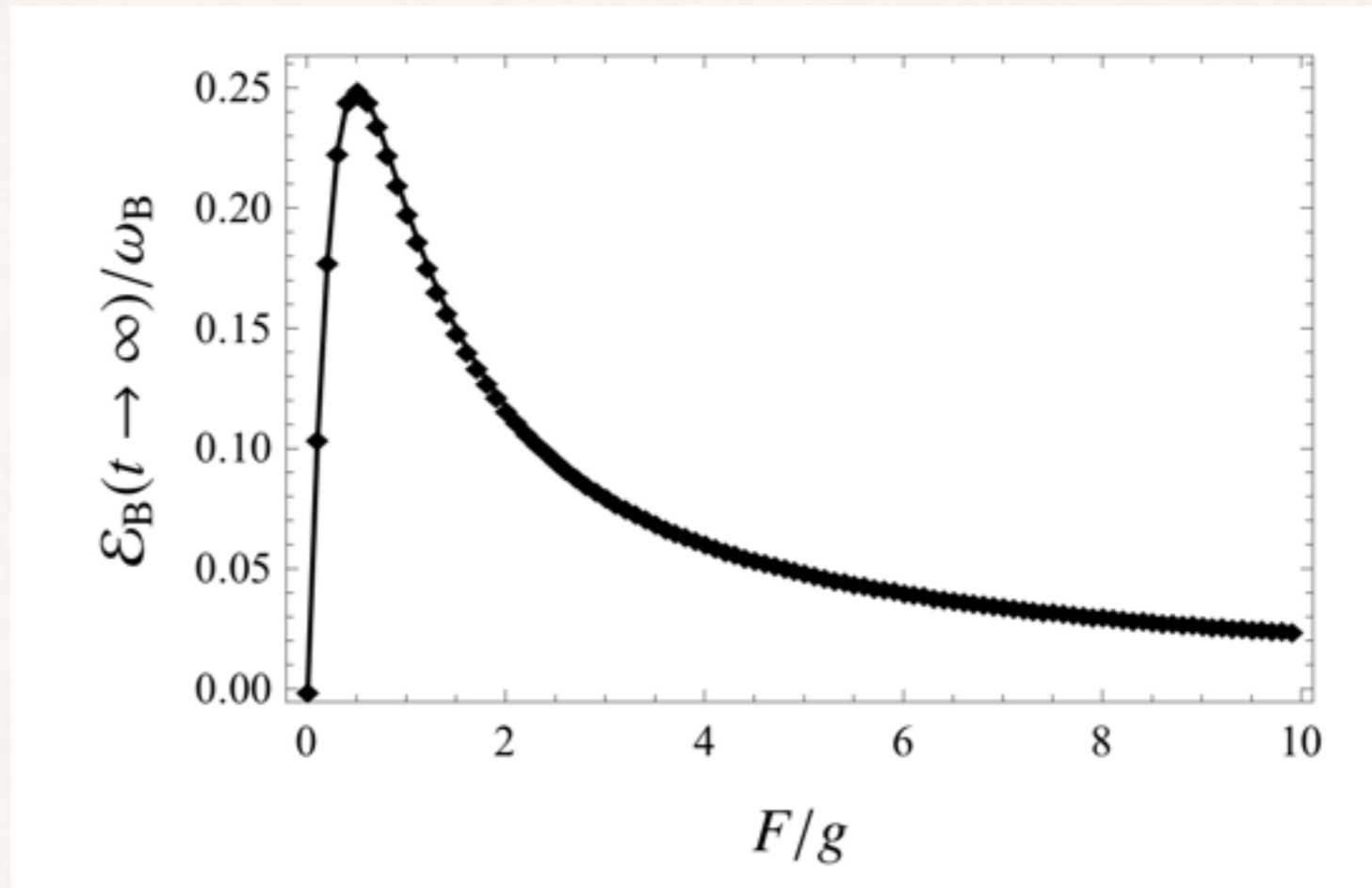
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# TLS at Resonance - Steady State

$$\omega_C = \omega_d = \omega_B$$



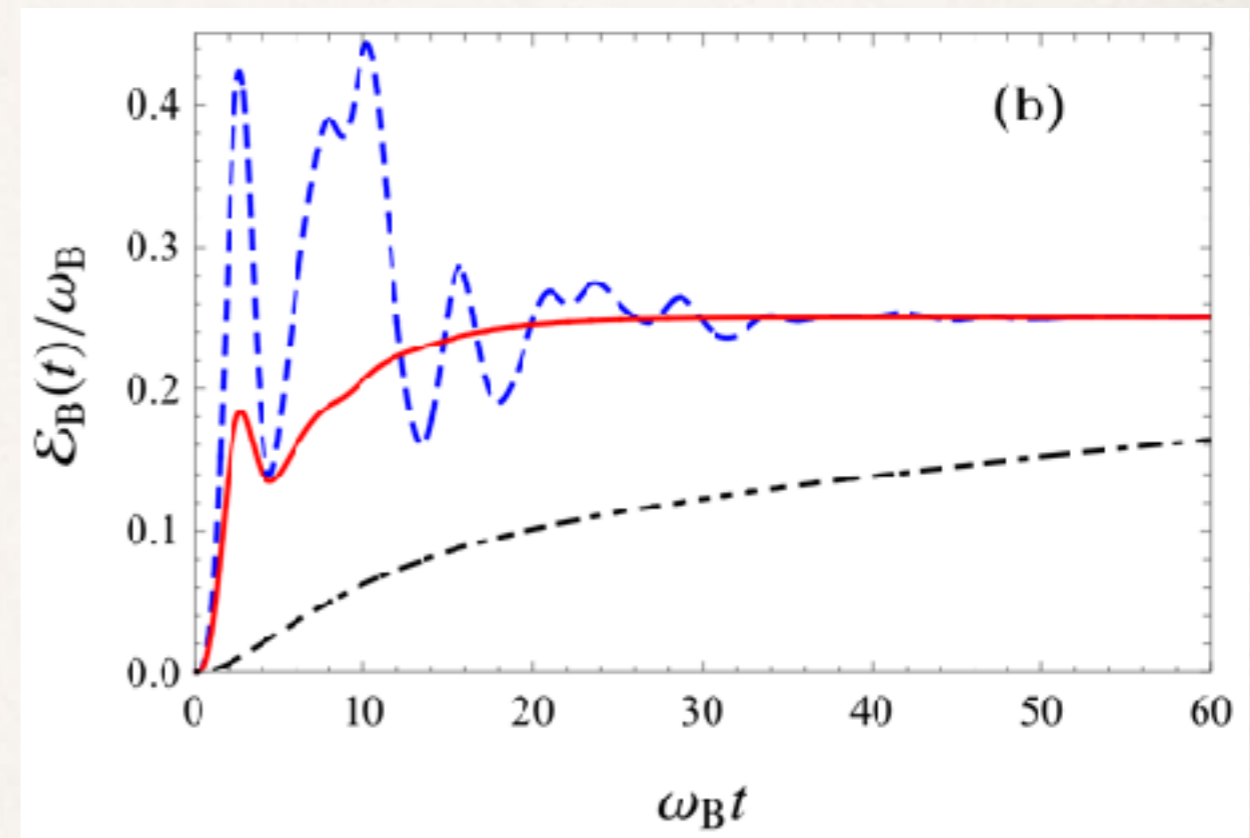
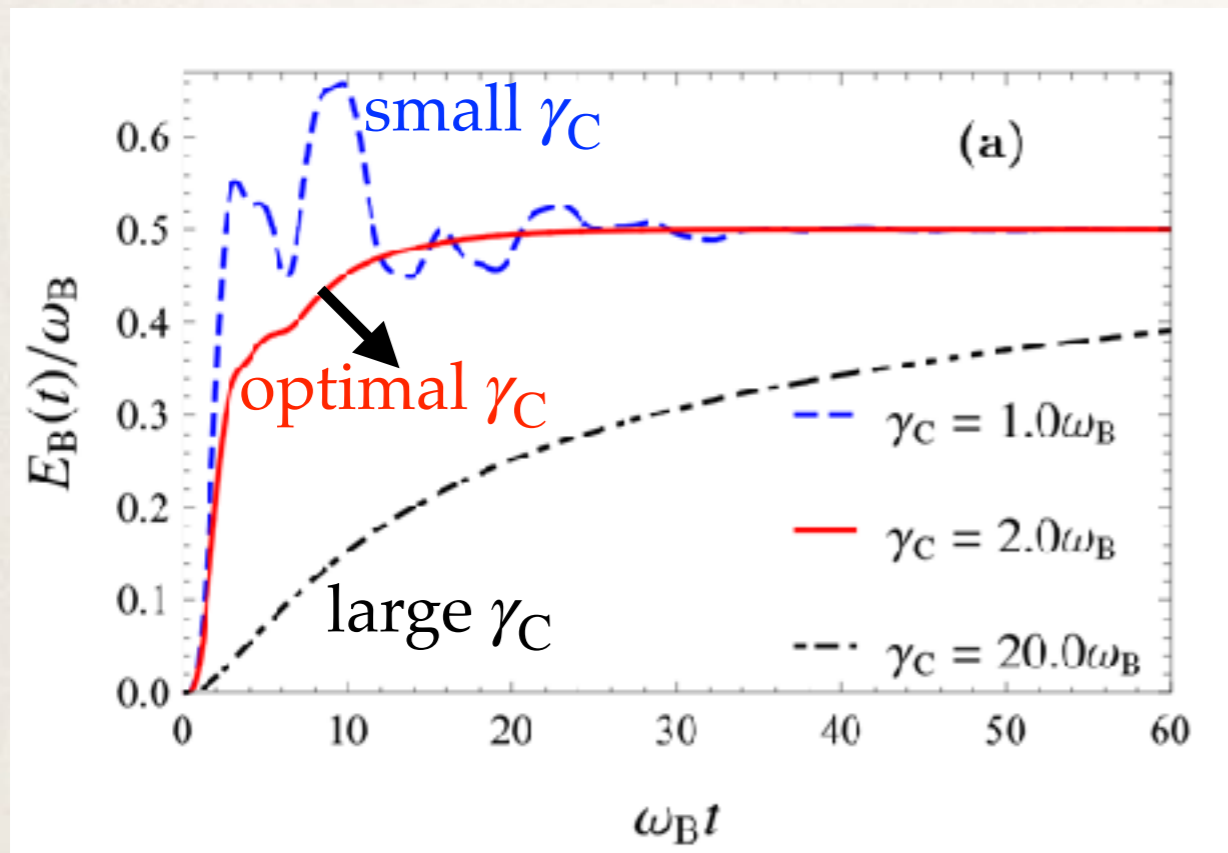
$$E_B(t \rightarrow \infty) = \frac{1}{2},$$

$$\mathcal{E}_B(t \rightarrow \infty) = \frac{\frac{F}{g}}{1 + 4\frac{F^2}{g^2}}.$$

# TLS at Resonance - Charging Dynamics

$F/g = 0.5$  (intermediate optimal driving)

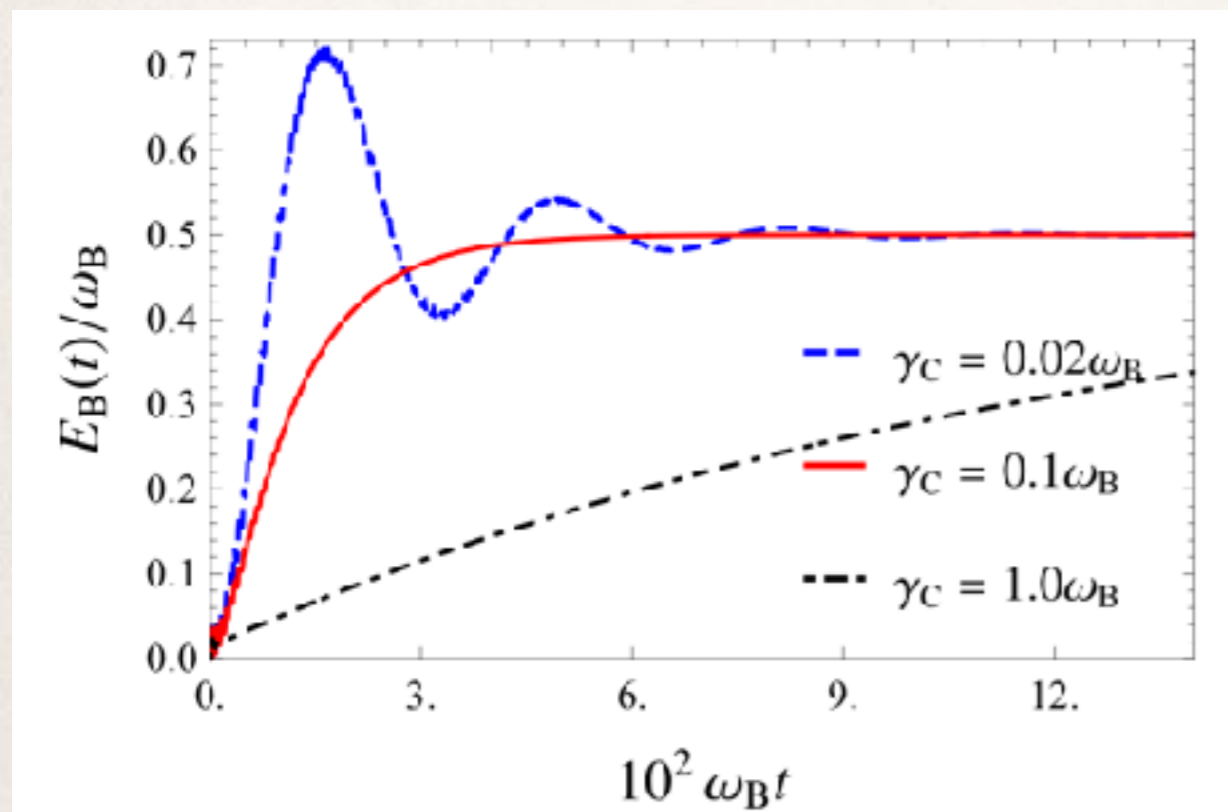
$g = \omega_B$  (throughout)



Moderate Dephasing Leads to Fast Charging!

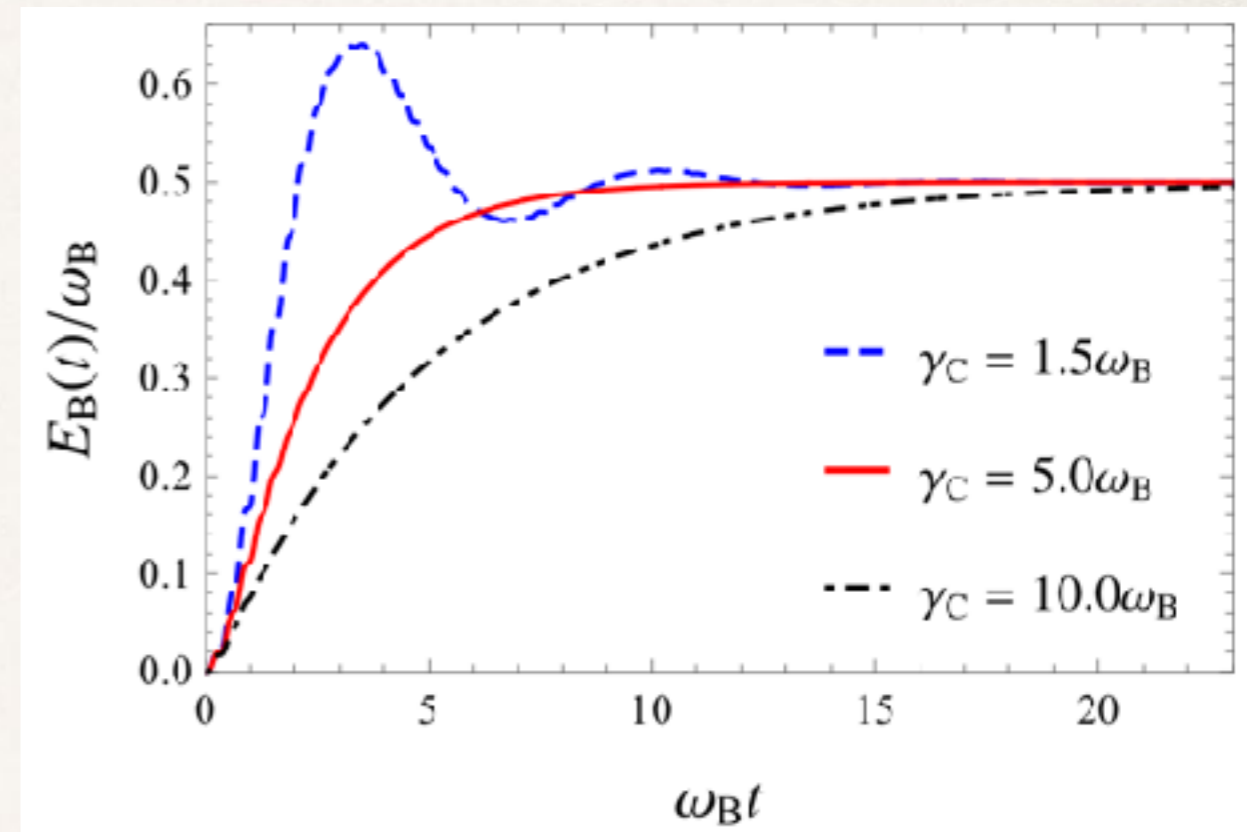


# TLS at Resonance - Charging Dynamics



$$F/g = 0.1$$

Weak Driving



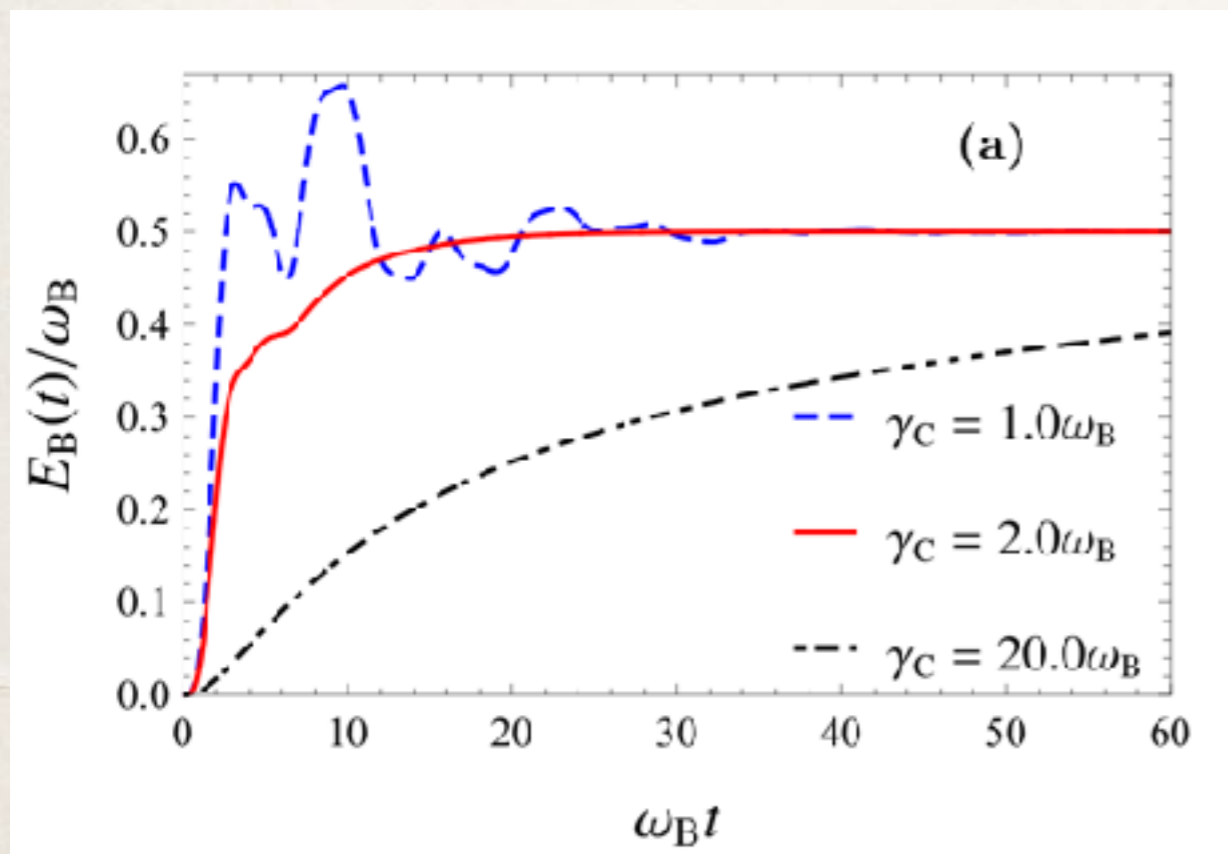
$$F/g = 10$$

Strong Driving

Moderate Dephasing Leads to Fast Charging!

# TLS at Resonance - Charging Dynamics

Moderate Dephasing Leads to Fast Charging



- Small Dephasing - Oscillation
- Large Dephasing - quantum Zeno

Transient Maxima - impractical, require fine control

# Universal Competition under Dephasing

Competition:

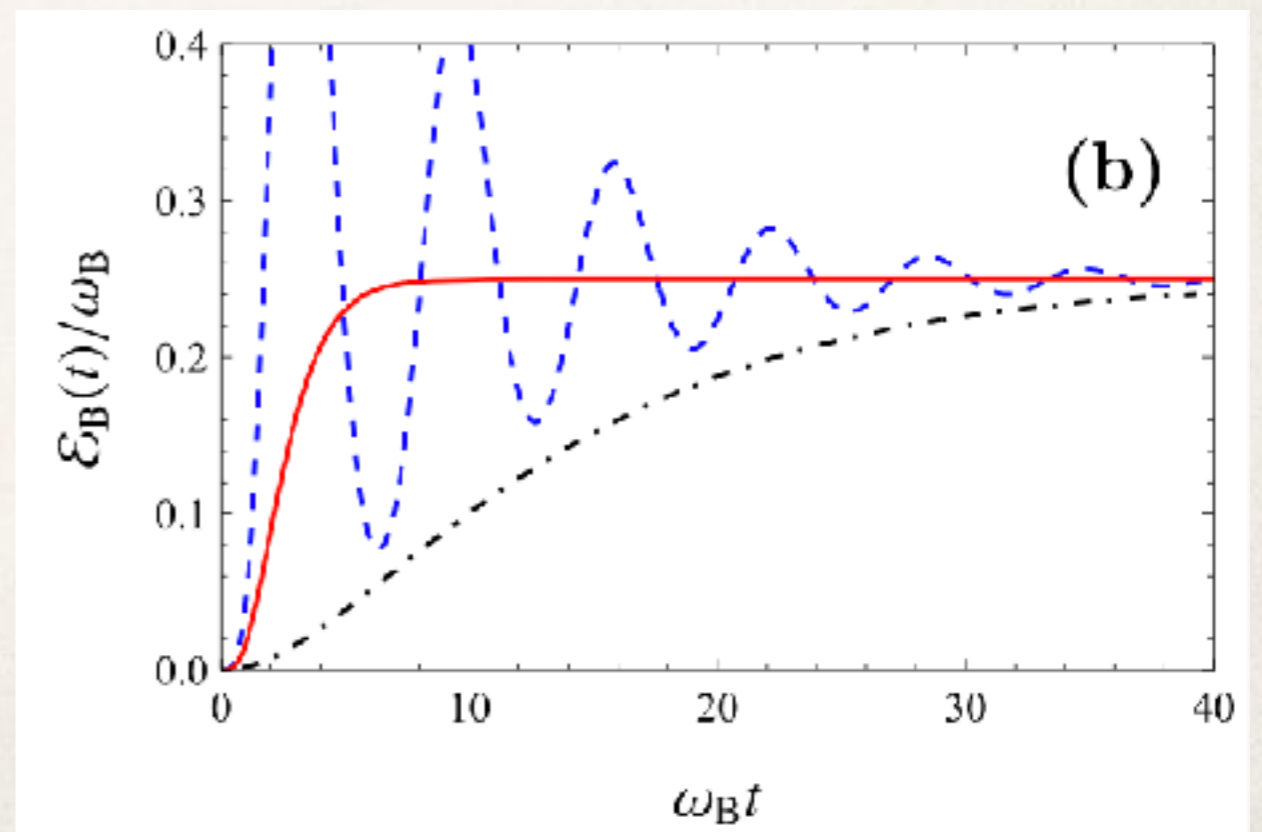
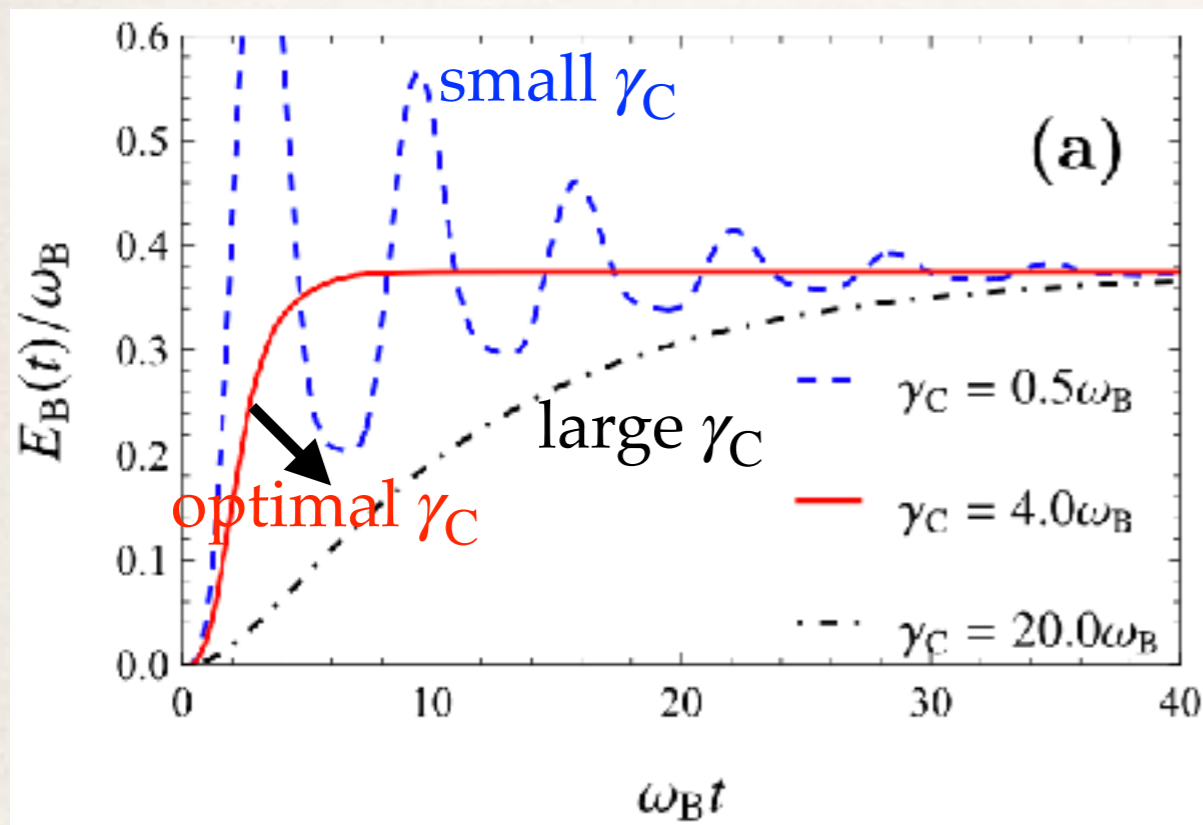
Coherent Oscillation

vs

Quantum Zeno Freezing

Universal for systems under dephasing

HOs at resonance



# Universal Competition under Dephasing

Competition:

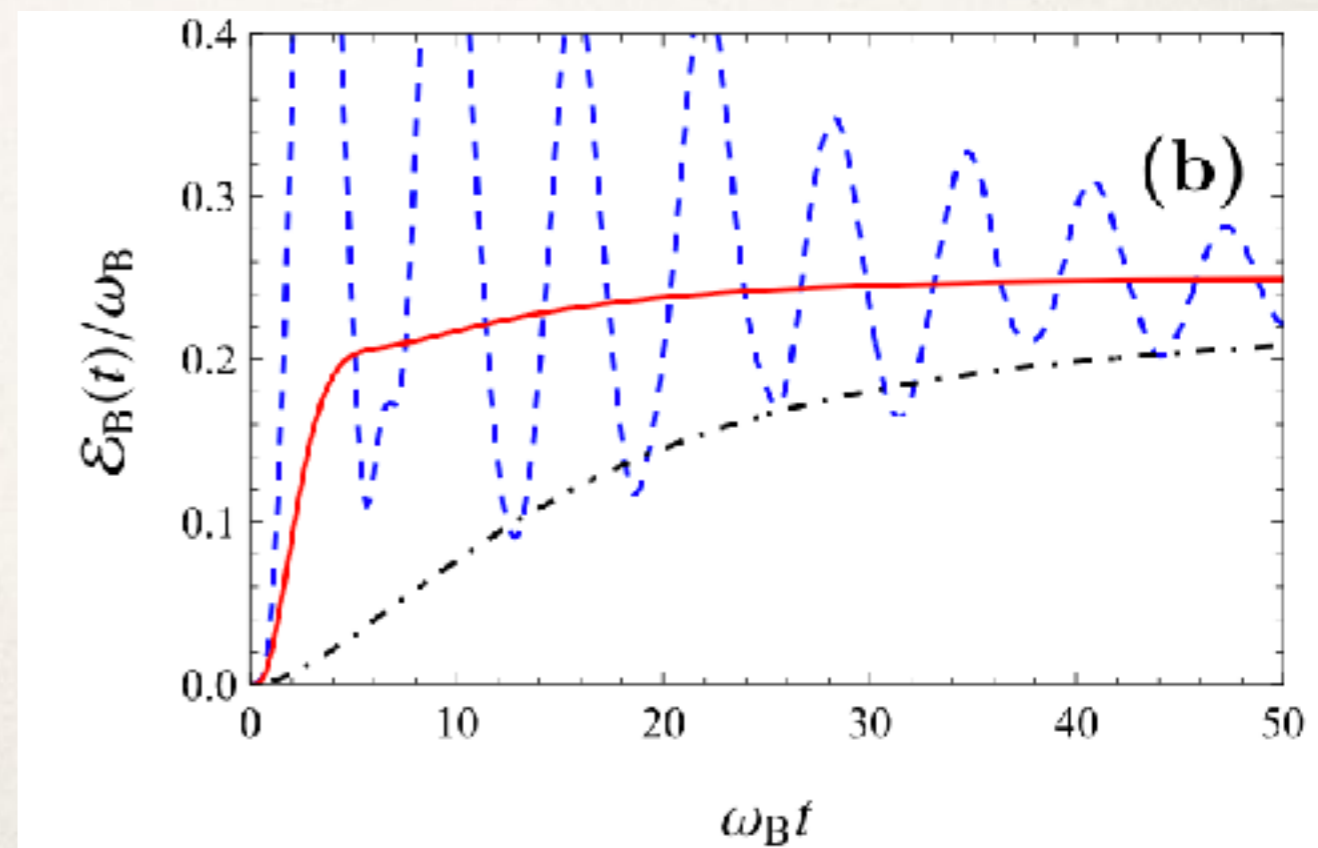
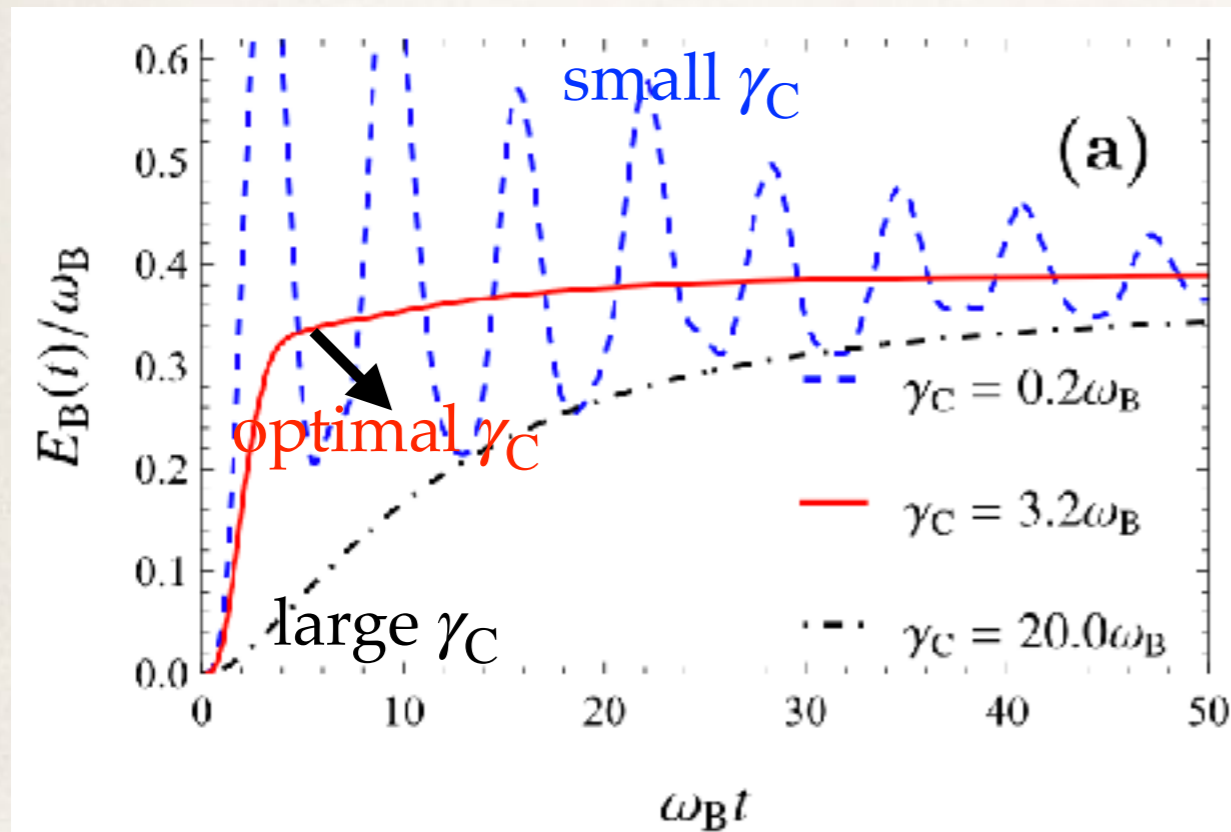
Coherent Oscillation

vs

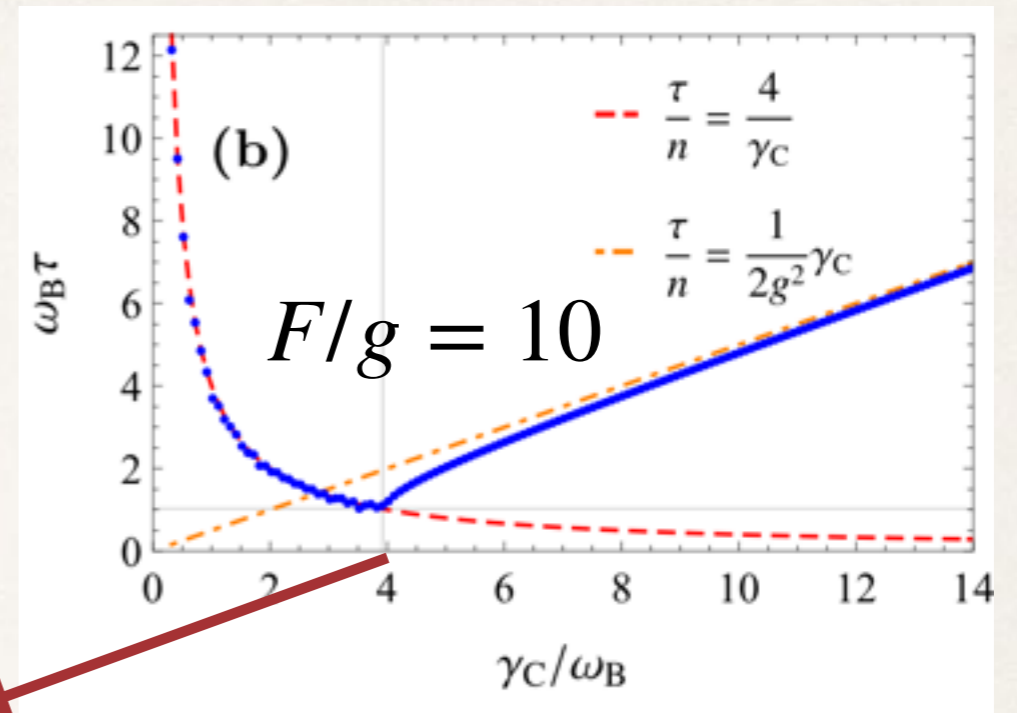
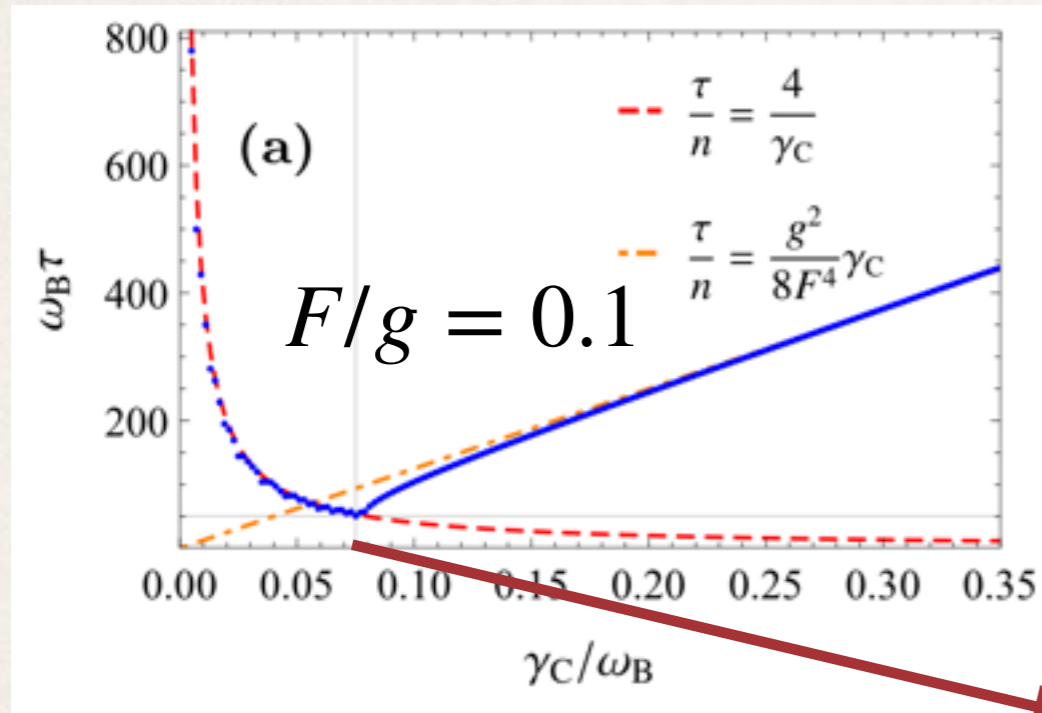
Quantum Zeno Freezing

Universal for systems under dephasing

TLS-HO at resonance

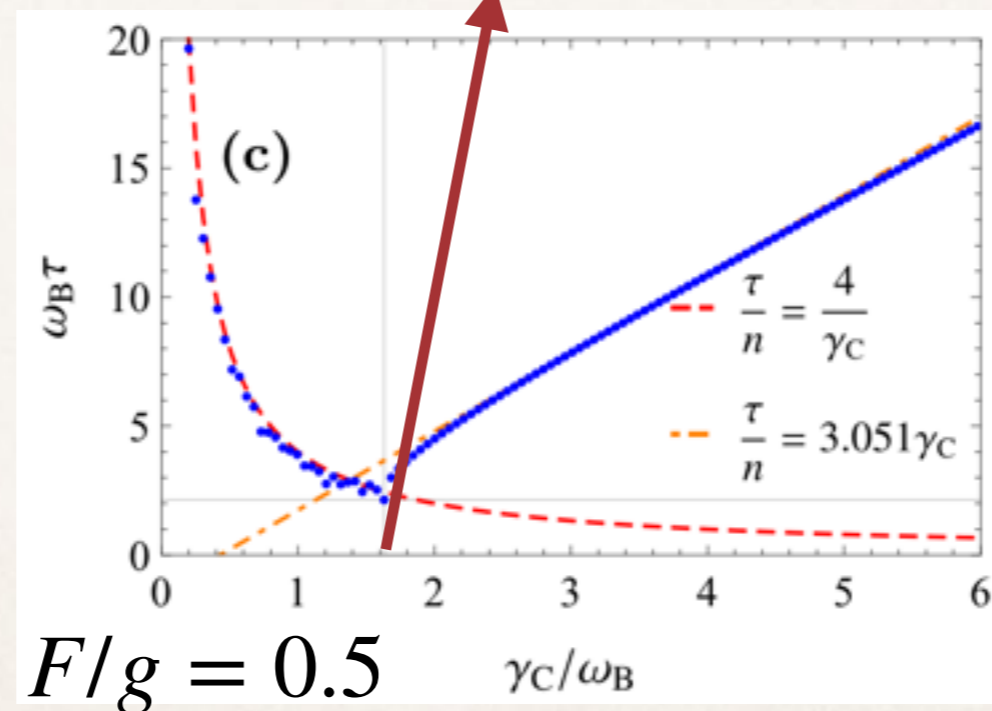


# TLS - Charging Time vs Dephasing



Optimal

$\gamma_C$



Weak Dephasing

$$\gamma_C \ll \{F, g\}$$

$$\tau \sim \frac{1}{\gamma_C}$$

Strong Dephasing

$$\gamma_C \gg \{F, g\}$$

$$\tau \sim \frac{\gamma_C}{[\omega^2]}$$

$$\text{Min } \tau \text{ when } \frac{1}{\gamma_C} \sim \frac{\gamma_C}{[\omega^2]}$$

# TLS - Optimal Dephasing Analytical Results

$$F/g \gg 1$$

$$F/g \ll 1$$

$$\gamma_C \ll g$$

$$\tau \sim \frac{4n}{\gamma_C}$$

$$\tau \sim \frac{4n}{\gamma_C}$$

$$\gamma_C \gg \{g, F\}$$

$$\tau \sim \frac{n}{2g^2} \gamma_C$$

$$\tau \sim \frac{ng^2}{F^4} \gamma_C$$

## Optimal Dephasing

$$\gamma_C^* \stackrel{F \ll g}{\approx} \frac{8F^2}{\sqrt{2}g},$$

$$\gamma_C^* \stackrel{F \gg g}{\approx} 2\sqrt{2}g,$$

$$\tau^* \sim \frac{4}{\gamma_C^*}$$

# Charging Strategy

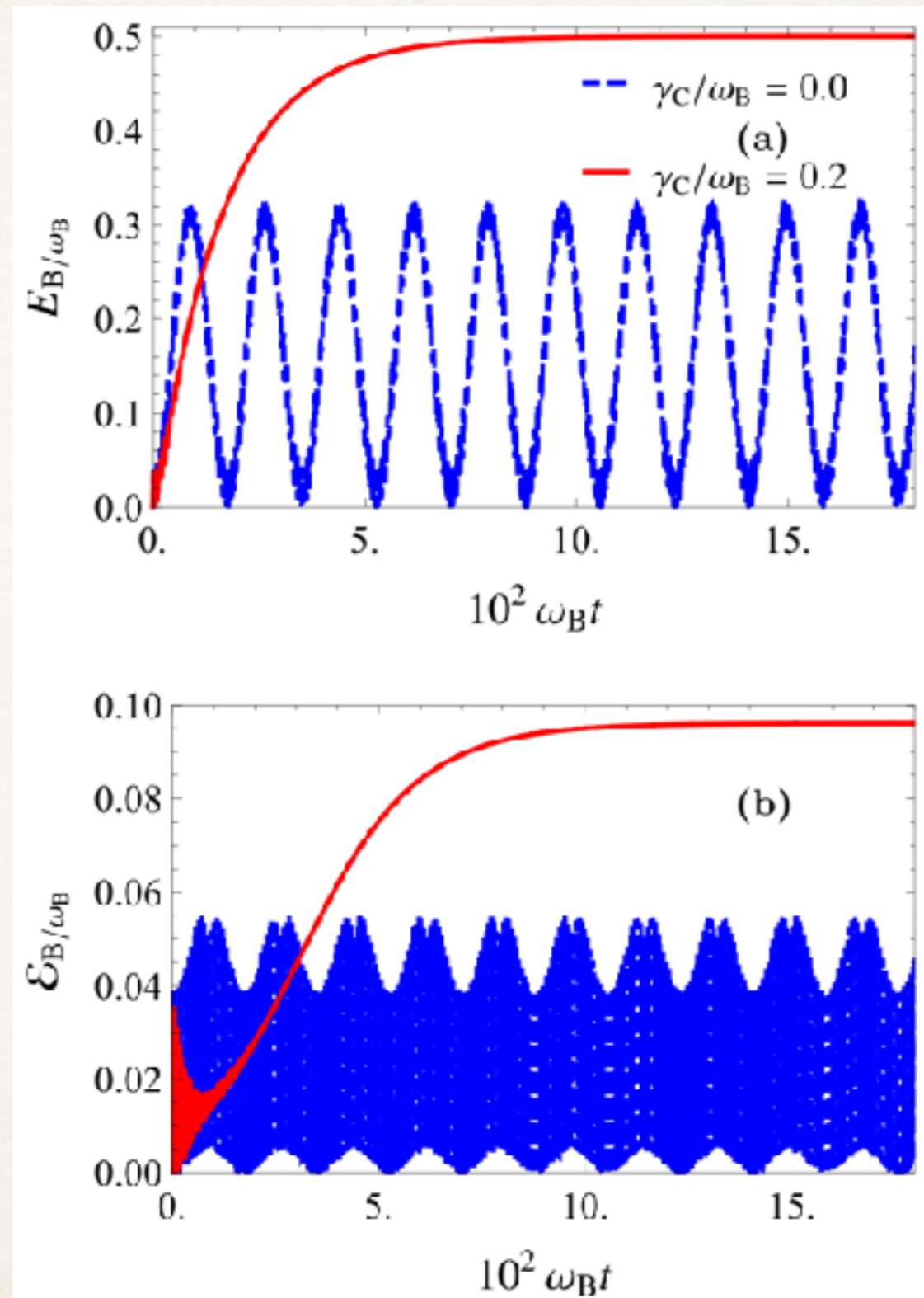
Choose  $F/g = 0.5$ , maximising steady state ergotropy

Since  $\gamma_C^* \sim a \frac{F}{g} F + bg$ , make  $\{F, g\}$  as large as possible

$$\gamma_C^* \stackrel{F \ll g}{\approx} \frac{8F^2}{\sqrt{2g}},$$

$$\gamma_C^* \stackrel{F \gg g}{\approx} 2\sqrt{2g},$$

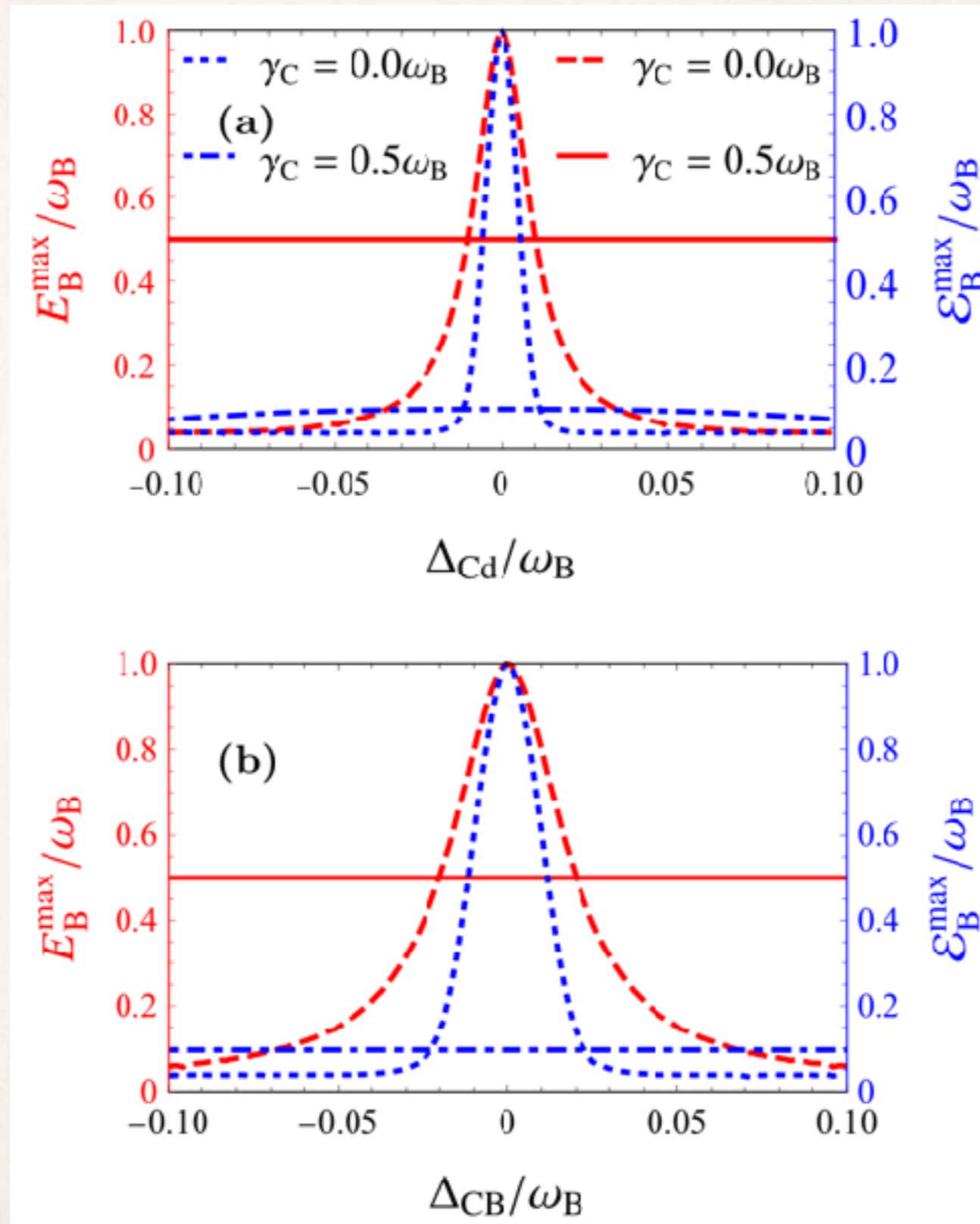
# TLS - Robustness To Detuning



$$\Delta_{Cd} = \Delta_{Bd} = 0.015\omega_B$$



# TLS - Robustness To Detuning



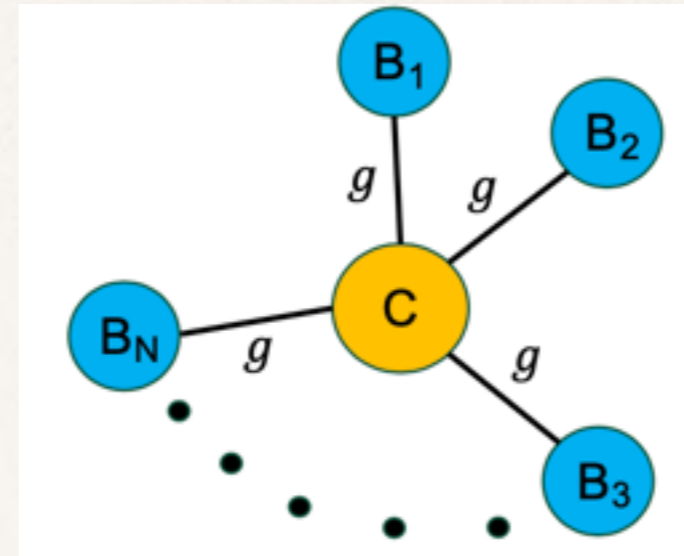
# Summary

- Moderate Dephasing of Charger leads to fast and stable battery charging
- General Behaviour - Coherent Oscillations vs Quantum Zeno Trade off
- Dephasing - Robustness w.r.t. Finite Detuning

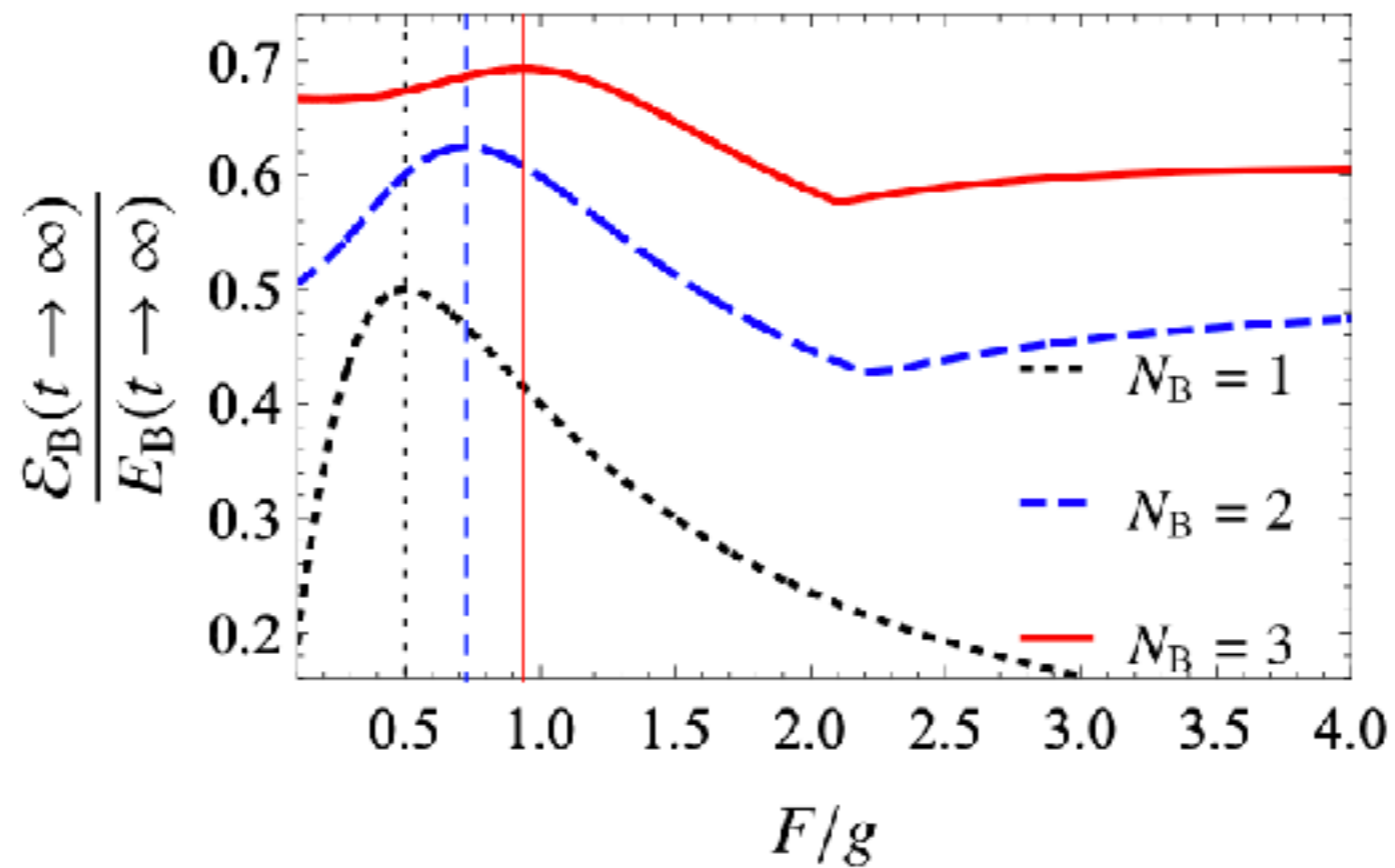
R. Shastri et.al., npj Quantum Inf **11**, 9 (2025)



# Outlook- Enhancing Ergotropy

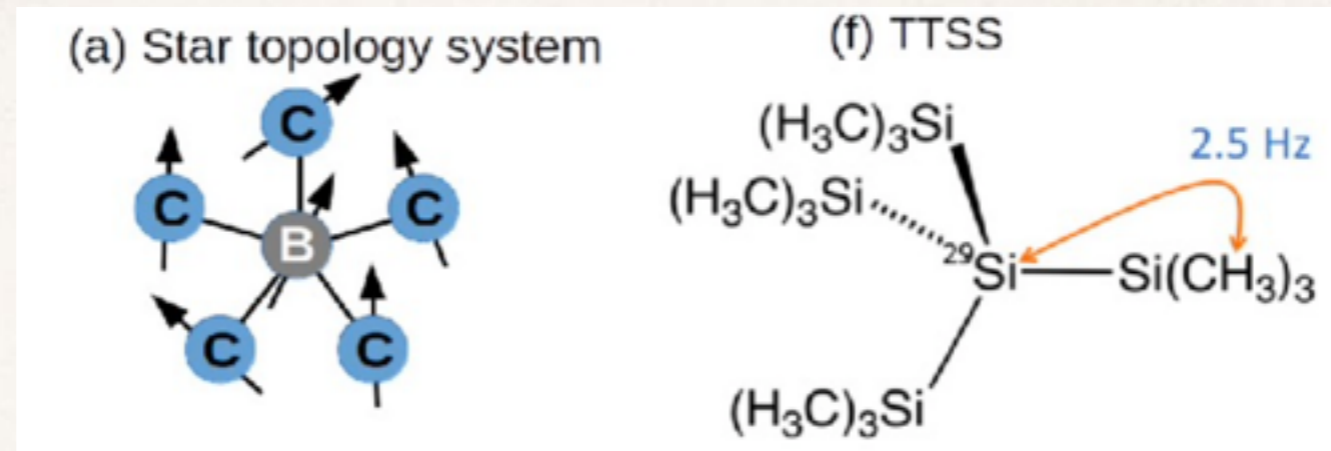


Star Config



# Outlook- Experimental Implementation

NMR System



1 Battery - coupled to multiple charger

Battery - Central Nuclear Spin ( $I = 1/2$ )

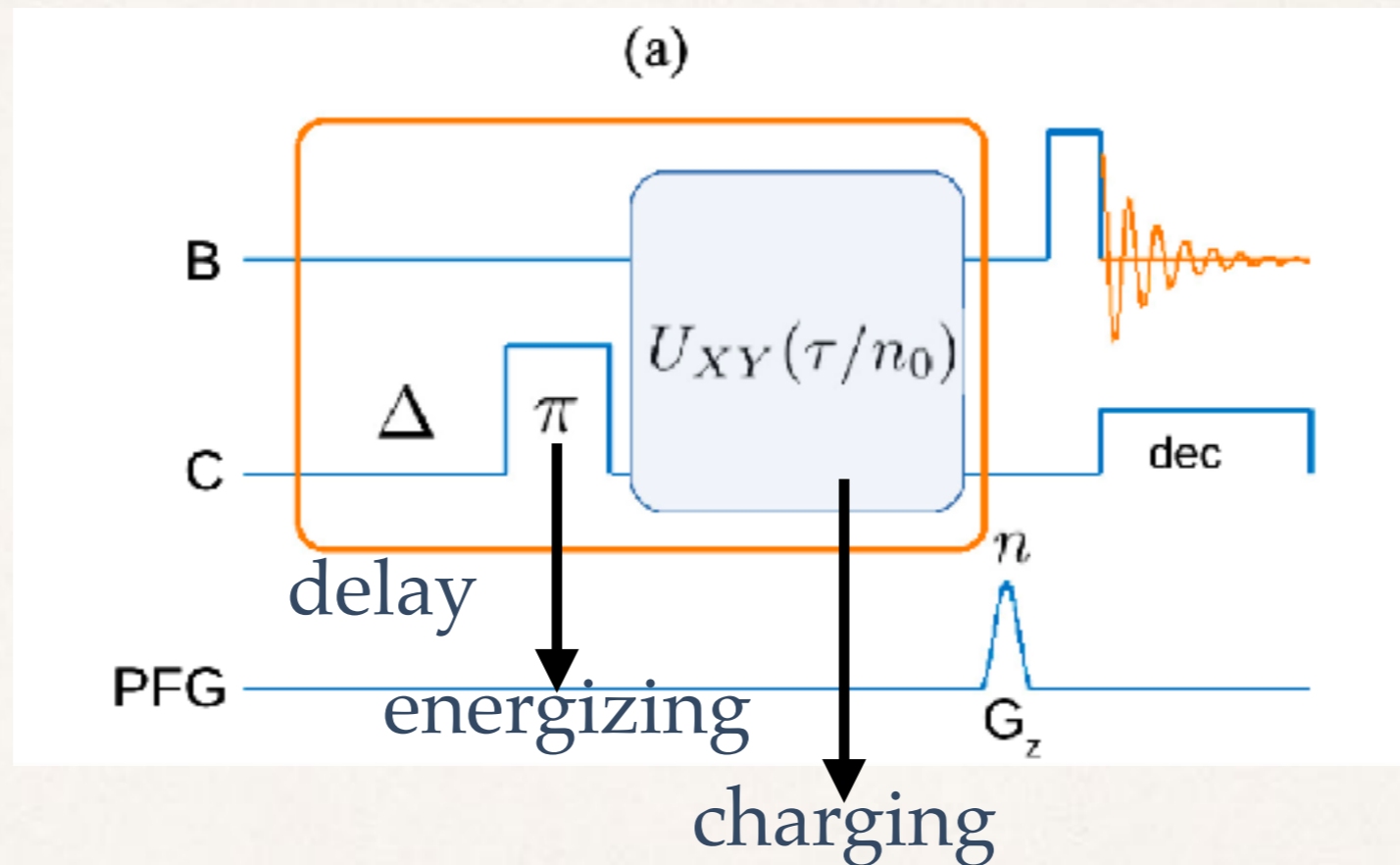
Charger -  $^1\text{H}$  Nuclear Spin

$\hat{S}_{x,y,z}$  - Battery spin

$$\hat{H} = J(\hat{S}_x \hat{I}_x + \hat{S}_y \hat{I}_y)$$

$$\hat{I}_{x,y,z} = \sum_{i=1}^N I_{x,y,z}^i - \text{Battery spin}$$

# Outlook- Experimental Implementation



storage times: battery  $T_1^B \gg T_1^C$  charger

Asymptotic Charging -iteratively re-energizing chargers after delay

# Other Topics of Current Interest

Collective Dissipation of Oscillator Dipoles  
Strongly Coupled to 1-D Electromagnetic Reservoirs

Poster



Ipsita Bar

Indicators of Chaos in the Closed and Open Dicke Model

Poster



Prasad Pawar

Thank you

# *Additional Slides*



tem as  $\rho$ , the maximum amount of work which can be extracted from it during this cyclic process is known as *ergotropy* [9, 199, 279–282], defined as

$$\begin{aligned} \mathcal{E}^1 &= \overline{W}_{U, \max}^1 = \text{Tr} [\rho h_0] - \min_{U(\tau) \in SU(d)} \{ \text{Tr} [U(\tau) \rho U(\tau)^\dagger h_0] \} \\ &= \text{Tr} [\rho h_0] - \text{Tr} [\sigma_\rho h_0], \end{aligned} \quad (139)$$

where the superscript ‘1’ denotes that we are working with a single *cell* or single copy of the system,  $\overline{W}_{U, \max}^1$  denotes the maximum work that can be extracted (negative of the work done,  $\overline{W} = -W$ ) using unitary operations, and  $U(\tau)$  are unitary time-evolution operators acting for the duration of time  $\tau$ . Note that the

state corresponding to  $\rho$ , defined as the state having zero ergotropy. Hence, no energy can be extracted from  $\sigma_\rho$  through cyclic unitary processes, that is

$$\begin{aligned} \Delta E &= -\overline{W}_U^1 \\ &= \text{Tr} [U \sigma_\rho U^\dagger h_0] - \text{Tr} [\sigma_\rho h_0] \geq 0, \quad \forall U \in SU(d), \end{aligned} \quad (140)$$

where a positive value of  $\Delta E$  corresponds to a negative work extraction, or equivalently, a work deposition on the system. In general, it can be shown that a state is passive if it is diagonal in the energy eigen-basis of the system with non-decreasing diagonal elements (populations), when arranged in the order of non-increasing energies, i.e.,  $\sigma = \sum_{j=1}^d s_j |j\rangle \langle j|$  with  $s_j \geq s_{j+1}$  for  $\varepsilon_j \leq \varepsilon_{j+1}$  [197, 198]. Consequently, the passive state  $\sigma_\rho$  which can be attained by means of local unitary operations on  $\rho$  is unique in nature. Unless otherwise mentioned, we shall use the notation  $\sigma_\rho$  to denote the unique passive state corresponding to  $\rho$  in the rest of the article.

Response: We would like to first point out a scaling property present in the dynamics of the charger-battery system described via Eqs. (11) and (12) in the previous version or Eqs. (A3) and (A4) in the modified version of the draft in the resonant case. The solutions only depend on the ratios of quantities  $F/g$  and  $\gamma_C/g$ . While we have chosen  $g = \omega_B$  in the results for the sake of presentation, the results hold good for even smaller coupling  $g$ , as long as we make a proportional change in  $\gamma_C$ . Thus we can always understand our results as being calculated for small values of  $g/\omega_B$  such that the local master equation is valid. Moreover, since this leads even smaller values of  $\gamma_C/\omega_B$  as per the scaling we described, the weak dissipation limit required for usual GKLS master equation is well satisfied. Note that in this case,  $\gamma_C/g$  can still be significantly large but since the system's energy scale is determined by  $\{\omega_B = \omega_C\} \gg g$  in this regime, this does not violate the Born approximation required for the GKLS master equations. To illustrate this, in Fig. R2 we have plotted the charging dynamics for the same values of  $\gamma_C/g$  and  $F/g$  as in Fig. 2 of the paper, but with much smaller value of  $g = 0.1\omega_B$ . It is clear that the results remain exactly the same. In the revised draft, we have added a remark in the

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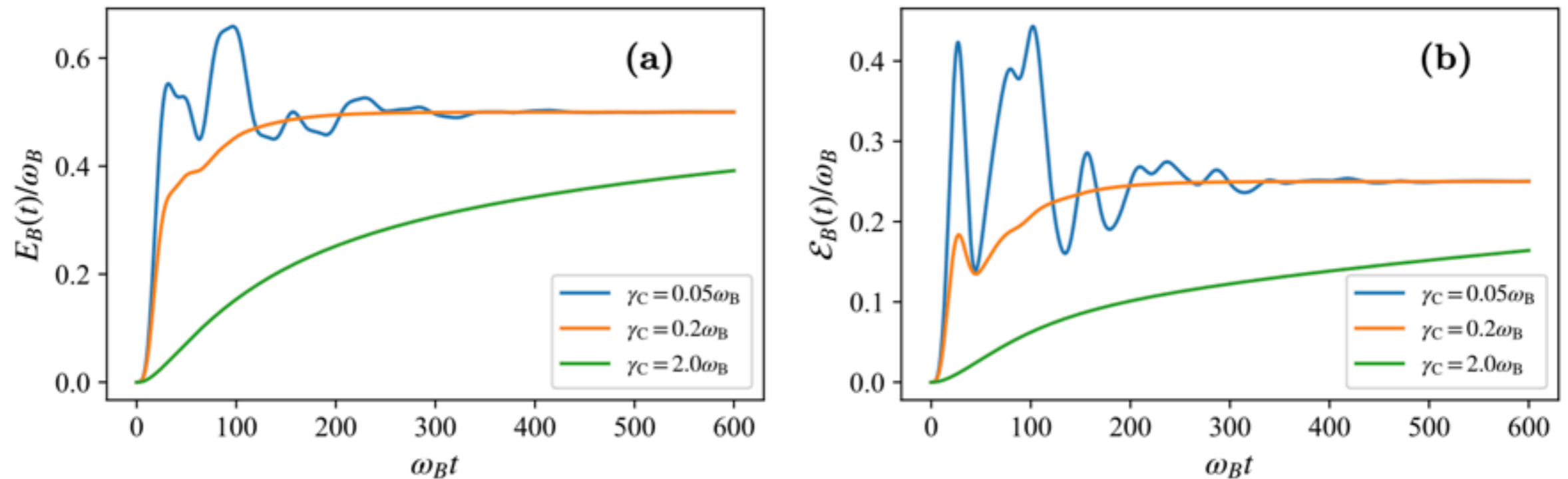


Figure R2: Figure 2 of main paper with  $F = 0.05\omega_B$ , and  $g = 0.1\omega_B$  ( $F/g = 0.5$ ).