

Entropy, Energy, and Temperature in Small Systems: Impact of a Relative Energy Window on Microcanonical Statistical Mechanics

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Challenges in defining entropy in small systems

Microcanonical (NVE) ensemble in statistical mechanics:¹ less investigated but important

An intense debate in microcanonical statistical mechanics:

What w do we choose in $S = -\sum_i p_i \ln p_i = k_B \ln w$?²⁻⁷

Gibbs volume entropy

$$S = k_B \ln \Omega$$

Boltzmann surface entropy

$$S = k_B \ln \left(\underbrace{\frac{\partial \Omega}{\partial E}}_{\omega} \epsilon \right)$$

Isn't energy fixed to a specified value E for an isolated system, and one only looks at degeneracies?

Pathria: “almost every system has some contact with its surroundings, however little it may be; as a result, its energy cannot be defined sharply”

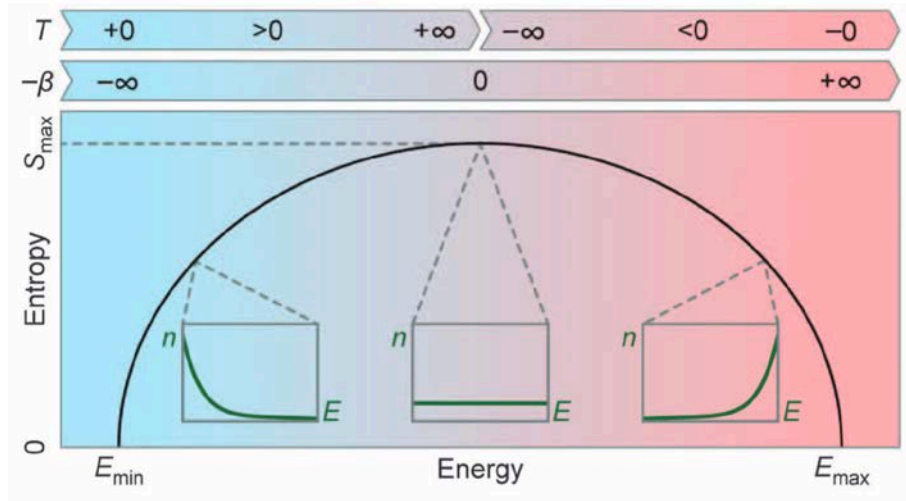
¹Govind Rajan. *J. Chem. Educ.* **2024**, 101, 2448

²Gibbs, J. W. Charles Scribner's Sons, NY, **1902** | ³Hertz, P. *Ann. Phys. (Leipzig)* **1910**, 338, 225 | ⁴Pearson et al. *Phys. Rev. A* **1985**, 3030

⁵Dunkel and Hilbert. *Nat. Phys.* **2014**, 10, 67 | ⁶Frenkel and Warren. *Am. J. Phys.* **2015**, 83, 163 | ⁷Corti et al. *J. Phys. Chem. B* **2023**, 127, 3431

Main point of contention – possibility of negative absolute temperatures

- $T < 0$ K possible in systems with bounded energy as higher energies involve more particles occupying highest-energy level, thus reducing entropy ($\Delta S < 0$) with increase in energy ($\Delta E > 0$)
 - Nuclear spin systems
 - Lasers } **Population inversion**
- Previous measurements based, e.g., on
 - $T = \Delta E / \Delta S$ with ΔS calculated from measured loss of nuclear polarization (*Science* **1994**, 265, 1821)
 - Band imaging of kinetic energy fit to Bose-Einstein distribution (*Science* **2013**, 339, 52)
- Negative absolute temperatures hotter than positive infinity kelvin – temperature scale is discontinuous, but “hotness” $-\beta = -\frac{1}{k_B T}$ is continuous



Issues plaguing Gibbs and Boltzmann entropy

- Gibbs entropy entirely disallows negative absolute temperatures, in disagreement with experiments [Frenkel and Warren. *Am. J. Phys.* **2015**]
- Boltzmann entropy predicts unphysical negative/infinite absolute temperatures for small systems with unbounded energy spectrum [Dunkel and Hilbert. *Nat. Phys.* **2014**]
 - 1D/2D ideal gas of a single particle
 - Single quantum harmonic oscillator
 - Particle in a 1D box

Relative-energy-constrained Boltzmann surface entropy

$$S = k_B \ln \left(\underbrace{\frac{\partial \Omega}{\partial E} \epsilon_r E}_{\omega_r} \right)$$

Justified by the Heisenberg energy-time
uncertainty principle and an eigenstate
thermalization time inversely proportional
to system energy

Predictions for an ideal gas and a particle in a box

At high temperatures/energies, the energy states form a continuum

$$\omega_r = \frac{\partial \Omega}{\partial E} \epsilon_r E \quad S_{Br} = k_B \ln \omega_r \quad T_{Br} = \frac{\partial E}{\partial S_{Br}}$$

At low temperatures/energies, the discreteness of states is apparent

$$\omega_r = \frac{\Delta \Omega}{\Delta E} \epsilon_r E \quad T_{Br} = \frac{\Delta E}{\Delta S_{Br}}$$

System ↓ \ Entropy →	Degeneracy	Gibbs	Boltzmann	Proposed	Canonical
3D ideal gas (high T)	not defined	$E = \frac{3N}{2} k_B T_G$	$E = \left(\frac{3N}{2} - 1\right) k_B T_B$	$E = \frac{3N}{2} k_B T_{Br}$	$E = \frac{3N}{2} k_B T_C$
Particle in a 1D box (high T)	$T_D \rightarrow \infty$	$E = \frac{k_B T_G}{2}$	$E = -\frac{k_B T_B}{2}$	$E = \frac{k_B T_{Br}}{2}$	$E = \frac{k_B T_C}{2}$

Boltzmann/degeneracy entropy seem to fail for these two systems

The new definition agrees well with canonical predictions, without being forced to

But Gibbs' definition seems to agree too?

Examining a quantum harmonic oscillator

Unbounded energy spectrum

$$E_n = h\nu \left(n + \frac{1}{2} \right)$$

Total number of states

$$\Omega = 1 + n = \frac{E}{h\nu} + \frac{1}{2}$$

Gibbs entropy misses zero-point energy

$$\frac{E - E_0}{h\nu} = \frac{3 - \exp\left(\frac{h\nu}{k_B T_G}\right)}{2 \exp\left(\frac{h\nu}{k_B T_G}\right) - 2}$$

Relative-energy-constrained

Boltzmann entropy correct $\forall T \in [0, \infty)$

$$\frac{E - E_0}{h\nu} = \frac{1}{\exp\left(\frac{h\nu}{k_B T_{Br}}\right) - 1}$$

Two-level system with negative temperatures

Bounded energy spectrum

$$E = N_1 E_1; \quad E_{\max} = N E_1$$

Proposed partition function

$$\omega_r = \frac{\Delta\Omega}{\Delta E} \epsilon_r E = \frac{N!}{(N_0 - 1)! (N_1 + 1)!} \frac{\epsilon_r E}{E_1}$$

Density of states

$$\Omega = \sum_{N'_0=1}^{N_1} \frac{N!}{N'_0! N'_1!}$$

**Gibbs entropy misses negative temperatures;
Boltzmann entropy close to degeneracy**

$$T_B = \frac{\Delta E}{\Delta S_B} = \frac{E_1}{k_B \ln \left(\frac{N_0 - 1}{N_1 + 2} \right)}$$

**Proposed entropy again agrees with
canonical predictions**

$$T_{Br} = \frac{\Delta E}{\Delta S_{Br}} = \frac{E_1}{k_B \ln \left(\frac{(N_1 + 1)(N_0 - 1)}{N_1(N_1 + 2)} \right)}$$

Is there precedent for a relative-energy window?

- An energy tolerance that makes the microcanonical and canonical (C) ensembles agree:
 - Proposed by Gurarie (*Am. J. Phys.* **2007**, 75, 747)
 - Used recently by Burke and Haque (*Phys. Rev. E* **2023**, 107, 034125)

$$\epsilon = \sqrt{2\pi k_B T_C^2 C_C} \quad \text{versus} \quad \epsilon = \epsilon_r E$$

- Works for an ideal gas, but gives incorrect results for a harmonic oscillator

$$T = \frac{h\nu}{S(E + h\nu - E_0) - S(E - E_0)} = \frac{h\nu}{k_B \ln \left(\sqrt{\frac{\frac{h\nu}{E + h\nu - E_0} + 1}{\frac{h\nu}{E - E_0} + 1}} \left(\frac{E + h\nu - E_0}{E - E_0} \right) \right)}$$

versus

$$T = \frac{h\nu}{k_B \ln \left(\frac{E - E_0 + h\nu}{E - E_0} \right)}$$

Conclusions

$$\sigma_E \Delta t \geq \frac{\hbar}{2}$$

$$\Delta t \sim \frac{h}{E}$$



$$\sigma_E \sim E$$

(1) Entropy definition for isolated systems considering a relative energy window

(2) Motivated by Heisenberg uncertainty principle and a thermalization time inversely proportional to energy

(3) Preserves zero-point energy of a quantum harmonic oscillator and $T > 0$ K for unbounded energy systems

(4) Predicts $T < 0$ K when appropriate; closest to canonical ensemble

Thank you for listening

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