

Two applications of the bootstrap in QCD
Martin Kruczenski (Purdue University)

"Rigidity" of a theory or mathematical structure
↳ Determined by few parameters

- o o o o
 - o o o o
 - o o o o
 - o crystal
 - o analytic function
 - o integrable theory
- Can sometimes be determined from general constraints

Interesting ~~of~~ constraint: Positivity. ~~simple~~

$$A_{ij} = \langle x_i x_j \rangle$$

$$A \succeq 0$$

↑ random
variables

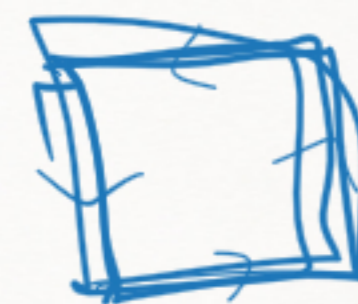
$$\sum_i A_{ij} \xi_j =$$

$$= \langle (\xi_i x_i) (\xi_j x_j) \rangle \geq 0$$

Single plaquette model

$$Z = \int dU \ e^{-\frac{N}{2\lambda} \text{Tr}(U+U^\dagger)}$$

↓ $SU(N)$



$$\underline{W_n} = \frac{1}{N} \frac{1}{Z} \int dU \operatorname{Tr}(U^n) e^{-S}$$

$$W_1 = u = \begin{cases} 1 - \lambda/2 & \lambda \leq 1 \\ 1/(2\lambda) & \lambda \geq 1 \end{cases} \quad \begin{matrix} 3^{\text{rd}} \\ \text{order} \end{matrix}$$



Gross-Witten

Wadia '80

$$A \quad \text{Tr} (A^+ A) = \sum_{ij} |a_{ij}|^2 \geq 0$$

$$A = \sum_{n=0}^L c_n U^n \Rightarrow \text{Tr} \sum_{n,m=0}^L c_m^* c_n (U^+)^m U^n \geq 0$$

$$\sum_{n,m=0}^L c_m^* c_n \langle \text{Tr} U^{-m+n} \rangle \geq 0$$

$\sqrt{\frac{1}{n-m}}$
 $\sqrt{\frac{1}{n-m}}$

$$\rho = \begin{pmatrix} w_0 & w_1 & w_2 & w_3 \\ w_1 & w_0 & w_1 & w_2 \\ w_2 & w_1 & w_1 & w_1 \\ w_3 & w_2 & w_1 & w_1 \end{pmatrix} \succeq 0$$

$$w_0 = \frac{1}{N} \text{Tr} \rho = 1$$

$$\begin{vmatrix} 1 & w_n \\ w_n & 1 \end{vmatrix} \quad |w_n| \leq 1$$

loop equation

$$f = \int dU \ (U^n)_{ab} \ e^{-\frac{N}{2\lambda} \text{Tr}(U+U^\dagger)}$$

$$U \rightarrow (1+i\varepsilon)U$$

$$U^\dagger \rightarrow U^\dagger(1-i\varepsilon)$$

$\varepsilon \Rightarrow$ hermitian
traceless

$$\delta_\varepsilon f = 0$$

$$(1+i\varepsilon)U^n + U(1+i\varepsilon)U^{n-1} + \dots + U^{n-1}(1+i\varepsilon)U$$

$$i(\varepsilon U^n + U\varepsilon U^{n-1} + \dots + U^{n-1}\varepsilon U) \leftarrow$$

$$\delta_\epsilon f = i \int dU e^{-S} \left[\underline{\underline{(\epsilon U^n + U \epsilon U^{n-1} + \dots + U^{n-1} \epsilon U)_{ab}}} + U_{ab}^n \left(-\frac{N}{2\lambda} \text{Tr}(\epsilon U - \epsilon U^\dagger) \right) \right] = 0$$

$$\underbrace{\epsilon_{cd} B^{cd}} = 0$$

$$B^{cd} = \frac{1}{N} \text{Tr} B \delta^{cd}$$

$$\left\langle \cancel{\epsilon_{cd}} \left(\underbrace{\delta_{ac}}_{\downarrow} \underbrace{U_{db}}_{\uparrow} + \underbrace{U_{ac}}_{\downarrow} \underbrace{U_{cd}}_{\uparrow} + \dots + \underbrace{U_{ac}}_{\downarrow} \underbrace{U_{db}}_{\uparrow} - \frac{N}{2\lambda} \underline{\underline{U_{ab}^n (U_{dc} - U_{dc}^\dagger)}}} \right) \right\rangle = 0 \quad (N \rightarrow \infty)$$

$$N \langle \text{Tr} U^n \rangle + \langle \text{Tr} U \rangle \langle \text{Tr} U^{n-1} \rangle + \dots + \langle \text{Tr} U^{n-1} \rangle \langle \text{Tr} U \rangle$$

$$= \frac{N}{2\lambda} (\text{Tr} U^{n+1} - \text{Tr} U^{n-1}) = 0$$

$$2W_n + 2W_1 W_{n-1} + \dots + 2W_{n-1} W_1 - \frac{1}{2\lambda} (W_{n+1} - W_{n-1}) = 0$$

$$W_{n+1} - W_{n-1} = 2\lambda W_1 + 2\lambda \sum_{p=1}^{n-1} W_p W_{n-p}$$

$$\begin{pmatrix} w_0 & w_1 & w_E & - \\ L_1 & & & \\ & & & w_c \\ & & & w_1 \\ & & w_n & w_0 \end{pmatrix} \geq 0$$

- o) Henry Lin
- oo) P. Anderson & M.K.

Positivity

$$A_{ij} = \langle x_i, x_j \rangle$$

$$\sum_i \langle x_i, x_j \rangle \xi_j = \langle (\xi_i, x_i)^2 \rangle \geq 0$$

$$\boxed{A \succeq 0}$$

$$\boxed{\omega_a}$$

$$A = A^{(0)} + A_a^{(1)} \omega_a$$

$$F = \sum_a \alpha_a \omega_a$$

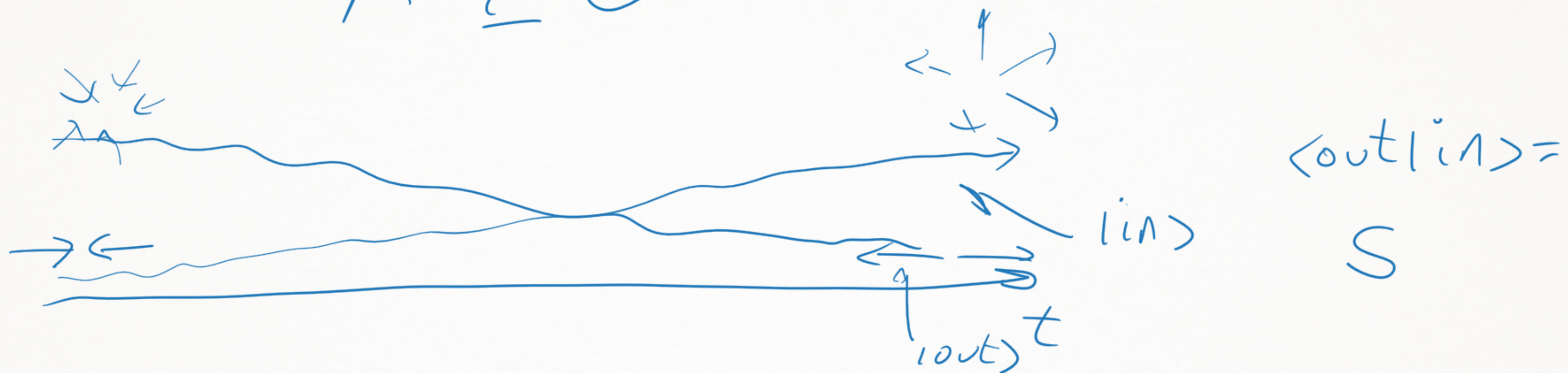
MOSEK
Gurobi
sdpa → sdpb



$$\{|\psi_i\rangle\} \quad A_{ij} = \overline{\langle \psi_i | \psi_j \rangle}$$

$$\sum_i^* A_{ij} \xi_j = \|\xi_i |\psi_i\rangle\|^2 \geq 0$$

$$A \geq 0$$



$$\begin{matrix} \langle in | & |in\rangle \\ \langle out | & |out\rangle \end{matrix} \begin{pmatrix} \mathbb{1} & S^\dagger \\ S & \mathbb{1} \end{pmatrix} \geq 0$$

$$\begin{pmatrix} 1 & S \\ S^\dagger & 1 \end{pmatrix} \geq 0$$

$$\underline{\underline{|S|^2 \leq 1}}$$

$$\begin{matrix} \langle in | \\ \langle out | \\ \langle 0 | \theta^\dagger \end{matrix} \begin{pmatrix} |in\rangle & |out\rangle & |0\rangle \\ \mathbb{1} & S^\dagger & \langle in | \theta | 0 \rangle \\ S & \mathbb{1} & \langle out | \theta | 0 \rangle \\ \langle 0 | \theta^\dagger | 0 \rangle & & \end{pmatrix} \geq 0$$

form-factor

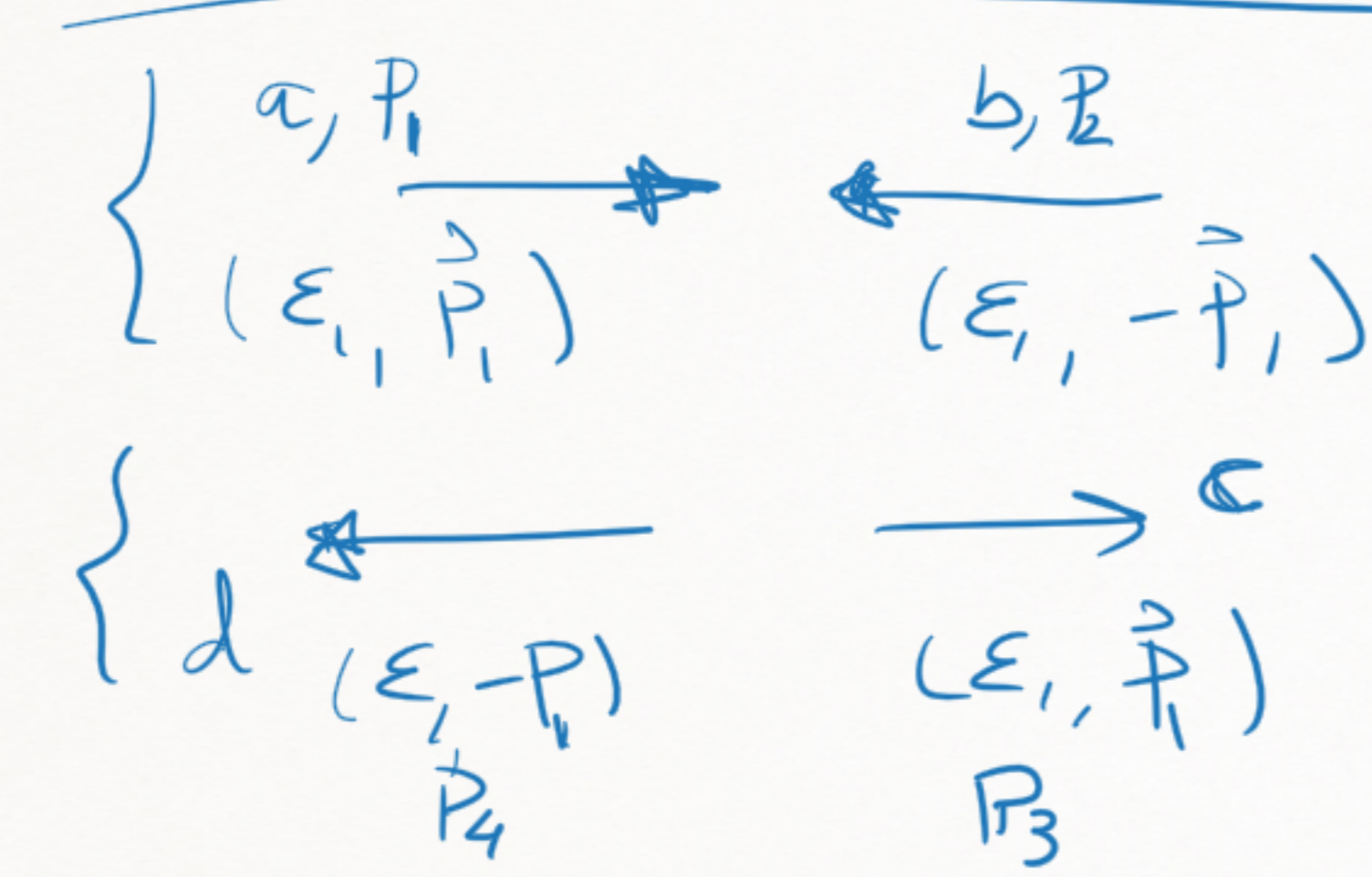
$$S = \frac{1}{2g^2} \int d^2x (\partial_\mu \hat{n})^2$$

$$\frac{\hat{n}^2 = 1}{N}$$

$O(N)$

N massive scalar

$$m = \mu e^{-\frac{2\pi}{g^2 N}}$$

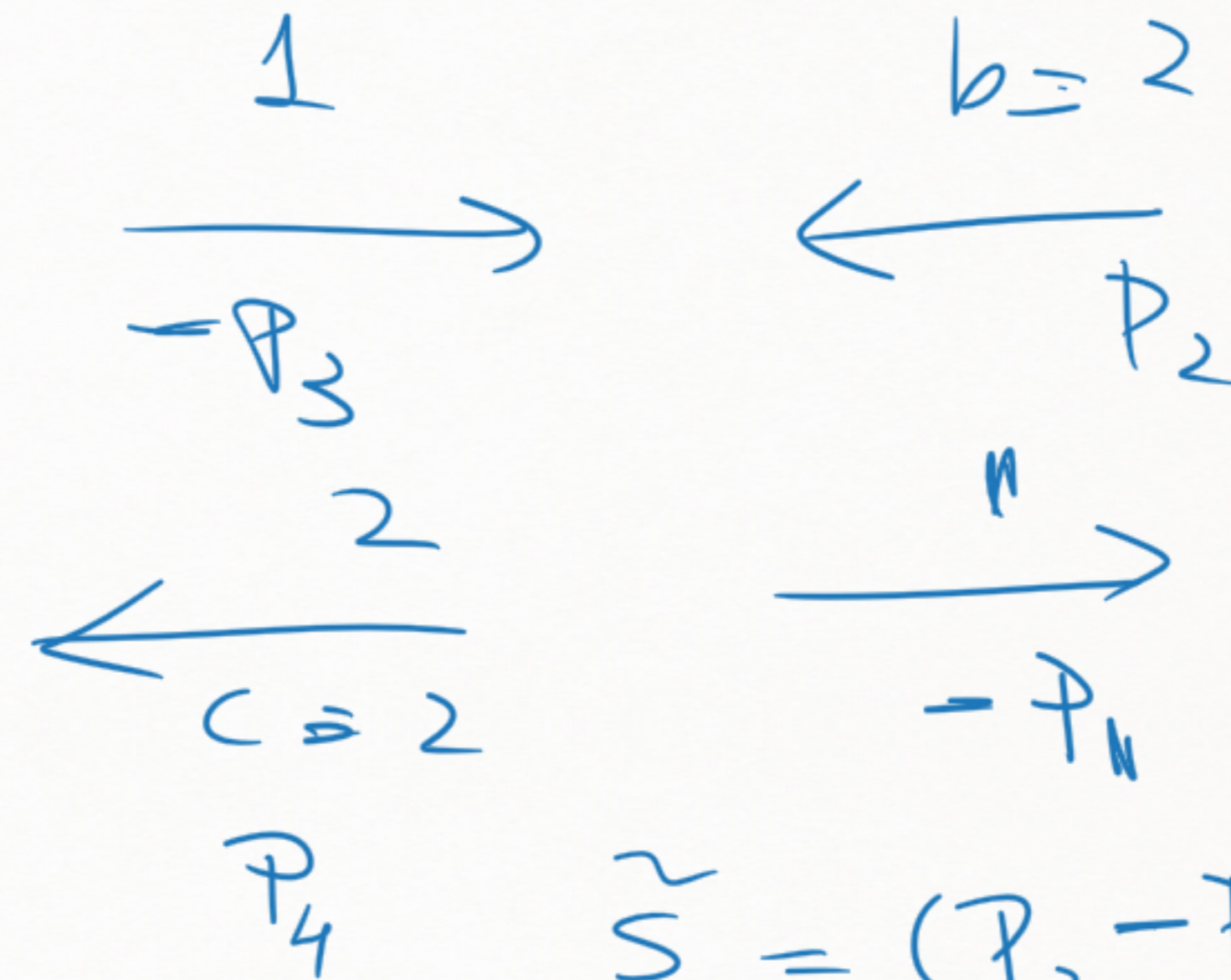
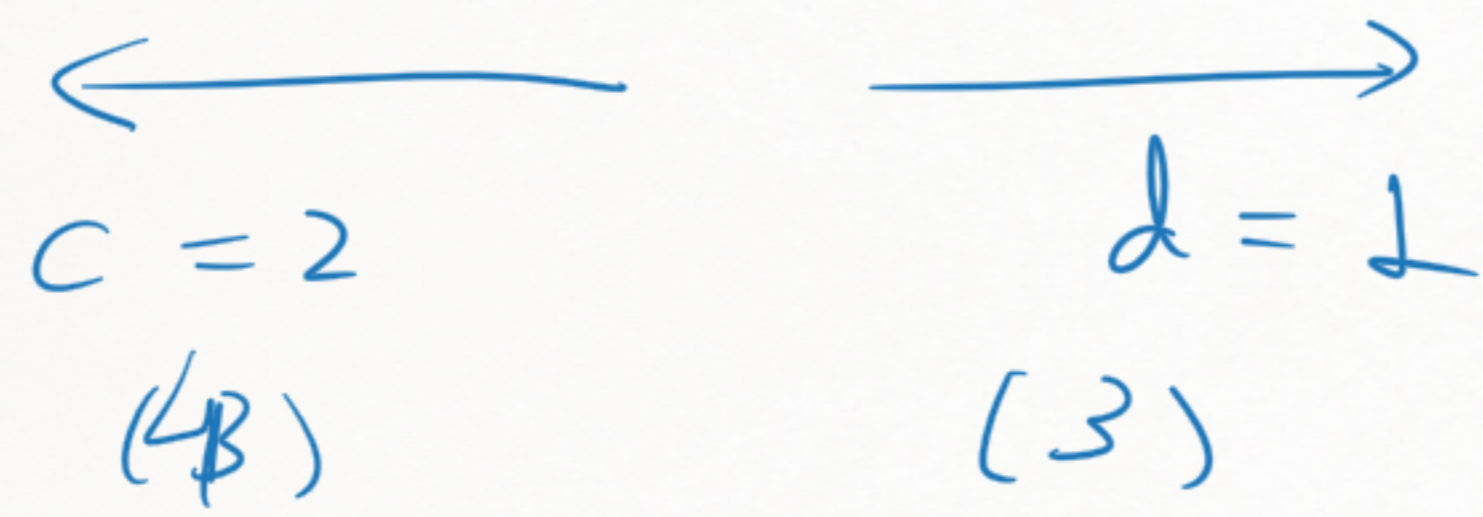
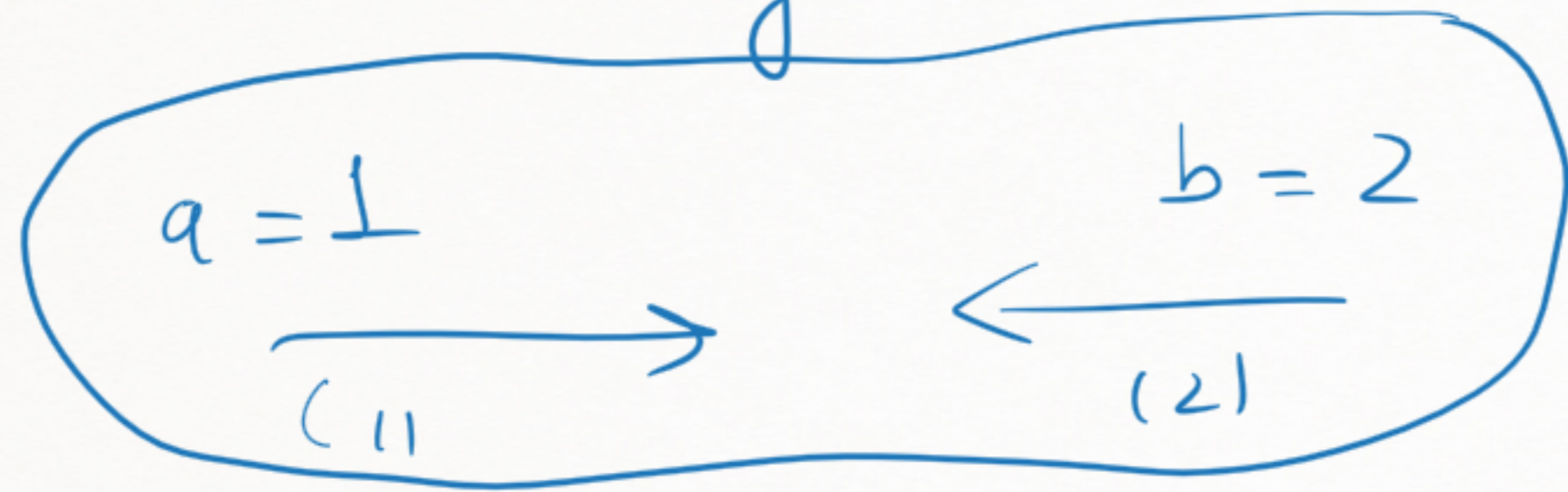


$$S_{ab \rightarrow cd} = S_T \delta_{ac} \delta_{bd} + S_R \delta_{bc} \delta_{ad}$$

$$+ S_A \delta_{ab} \delta_{cd} + (\text{crossed}) (\text{crossed})$$

$$s = (P_1 + P_2)^2 = 4\epsilon_1^2, \quad u = (P_1 - P_3)^2 = 0, \quad t = (P_2 - P_3)^2 = -4P_1^2 = 4m^2 - s$$

Crossing



$$\tilde{S} = (P_2 - P_3)^2 = 4m^2 - S$$

$$S_T(s) = S_T(\tilde{S} = 4m^2 - S)$$

$$S_A(s) = S_R(s)$$

$$ab \rightarrow \#, -, I$$

$$|I=0\rangle = \frac{1}{\sqrt{2}} (|11\rangle + |22\rangle + \dots - |NN\rangle)$$

$$\langle I=0 | S | I=0 \rangle = S_T + S_R + N S_A \approx S_I$$

$$\boxed{|S_{\pm}|^2 \leq 1}$$

$$S_{\pm} = S_T \pm S_R$$

$$\boxed{|S_{+}|^2 \leq 1}$$

$$\boxed{|S_{-}|^2 \leq 1}$$

$S_A(s), S_T(s), S_R(s)$ ←

$\sqrt{s} = k \epsilon^2$

t-cut
↓



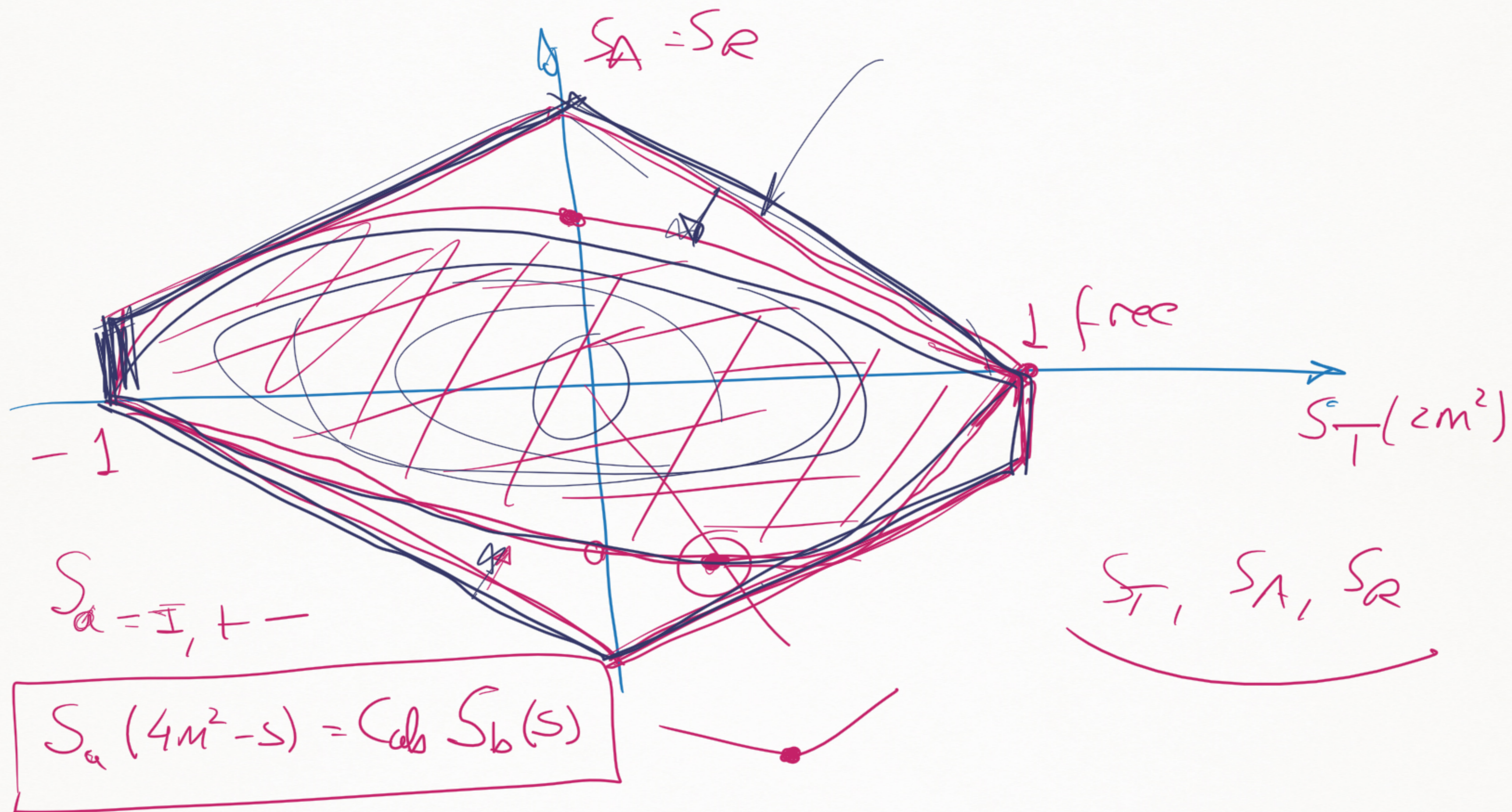
$|S_{\pm}|^2 \leq 1$
physical line.

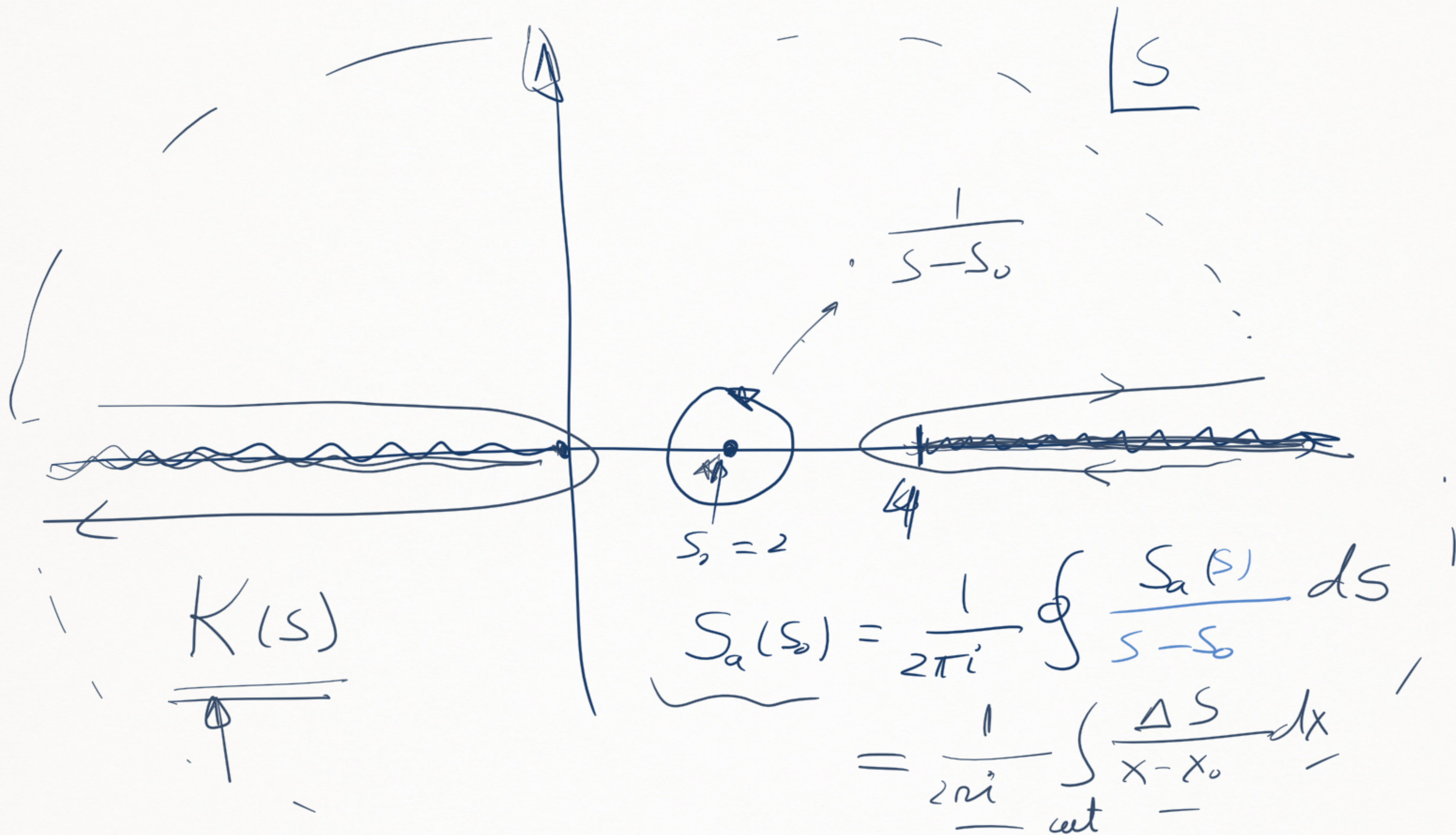
$k = \sqrt{\frac{s - 4m^2}{4}}$

$\left\{ \begin{aligned} S_T(4m^2 - s) &= S_T(s) \\ S_A(4m^2 - s) &= S_R(s) \end{aligned} \right.$

$\left\{ \begin{aligned} S_T(2m^2) \\ S_A(2m^2) = S_R(2m^2) \end{aligned} \right.$

$\begin{matrix} e^{-ikr} & + & e^{ikr} & e^{ikr} \\ e^{ikr} & + & e^{-ikr} & e^{-ikr} \end{matrix}$



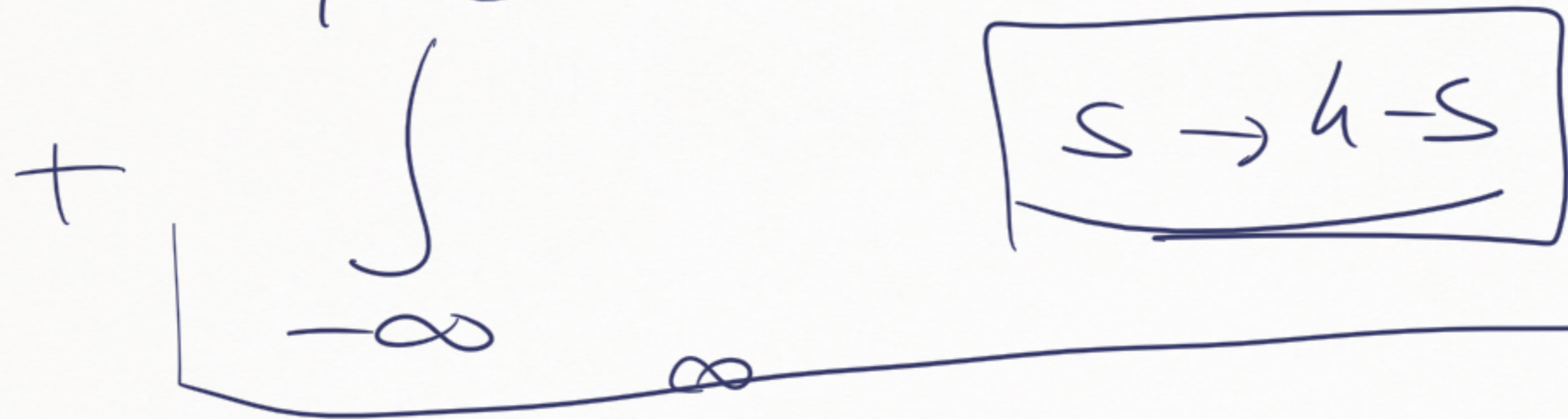


max F

$$F = \sum_{\infty} v_a S_a(s \neq 2) = \frac{1}{2\pi i} \oint K_a(s) S_a(s)$$

pole at $s=2$
residue v_a

$$= \frac{1}{2\pi i} \int_{\gamma} (K_a(s^+) S_a(s^+) - K_a(s^-) S_a(s^-))$$



$$K_a(s) = -K_b(h-s)$$

$$K_a(s) = (K_b(s^*))^*$$

$$F = \frac{2}{i\pi} \int_{\gamma} \text{Im}(K_a S_a)$$

$$F = \sum_{\infty} V_a S_a(s=2) = \frac{2}{\pi} \int_{\eta}^{\infty} \text{Im}(K_a S_a) dx \leq$$

$$\leq \frac{2}{\pi} \int_{\eta}^{\infty} |K_a| |S_a| dx \leq \frac{2}{\pi} \int_{\eta}^{\infty} |K_a| dx$$

\uparrow
 ≤ 1

$a = I, +, -$

$F \leq \min_{\{K_a\}} \frac{2}{\pi} \int_{\eta}^{\infty} |K_a| dx$

$$K_a = \frac{V_a}{s-2} \frac{2}{\sqrt{s} \sqrt{4-s}}$$

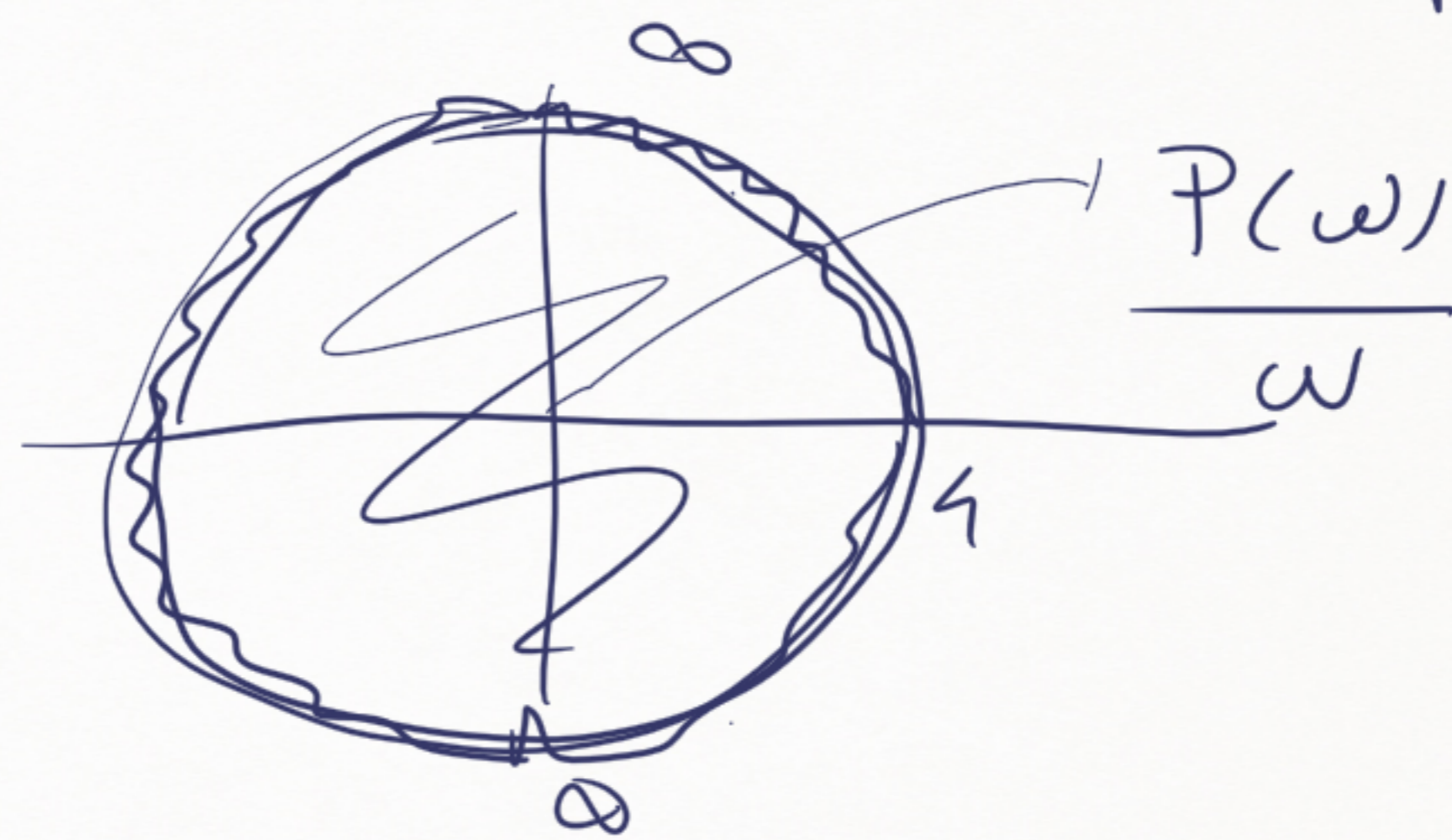
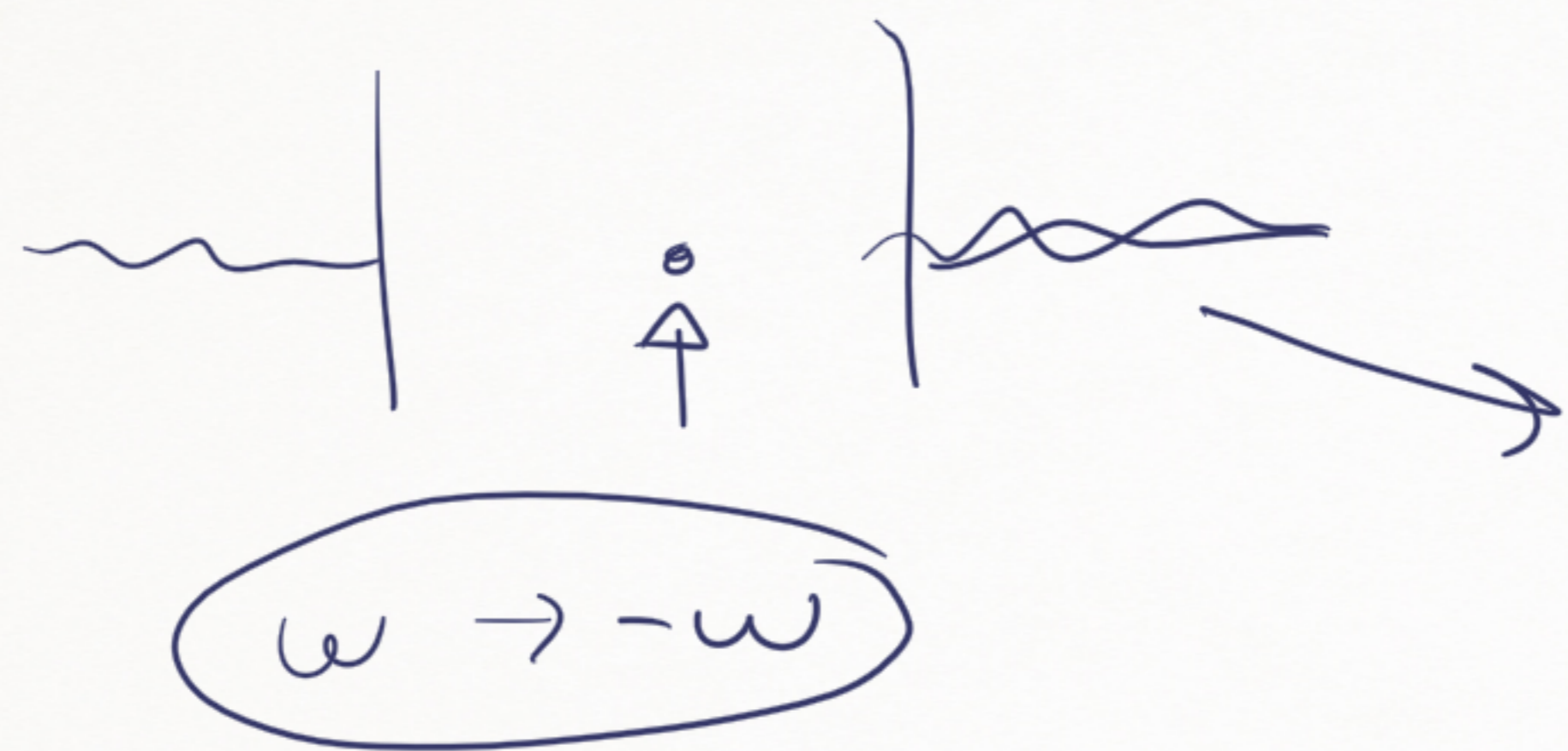
$F \leq \sum_a |V_a|$

$$S_a = \frac{i\bar{k}_a}{|k_a|}$$

$$\Downarrow = t - v$$

minimizer

$$w = \frac{\sqrt{s} - \sqrt{4-s}}{\sqrt{s} + \sqrt{4-s}}$$



- SNOWMASS paper (w/ Penedones, van Roes)

2203.02421

- Aninda Sinha (Bangalore)

-

Review

Positivity

$$D \begin{pmatrix} D & R \\ S_D & S_1 \\ S_2 & S_3 \end{pmatrix} \begin{pmatrix} S_D^+ & S_2^+ \\ S_1^+ & S_3^+ \end{pmatrix} = \begin{pmatrix} S_D S_D^+ + S_1 S_1^+ & \\ & \\ & \\ & \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$S_D S_D^+ + S_1 S_1^+ = \mathbb{1} \quad \begin{pmatrix} \mathbb{1} & S_D \\ S_1^+ & \mathbb{1} \end{pmatrix} \succeq 0$$

$$S S^+ = \mathbb{1}$$

$$\underbrace{S_D S_D^+ - \mathbb{1}} = \underbrace{-S_1 S_1^+}_{\leq 0}$$

$$A_1 \geq 0 \quad A_2 \geq 0$$

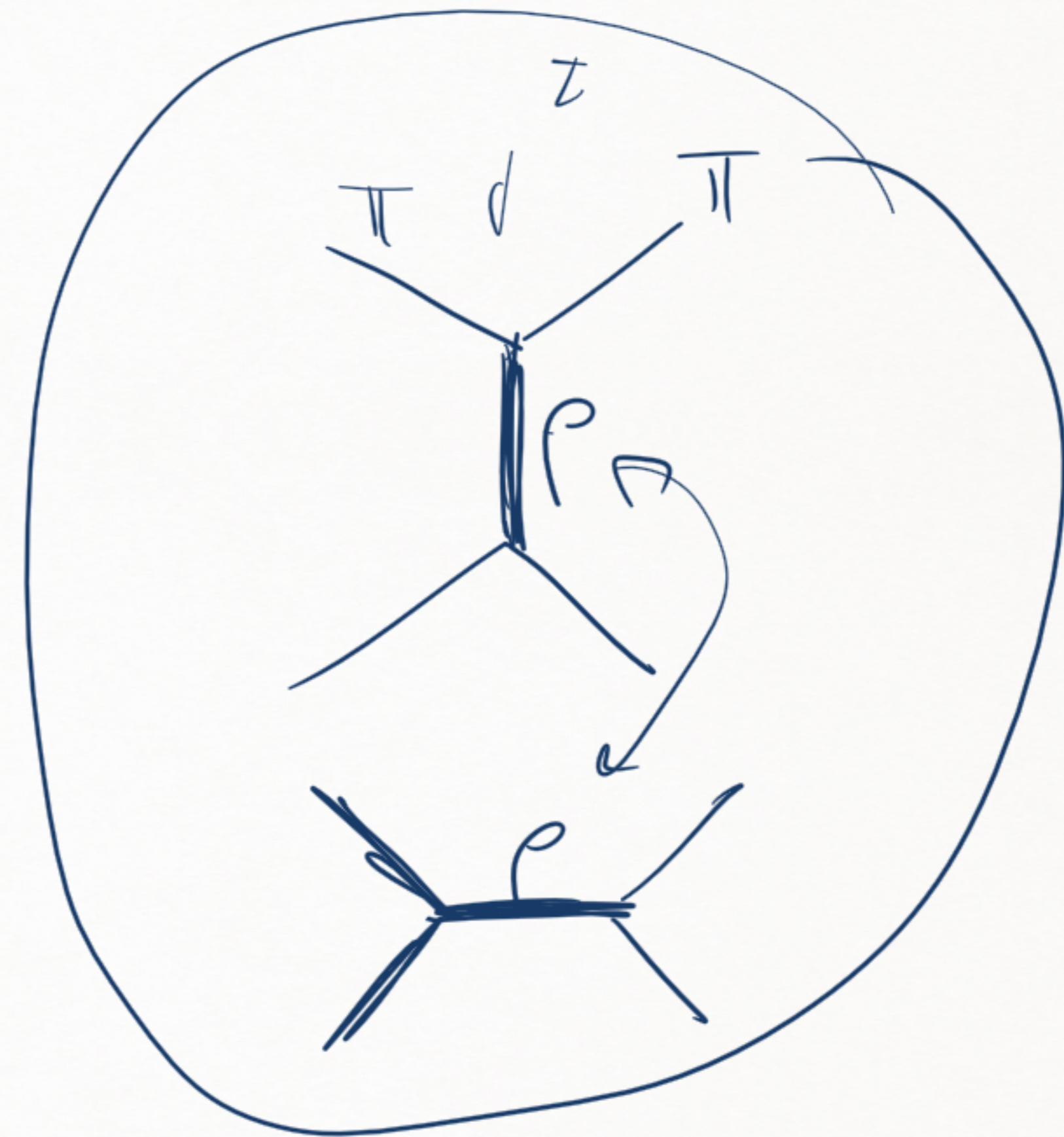
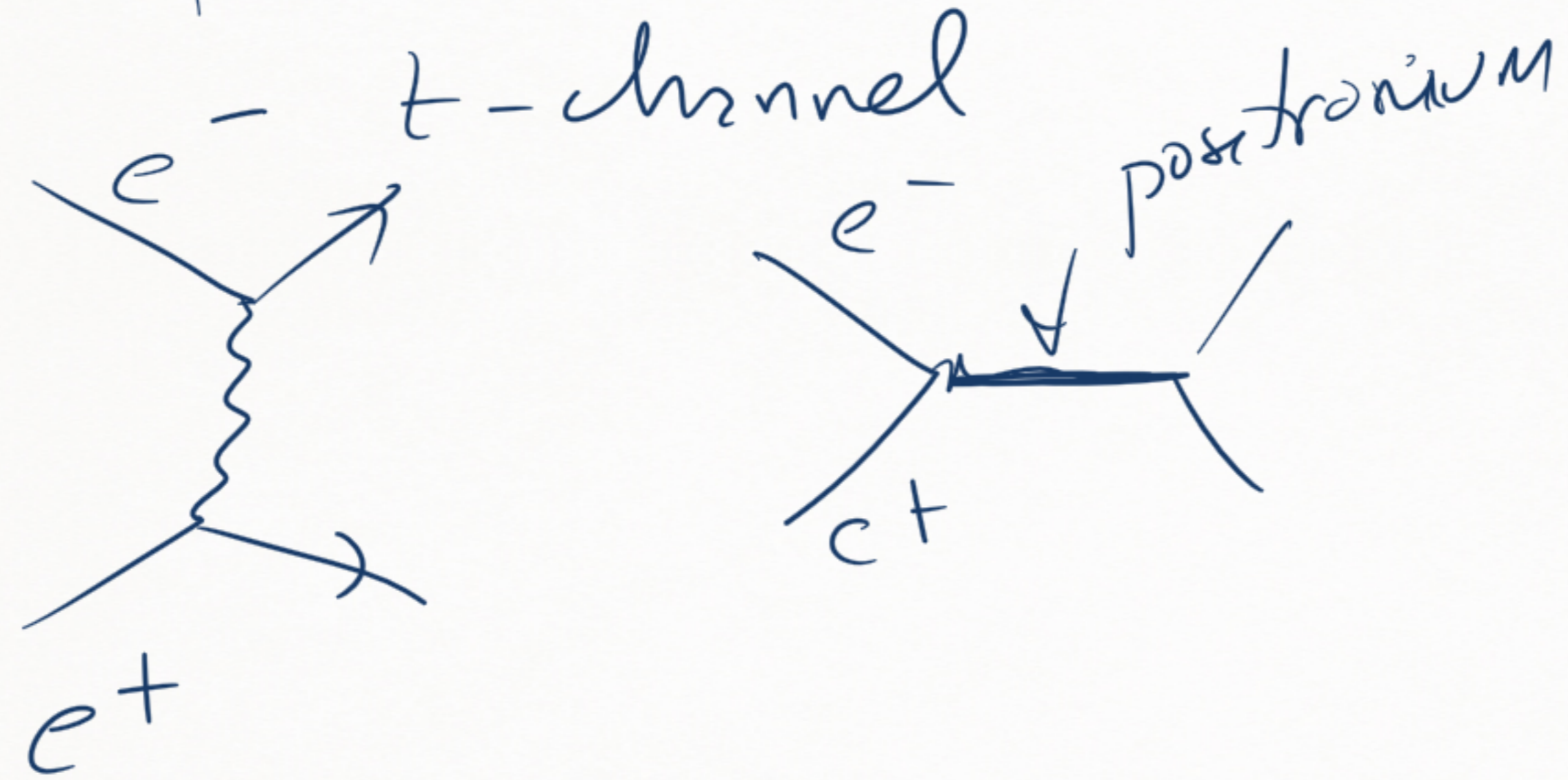
$$0 \leq \alpha \leq 1 \quad \alpha A_1 + (1-\alpha) A_2 \geq 0$$

$$\alpha \underbrace{\langle v | A_1 | v \rangle}_{\geq 0} + (1-\alpha) \underbrace{\langle v | A_2 | v \rangle}_{\geq 0} \geq 0$$

Space of positive semi-definite matrices is convex

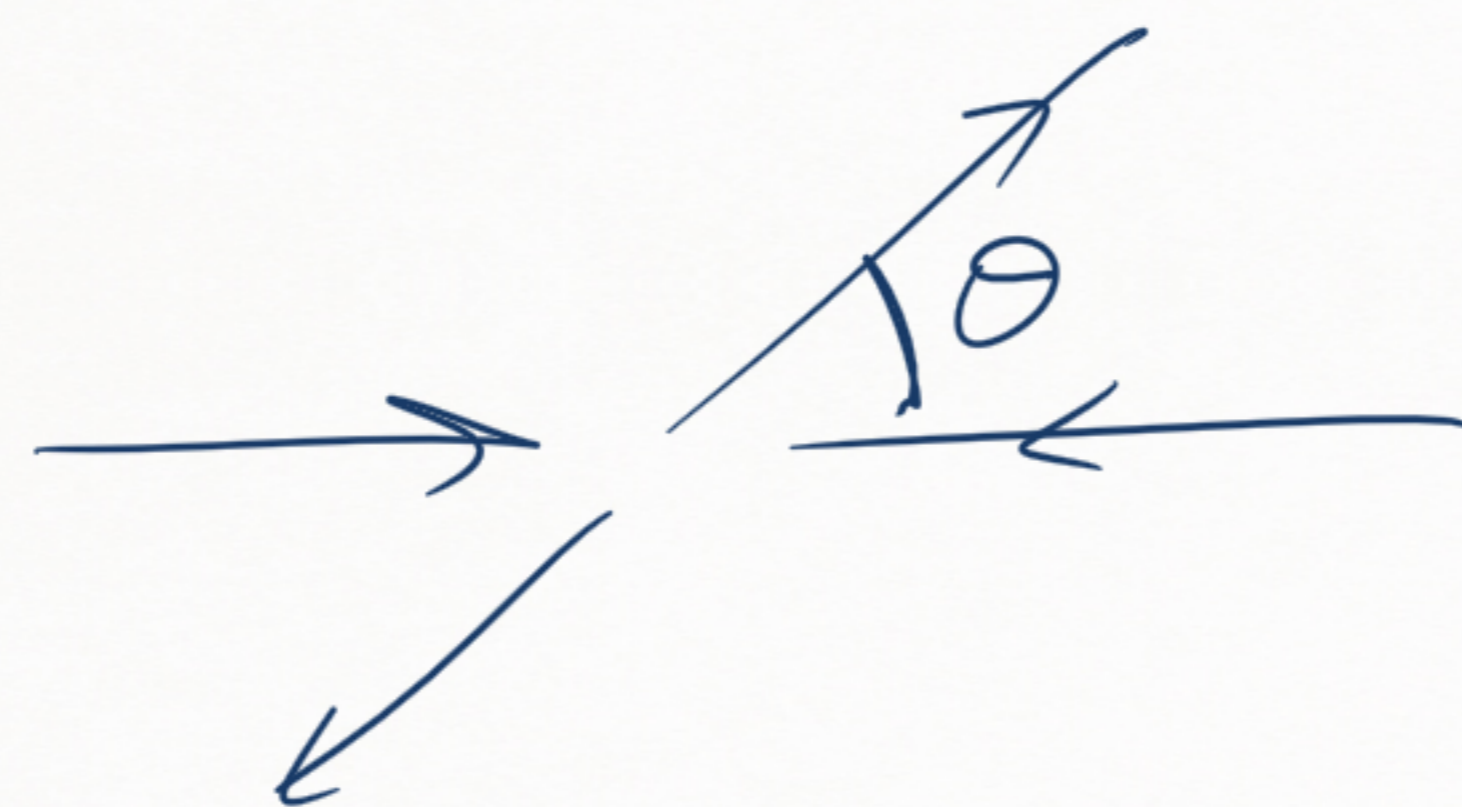
$SU(3), 2 \rightarrow$ theory of pions

$SU(N), N_f \rightarrow$



$$\begin{array}{ccc}
 \pi^0 + \pi^0 & \longrightarrow & \pi^+ + \pi^- \\
 P_1 \quad P_2 & & P_3 \quad P_4 \\
 m^2 = 1 & &
 \end{array}$$

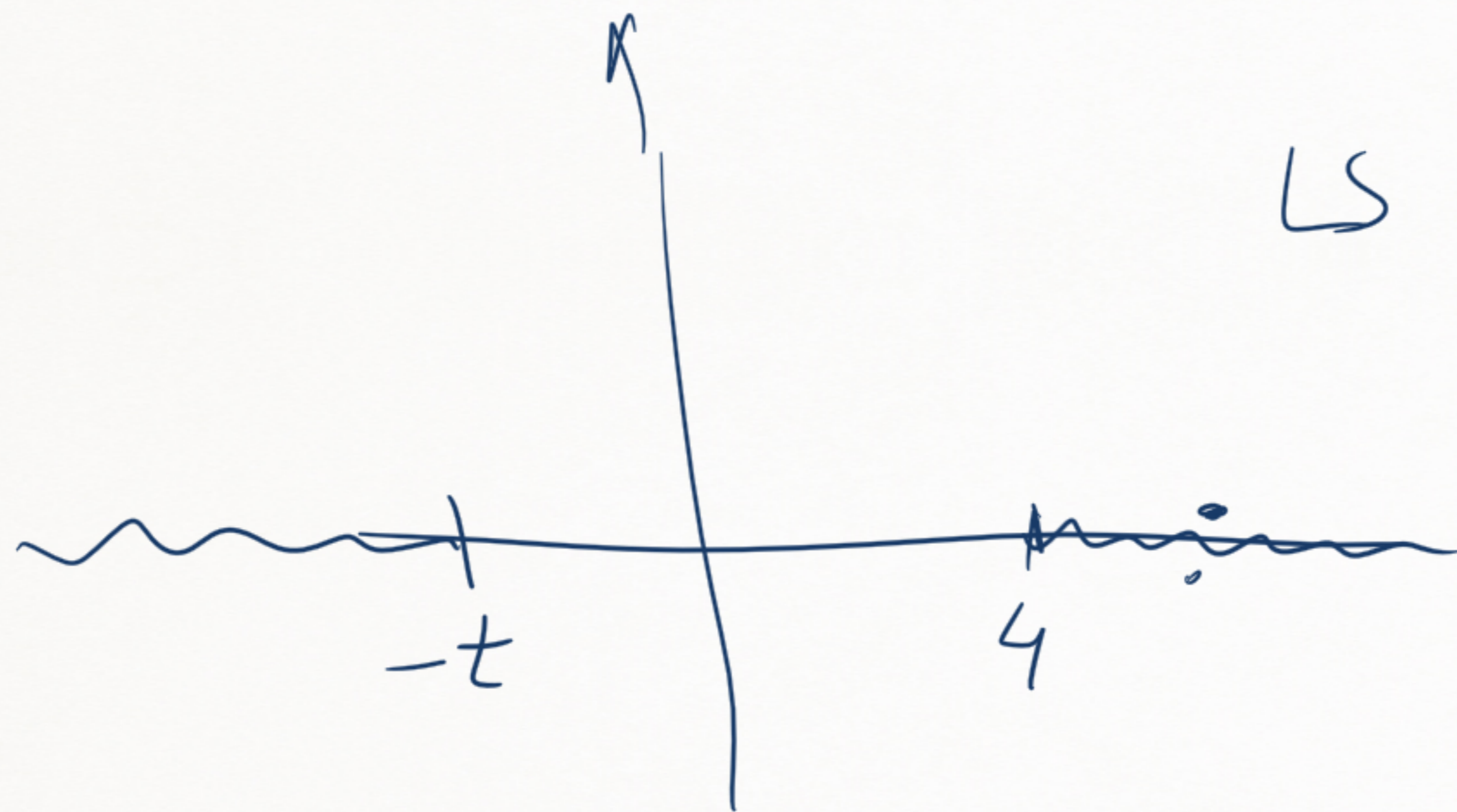
$$\begin{aligned}
 s &= (P_1 + P_2)^2 \\
 t &= (P_1 - P_3)^2 \\
 u &= (P_1 - P_4)^2
 \end{aligned}
 \quad s + t + u = 4$$



$$S_e(s) = \underbrace{1}_1 e^{i\delta_e}$$

$$\begin{aligned}
 &\Lambda_{\text{QCD}} \\
 &m/\Lambda_{\text{QCD}} \quad \underline{N} \\
 S_e &= 1 + i \underline{h_e(s)} \quad |S_e|^2 \leq 1 \\
 &\quad \boxed{|h_e(s)|^2 \leq 2 \Im h_e(s)}
 \end{aligned}$$

$$\left[F(s^+, t) = \sum_{l \text{ even}} \frac{2}{\pi} \sqrt{\frac{s}{s-4}} (2l+1) \operatorname{he}(s)_x \right. \\
 \left. \times \underset{\cos \theta}{\mathcal{P}}_l \left(1 + \frac{2t}{s-4} \right) \right]$$



$$s \leftrightarrow t \leftrightarrow u$$

$$\left\{ \begin{array}{l} s > 4 \\ t > 4 \\ u > 4 \end{array} \right.$$

$$s + t + u = 4$$

$$s + t < 0$$

$$\begin{aligned}
 \mathbb{F}(s, t, u) &= f_0 + \frac{1}{\pi} \int_{\mathfrak{h}}^{\infty} dx \, \sigma(x) \left(\frac{1}{x-s} + \frac{1}{x-t} + \frac{1}{x-u} \right) \\
 &\quad + \frac{1}{\pi^2} \int_{\mathfrak{h}}^{\infty} dx \int_{\mathfrak{h}}^{\infty} dy \, \rho(x, y) \left(\frac{1}{(x-s)(y-t)} + \frac{1}{(x-s)(y-u)} + \frac{1}{(x-t)(y-u)} \right)
 \end{aligned}$$

$\frac{1}{x-s-i\epsilon} = \text{p.o.} \frac{1}{x-s} + i\pi \delta(x-s)$

Study the space of such functions

→ unitarity saturated \Rightarrow ρ vanishes in certain region

o) " " at all energies \Rightarrow free theory
Aks theorem

$$\Rightarrow \sum_{n=0}^{\infty} t^n P_n(z) = \frac{1}{\sqrt{1-2zt+t^2}} \quad |z| \leq \frac{1+t^2}{2t}$$

$$\sum_{l=0}^{\infty} a_l P_l(\chi\mu) < \infty \quad \lim_{l \rightarrow \infty} |a_l| = e^{-\alpha} \quad |\mu| < \alpha$$

$$F = \frac{2}{\pi} \sqrt{\frac{s}{s-4}} \sum_{l \text{ even}} (2l+1) h_l(s) P_l(\text{ch}\mu)$$

$\hookrightarrow 1 + \frac{2t}{s-4}$

$$t=4 \quad \text{ch}\mu = \text{ch}\alpha_1(s) = \frac{s+4}{s-4}$$

$$\lim_{l \rightarrow \infty} |h_l(s)|^{1/l} \leq e^{-\alpha_1(s)}$$

saturate

$$|h_l(s)| \sim \left(\frac{\sqrt{s-2}}{\sqrt{s+2}} \right)^l$$

$$\Im h_l(s) = \frac{1}{2} \underbrace{|h_l(s)|^2}$$

$$\sum_{l=0}^{\infty} (2l+1) \underbrace{\text{Im } h_l(s)}_{|u| < 2\alpha_1} P_l(\text{ch } \mu) = \sigma(s) + \int_4^{\infty} dy \rho(s, y) \left(\frac{1}{y-t} + \frac{1}{y-u} \right)$$

$\boxed{4 < s < 16}$

$$1 + \frac{2t_2(s)}{s-4} = \text{ch}(2\alpha_1(s))$$

$$\boxed{t < t_2(s)}$$

$$\underline{4 < s < 16}$$

$$\rho(s, y) = 0$$

s saturating unitarity

$$\boxed{y < t_2(s)} = 16 + \frac{64}{s-4}$$



Suppose no particle production

$$p(s, y) = 0$$

$$y > 4$$

$$4 < s < 16$$

$$p(s, y) = 0$$

$$\underline{s > 4} \quad \boxed{4 < y < 16}$$

$$1 + \frac{2t}{s-4} = c\theta$$

$$\underline{\theta=0} \rightarrow t=0$$

$$F(s, 0, 4-s) = F(s) \quad \text{forward amplitude}$$

$$F(s) = F(4-s)$$

