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## Lecture 2: Quantum state reduction and feedback stabilization

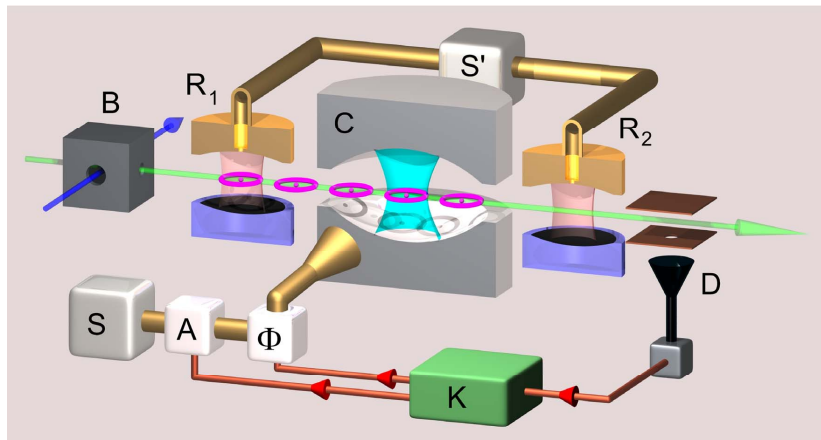
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# The closed-loop QED experiment<sup>2</sup>

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Sampling time ( $\sim 100 \mu\text{s}$ ) long enough for feedback computations.

<sup>1</sup>Courtesy of Igor Dotsenko

<sup>2</sup>C. Sayrin et al., Nature, 1-September 2011

## LKB photon box: Controlled Markov chain

- ▶  $\rho_{k+1} = \frac{D(\alpha_k)M_{\mu_k}\rho_kM_{\mu_k}^\dagger D(-\alpha_k)}{\text{tr}(M_{\mu_k}\rho_kM_{\mu_k}^\dagger)}$  with  $\mu_k = g$  or  $e$ .
- ▶ Maximum of photon number in the cavity is  $n^{\max}$
- ▶  $M_g = \cos(\frac{\phi_R + \phi}{2} + N\phi)$  and  $M_e = \sin(\frac{\phi_R + \phi}{2} + N\phi)$ , where  $N = a^\dagger a$  is the photon number operator,  $N = (\text{diag}(n))_{0 \leq n \leq n^{\max}}$ ,  $a$  is an upper-triangular matrix with  $(\sqrt{n})_{1 \leq n \leq n^{\max}}$  as upper diagonal
- ▶  $D(\alpha_k)$  is the displacement operator describing the coherent pulse injection  $D(\alpha_k) = \exp(\alpha_k(a^\dagger - a))$  with  $\alpha_k$  a real parameter corresponding to the control input

**Remark.** Measurement operators are diagonal in the basis  $|n\rangle \langle n|$ , i.e., Quantum Non-Demolition (QND) measurements

# Feedback stabilization

**Aim.** Stabilize a particular photon number state  $\bar{\rho} = |\bar{n}\rangle \langle \bar{n}|$ .

Define  $V(\rho_k) = 1 - \text{tr}(\bar{\rho}\rho_k)$ .

**Theorem (Mirrahimi, Dotsenko, Rouchon, 2009).** Take the feedback controller

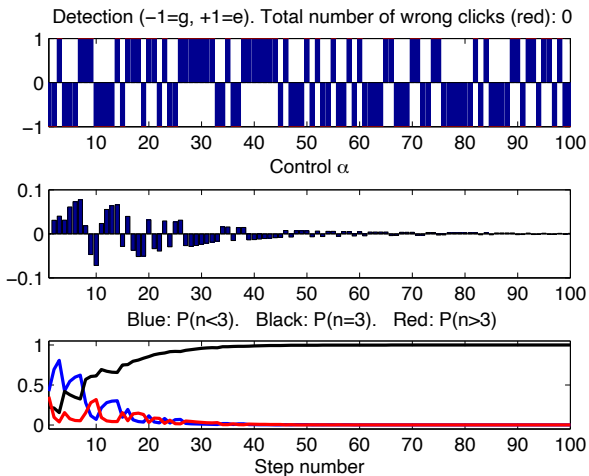
$$u_n = \begin{cases} \operatorname{argmax}_{\alpha \in [-\bar{\alpha}, \bar{\alpha}]} \operatorname{tr}(\bar{\rho} D(\alpha) \rho_{k+1/2} D(-\alpha)) & \text{if } V(\rho) > 1 - \epsilon \\ c_1 \operatorname{tr}([\bar{\rho}, \mathbf{a}^\dagger - \mathbf{a}] \rho_{k+1/2}) & \text{if } V(\rho) \leq 1 - \epsilon \end{cases}$$

Assume that for all  $n \in \{0, \dots, n^{\max}\}$ , we have that  $\frac{\phi_R + \phi}{2} + n\phi \neq 0 \pmod{\pi/2}$ . For small enough  $c_1 > 0$  and  $\epsilon > 0$ , the trajectories converge almost surely towards the target Fock state  $\bar{\rho}$ .

## Sketch of the proof

- ▶ **Step 1:** the trajectories starting within the set  $\{\rho \mid V(\rho) > 1 - \epsilon\}$  reach in one step the set  $\{\rho \mid V(\rho) \leq 1 - 2\epsilon\}$  with probability one
- ▶ **Step 2:** the trajectories starting within the set  $\{\rho \mid V(\rho) \leq 1 - 2\epsilon\}$  will never hit the set  $\{\rho \mid V(\rho) > 1 - \epsilon\}$  with a uniformly non-zero probability  $p > 0$
- ▶ **Step 3:** quantum trajectories converge toward the state  $\bar{\rho}$

# Monte-Carlo simulations



A better convergence

# The Markov model

The random evolution of the state  $\rho_k$  at time step  $k$ ,

$$\rho_{k+1} = \mathbb{U}_{u_k}(\mathbb{M}_{\mu_k}(\rho_k))$$

- ▶ Consider a finite dimensional quantum system (the Hilbert space  $\mathcal{H} \in \mathbb{C}^d$  is of dimension  $d > 0$ ) and  $\mathcal{D}(\mathcal{H}) := \{\rho \in \mathbb{C}^{d \times d} \mid \rho = \rho^\dagger, \text{tr}(\rho) = 1, \rho \geq 0\}$ ,
- ▶  $u_k \in \mathbb{R}$  is the control at step  $k$ ,
- ▶  $\mu_k$  is a random variable taking values  $\mu$  in  $\{1, \dots, m\}$  with probability  $p_{\mu, \rho_k} = \text{tr}(M_\mu \rho_k M_\mu^\dagger)$ ,



- ▶ The set of Kraus operator satisfies the constraint

$$\sum_{\mu=1}^m M_{\mu}^{\dagger} M_{\mu} = I,$$

- ▶  $\mathbb{U}_U$  is the super-operator

$\mathbb{U}_U : \mathcal{D}(\mathcal{H}) \ni \rho \mapsto U_U \rho U_U^{\dagger} \in \mathcal{D}(\mathcal{H})$  where  $U_U = \exp(-iU H)$ , where  $H$  is an Hermitian operator  $H \in \mathbb{C}^{d \times d}$  with  $H^{\dagger} = H$ ,

- ▶ For each  $\mu$ ,  $\mathbb{M}_{\mu}$  is the super-operator

$\mathbb{M}_{\mu} : \rho \mapsto \frac{M_{\mu} \rho M_{\mu}^{\dagger}}{\text{tr}(M_{\mu} \rho M_{\mu}^{\dagger})} \in \mathcal{D}(\mathcal{H})$  defined for  $\rho \in \mathcal{D}(\mathcal{H})$  such

that  $p_{\mu, \rho} = \text{tr}(M_{\mu} \rho M_{\mu}^{\dagger}) \neq 0$ .

## Assumption 1 (QND measurement)

The measurement operators  $M_\mu$  are diagonal in the same ortho-normal basis  $\{|n\rangle \mid n \in \{1, \dots, d\}\}$ , therefore

$$M_\mu = \sum_{n=1}^d c_{\mu,n} |n\rangle \langle n| \text{ with } c_{\mu,n} \in \mathbb{C}.$$

Since  $\sum_{\mu=1}^m M_\mu^\dagger M_\mu = I$ , we have  $\sum_{\mu=1}^m |c_{\mu,n}|^2 = 1$  for all  $n \in \{1, \dots, d\}$ .

## Assumption 2

For all  $n_1 \neq n_2$  in  $\{1, \dots, d\}$ , there exists a  $\mu \in \{1, \dots, m\}$  such that  $|c_{\mu,n_1}|^2 \neq |c_{\mu,n_2}|^2$ .

## A fundamental theorem

**Theorem (Kushner, 1971).** Let  $X_k$  be a Markov chain on the compact state space  $S$ . Suppose, there exists a non-negative function  $V(X)$  satisfying

$$\mathbb{E}(V(X_{k+1})|X_k) - V(X_k) = Q(X_k),$$

where  $Q(X)$  is a positive continuous function of  $X$ , then the  $\omega$ -limit set  $\Omega$  (in the sense of almost sure convergence) of  $X_k$  is included in the following set

$$I := \{X \mid Q(X) = 0\}.$$

# Convergence of the open loop dynamics

When  $u_k = 0$ ,  $\forall k$ , the dynamics is simply given by

$$\rho_{k+1} = \mathbb{M}_{\mu_k}(\rho_k) = \frac{M_\mu \rho_k M_\mu^\dagger}{\text{tr}(M_\mu \rho_k M_\mu^\dagger)},$$

with the operator  $M_\mu = \text{diag}(c_{\mu,n})_{1 \leq n \leq d}$ .

**Theorem (A., Rouchon, Mirrahimi, 2011)** Consider a Markov process  $\rho_k$  obeying the dynamics given in above with an initial condition  $\rho_0$  in  $\mathcal{D}(\mathcal{H})$ . Then

- ▶ with probability one,  $\rho_k$  converges to one of the  $d$  states  $|n\rangle \langle n|$  with  $n \in \{1, \dots, d\}$ .
- ▶ the probability of convergence towards the state  $|n\rangle \langle n|$  depends only on the initial condition  $\rho_0$  and is given by

$$\text{tr}(\rho_0 |n\rangle \langle n|) = \langle n | \rho_0 | n \rangle .$$

## Elements of a proof:

- ▶ **Step 1.** Taking the following Lyapunov function

$$V(\rho) := - \sum_{n=1}^d f(\text{tr}(|n\rangle \langle n| \rho)), \quad \text{where } f(x) = \frac{x^2}{2}.$$

- ▶ **Step 2.** Show that  $\text{tr}(|n\rangle \langle n| \rho_{k+1})$  is a martingale.

$$\mathbb{E}(\text{tr}(|n\rangle \langle n| \rho_{k+1}) | \rho_k) = \text{tr}(|n\rangle \langle n| \rho_k).$$

- ▶ **Step 3.** The function  $-f$  being concave, thus  $V(\rho)$  is a super-martingale

$$\mathbb{E}(V(\rho_{k+1}) | \rho_k) \leq V(\rho_k).$$

More precisely, we obtain

$$\begin{aligned} \mathbb{E}(V(\rho_{k+1})|\rho_k) - V(\rho_k) = \\ -\frac{1}{4} \sum_{n=1}^d \sum_{\mu, \nu} \text{tr}(M_\mu \rho_k M_\mu^\dagger) \text{tr}(M_\nu \rho_k M_\nu^\dagger) \times \\ \left( \frac{|c_{\mu, n}|^2 \langle n | \rho_k | n \rangle}{\text{tr}(M_\mu \rho_k M_\mu^\dagger)} - \frac{|c_{\nu, n}|^2 \langle n | \rho_k | n \rangle}{\text{tr}(M_\nu \rho_k M_\nu^\dagger)} \right)^2 := Q_1(\rho_k). \end{aligned}$$

The  $\omega$ -limit set  $\Omega$  is a subset of the following set

$$\{\rho_\infty \mid \frac{|c_{\mu, n}|^2 \langle n | \rho_\infty | n \rangle}{\text{tr}(M_\mu \rho_\infty M_\mu^\dagger)} - \frac{|c_{\nu, n}|^2 \langle n | \rho_\infty | n \rangle}{\text{tr}(M_\nu \rho_\infty M_\nu^\dagger)} = 0\},$$

$\implies$  Assume  $\exists \bar{n}_1 \neq \bar{n}_2 \in \{1, \dots, d\} : \langle \bar{n}_1 | \rho_\infty | \bar{n}_1 \rangle > 0$  and  $\langle \bar{n}_2 | \rho_\infty | \bar{n}_2 \rangle > 0$ . Then

$$\forall \mu, \quad |c_{\mu, \bar{n}_1}|^2 = |c_{\mu, \bar{n}_2}|^2.$$

By Assumption 2,  $\exists \bar{\mu} \in \{1, \dots, m\} : |c_{\bar{\mu}, \bar{n}_1}|^2 \neq |c_{\bar{\mu}, \bar{n}_2}|^2$ . Then there exists a **unique**  $\bar{n}$  such that  $\langle \bar{n} | \rho_\infty | \bar{n} \rangle \neq 0$ .

- ▶ **Step 4.** The probability measure of the random variable  $\rho_k$  converges to  $\sum_{n=1}^d p_n \delta(|n\rangle \langle n|)$ , where,
  - ▶  $\delta(|n\rangle \langle n|)$  denotes the dirac distribution at  $|n\rangle \langle n|$ .
  - ▶  $p_n$  is the probability of convergence towards  $|n\rangle \langle n|$ .

In particular,

$$\mathbb{E}(\text{tr}(\rho_k |n\rangle \langle n|)) \longrightarrow p_n.$$

But  $\text{tr}(|n\rangle \langle n| \rho_k)$  is a martingale and  $\mathbb{E}(\text{tr}(|n\rangle \langle n| \rho_k)) = \mathbb{E}(\text{tr}(|n\rangle \langle n| \rho_0))$ . Then  $p_n = \langle n | \rho_0 | n \rangle$ .

# Feedback Stabilization

Recall:

$$\rho_{k+1} = \mathbb{U}_{u_k}(\mathbb{M}_{\mu_k}(\rho_k)).$$

**Goal.** Design a feedback law that globally stabilizes the Markov chain  $\rho$  towards a target state  $|\bar{n}\rangle \langle \bar{n}|$  with  $\bar{n} \in \{1, \dots, d\}$ .

**Astuce.** Construction of a strict Lyapunov function which is based on the connectivity of the graph attached to H and inverting a Laplacian matrix derived from H.



Lemma (A., Rouchon, Mirrahimi, 2011). Consider the  $d \times d$  real matrix  $R^H$ , defined by

$$R_{n_1, n_2}^H = 2|\langle n_1 | H | n_2 \rangle|^2 - 2\delta_{n_1, n_2} \langle n_1 | H^2 | n_2 \rangle.$$

Then  $R^H$  is a Laplacian matrix. Assume the graph  $G^H$  of the Laplacian matrix  $R^H$  is connected, Then, for any positive reals  $\lambda_n$ ,  $n \in \{1, \dots, d\}$ ,  $n \neq \bar{n}$ , there exists a vector  $\sigma = (\sigma_n)_{n \in \{1, \dots, d\}}$  of  $\mathbb{R}^d$  such that  $R^H \sigma = \lambda$  where  $\lambda$  is the vector of  $\mathbb{R}^d$  of components  $\lambda_n$  for  $n \neq \bar{n}$  and  $\lambda_{\bar{n}} = -\sum_{n \neq \bar{n}} \lambda_n$ .

Now choose  $\sigma$  by  $R^H \sigma = \lambda$  and construct the function  $W_0(\rho) = \sum_{n=1}^d \sigma_n \langle n | \rho | n \rangle$ , we have

$$\left. \frac{\partial^2 W_0(\mathbb{U}_u(|n\rangle\langle n|))}{\partial u^2} \right|_{u=0} = \sum_{l=1}^d \sigma_l \text{tr}([H, |n\rangle\langle n|][H, |l\rangle\langle l|]).$$

Since,

$$\text{tr}([H, |n\rangle\langle n|][H, |l\rangle\langle l|]) = R_{n,l}^H$$

Thus

$$\left. \frac{\partial^2 W_0(\mathbb{U}_u(|n\rangle\langle n|))}{\partial u^2} \right|_{u=0} = \sum_{l=1}^d R_{n,l}^H \sigma_l = \lambda_n.$$

Then  $W_0$  is convex for  $n \neq \bar{n}$  and concave for  $n = \bar{n}$ .

# The global stabilizing feedback

Consider the controlled Markov chain  $\rho_{k+1} = \mathbb{U}_{u_k}(\mathbb{M}_{\mu_k}(\rho_k))$ .  
Construct  $W_\epsilon(\rho) = \sum_{n=1}^d \left( \sigma_n \langle n | \rho | n \rangle + \epsilon (\langle n | \rho | n \rangle)^2 \right)$ , where  
the parameter  $\epsilon > 0$  should be not too large to ensure that  
 $\forall n \in \{1, \dots, d\} \setminus \{\bar{n}\}, \lambda_n + \epsilon (\langle n | H | n \rangle)^2 - \langle n | H^2 | n \rangle > 0$ .

**Theorem (A., Rouchon, Mirrahimi, 2011)** Denote by  
 $\rho_{k+\frac{1}{2}} = \mathbb{M}_{\mu_k}(\rho_k)$ . Take  $\bar{u} > 0$  and consider the following  
feedback law

$$u_k = K(\rho_{k+\frac{1}{2}}) = \operatorname{argmax}_{u \in [-\bar{u}, \bar{u}]} \left( W_\epsilon(\mathbb{U}_u(\rho_{k+\frac{1}{2}})) \right),$$

Then, for any  $\rho_0 \in \mathcal{D}(\mathcal{H})$ , the closed-loop trajectory  $\rho_k$   
converges almost surely to the pure state  $|\bar{n}\rangle \langle \bar{n}|$ .

# Elements of a proof

$$\mathbb{E}(W_\epsilon(\rho_{k+1})|\rho_k) - W_\epsilon(\rho_k) := Q_1(\rho_k) + Q_2(\rho_k).$$

With

$$Q_1(\rho_k) := \sum_{\mu} p_{\mu, \rho_k} \left( W_\epsilon(\mathbb{M}_\mu(\rho_k)) - W_\epsilon(\rho_k) \right),$$

and

$$Q_2(\rho_k) := \sum_{\mu} p_{\mu, \rho_k} \left( \max_{u \in [-\bar{u}, \bar{u}]} \left( W_\epsilon(\mathbb{U}_u(\mathbb{M}_\mu(\rho_k))) \right) - W_\epsilon(\mathbb{M}_\mu(\rho_k)) \right).$$

These functions are both positive continuous functions of  $\rho_k$ .

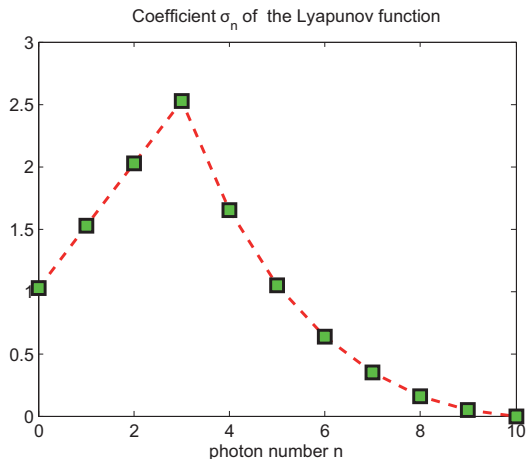
Then, the  $\omega$ -limit set  $\Omega$  is included in the following set

$$\{\rho \in \mathcal{D}(\mathcal{H}) \mid Q_1(\rho) = 0\} \cap \{\rho \in \mathcal{D}(\mathcal{H}) \mid Q_2(\rho) = 0\}.$$

▶  $Q_1(\rho) = 0 \implies \rho = |n\rangle \langle n|$

▶  $(Q_1(\rho) = 0) + (Q_2(\rho) = 0) \implies \rho = |\bar{n}\rangle \langle \bar{n}|.$

## The control Lyapunov function used for the LKB photon box



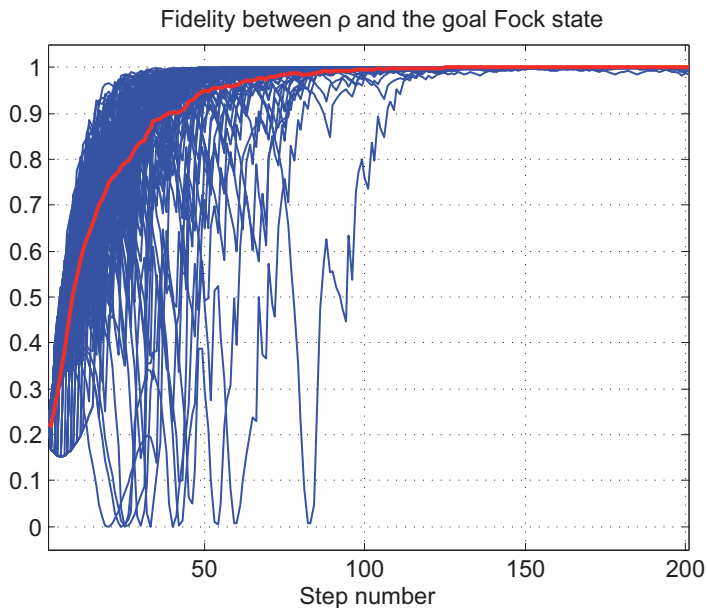
$$W_\epsilon(\rho) = \sum_{n=0}^9 \left( \epsilon \langle n | \rho | n \rangle^2 + \sigma_n \langle n | \rho | n \rangle \right) \text{ with } \epsilon = 0 \text{ in fact}$$

# LKB photon box

- ▶ Here  $\mathcal{H} = \mathbb{C}^{n^{\max}+1}$
- ▶ Take  $\theta/\pi$  irrational (then assumption there exists  $\mu$  such that  $|c_{\mu,n_1}|^2 \neq |c_{\mu,n_2}|^2$  is verified)
- ▶ The Hamiltonian is  $H = i(a^\dagger - a)$
- ▶ The graphs  $G^H$  is connected (and consequently  $R^H$ )
- ▶  $R_{nn}^H = 4n + 2$ ,  $R_{n-1,n}^H = -2n$ , and  $R_{n+1,n}^H = -2n - 2$
- ▶ Take  $\bar{n} = 3$
- ▶ Take  $\lambda_n = 1$  for any  $n \neq \bar{n}$
- ▶ Take  $\epsilon < \frac{1}{2n^{\max}+1}$
- ▶ For  $u$  small enough, by Baker-Campbell-Hausdorff:

$$\mathbb{U}_u(\rho) = \rho - iu[H, \rho] - u^2[H, [H, \rho]] + \mathcal{O}(|u|^3).$$

# Convergence toward the goal state $|3\rangle \langle 3|$



# LKB photon box: imperfections <sup>3</sup>

$$\rho_{n+1} = \frac{\Phi_{i_n}(\rho_n)}{\text{tr}(\Phi_{i_n}(\rho_n))}$$

with  $\Phi_i(\rho) = \sum_{i \in \mathcal{Y}} \eta_{i,j} V_j \rho V_j^\dagger$ . The elements  $\eta_{ij}$  of the correlation matrix for  $i, j \in \{no, g, e, gg, ge, ee\}$  is

$i \setminus j$	$no$	$g$	$e$	$gg$	$ee$	$ge$ or $eg$
$no$	1	$1 - \varepsilon_d$	$1 - \varepsilon_d$	$(1 - \varepsilon_d)^2$	$(1 - \varepsilon_d)^2$	$(1 - \varepsilon_d)^2$
$g$	0	$\varepsilon_d(1 - \eta_g)$	$\varepsilon_d \eta_e$	$2\varepsilon_d(1 - \varepsilon_d)(1 - \eta_g)$	$2\varepsilon_d(1 - \varepsilon_d)\eta_e$	$\varepsilon_d(1 - \varepsilon_d)(1 - \eta_g + \eta_e)$
$e$	0	$\varepsilon_d \eta_g$	$\varepsilon_d(1 - \eta_e)$	$2\varepsilon_d(1 - \varepsilon_d)\eta_g$	$2\varepsilon_d(1 - \varepsilon_d)(1 - \eta_e)$	$\varepsilon_d(1 - \varepsilon_d)(1 - \eta_e + \eta_g)$
$gg$	0	0	0	$\varepsilon_d^2(1 - \eta_g)^2$	$\varepsilon_d^2 \eta_e^2$	$\varepsilon_d^2 \eta_e(1 - \eta_g)$
$ge$	0	0	0	$2\varepsilon_d^2 \eta_g(1 - \eta_g)$	$2\varepsilon_d^2 \eta_e(1 - \eta_e)$	$\varepsilon_d^2((1 - \eta_g)(1 - \eta_e) + \eta_g \eta_e)$
$ee$	0	0	0	$\varepsilon_d^2 \eta_g^2$	$\varepsilon_d^2(1 - \eta_e)^2$	$\varepsilon_d^2 \eta_g(1 - \eta_e)$

$$\begin{aligned} V_{no} &= \sqrt{p_0} I, & V_g &= \sqrt{p_1} \cos \phi_N, & V_e &= \sqrt{p_1} \sin \phi_N \\ V_{gg} &= \sqrt{p_2} \cos^2 \phi_N, & V_{ge} &= V_{eg} = \sqrt{p_2} \cos \phi_N \sin \phi_N, \\ V_{ee} &= \sqrt{p_2} \sin^2 \phi_N, \end{aligned}$$

<sup>3</sup>See the proof of robustness in Amini et al., Automatica, 2013.



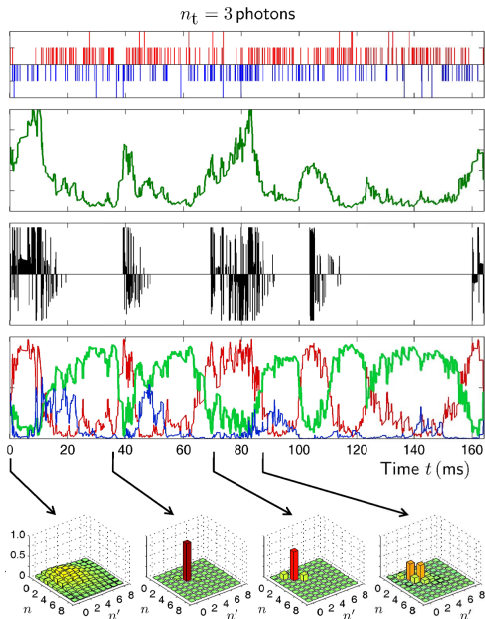
# Experimental data

## Stabilization around 3-photon state

An **open-loop trajectory** starting from coherent state with an average of 3 photons relaxes towards vacuum (decoherence due to finite photon life time around 70 ms) and a **closed-loop trajectory**

Detection efficiency 40%  
Detection error rate 10%  
Delay 4 sampling periods

Truncation to 9 photons



Thank you !