Dissipation induced by local non-Markovian baths

Laura Foini

CNRS, IPhT Saclay

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In collaboration with









Saptarshi Majumdar

Oscar Bouverot Dupuis

Alberto Rosso

Thierry Giamarchi

Where do we stand

Time scales of the ``environment"

Fast	Slow	Quenched disorder
Lindblad approach for Markovian evolution	Caldeira and Leggett type of bath Power law correlated	Interaction with static impurities
High T regime	Low T regime	

The initial motivation

Two types of localisations

- with quenched disorder (Anderson insulator or Bose glass/MBL phase with interactions)
- spin-boson model: single particle with slow bath

We want to study

Many-body system with slow bath \equiv annealed disorder

Can we talk about "localisation"? (at T = 0)

The system

Spin chain (S = 1/2):

$$H = J_{xy} \sum_{i} \left[S_{i}^{x} S_{i+1}^{x} + S_{i}^{y} S_{i+1}^{y} \right] + J_{z} \sum_{i} S_{i}^{z} S_{i+1}^{z}$$

Equivalent to a model of interacting fermions:

$$H = -J_{xy} \sum_{i} \left[c_{i}^{\dagger} c_{i+1} + c_{i+1}^{\dagger} c_{i} \right] + J_{z} \sum_{i} (n_{i} - 1/2)(n_{i+1} - 1/2)$$

Two cases:

- ► Half filling $\langle n_i \rangle = \pi q_F = \frac{1}{2} (\langle S^z \rangle = 0)$ (commensurate case)
- ► Doped system $1/2 < \langle n_i \rangle = \pi q_F < 1 \ (0 < \langle S^z \rangle < 1/2)$ (incommensurate case)

The system

Without the bath the system is in a *Luttinger liquid* phase (LL)

$$S_{LL} = \frac{1}{2\pi K} \int dx \int d\tau \left[u \left(\partial_x \phi \right)^2 + \frac{1}{u} \left(\partial_\tau \phi \right)^2 \right]$$
$$S^z \simeq -\frac{1}{\pi} \nabla \phi + \cos(\phi - 2q_F x)$$

- Critical phase with power law correlations
- ▶ Finite compressibility (susceptibility) and spin stiffness
- Perfectly conducting phase

Coupling to *local* baths

Caldeira Leggett bath

$$H_{B} = \sum_{i} \sum_{k} \left[\frac{1}{2} P_{i,k}^{2} + \frac{1}{2} \Omega_{k} X_{i,k}^{2} \right]$$



System-bath interaction

Coupling via the *density*

$$H_{SB} = \sum_{i} S_{i}^{z} \sum_{k} \lambda_{k} X_{i,k} = \sum_{i} S_{i}^{z} h_{i}(t)$$

In the limit of static bath

$$H_{tot} = J_{xy} \sum_{i} \left[S_{i}^{x} S_{i+1}^{x} + S_{i}^{y} S_{i+1}^{y} \right] + J_{z} \sum_{i} S_{i}^{z} S_{i+1}^{z} + \sum_{i} S_{i}^{z} h_{i}$$

paradigmatic model to study localisation with interaction

Nature of the bath

Spectral function

$$\operatorname{Im} \int e^{i\omega t} \langle [h_i(t), h_j(0)] \rangle = \delta_{ij} J(\omega)$$

$$J(\omega) = \pi \sum_{j} \frac{\lambda_k^2}{\Omega_k} \delta(\omega - \Omega_k) = \alpha \, \omega^s \quad \text{for } \omega < \Omega_D$$

▶ s < 1 sub Ohmic

▶ s = 1 Ohmic

▶ s > 1 super Ohmic

Caldeira, Leggett, Phys. Rev. Lett. 46, 211 (1981) Leggett, Chakravarty, Dorsey, Fisher, Garg, Zwerger, Rev. Mod. Phys. (1987)

Total action

... integrating out the bath (annealed average)

$$S_{tot} = S_{LL} + S_{Dis}$$

$$S_{Dis} = -\alpha \int \mathrm{d}x \int \mathrm{d}\tau \int \mathrm{d}\tau' \frac{\cos(\phi(x,\tau) - \phi(x,\tau'))}{|\tau - \tau'|^{1+s}}$$

Bath with long-range correlations in time and uncorrelated in space 2D classical field theory studied with

- Gaussian variational method
- ► RG
- Numerical simulations

Quantities of interest

Propagator
$$G(q, \omega_n) = \langle \phi(q, \omega_n) \phi(-q, -\omega_n) \rangle$$
 $\omega_n = \frac{2\pi n}{\beta}$

RG

$$\partial_l K(l) = -2K^2(l)\alpha(l)$$

$$\partial_l \alpha(l) = (2 - s - 2K(l))\alpha(l)$$

Critical point
$$K_c = 1 - \frac{s}{2} \quad \alpha = 0$$

BKT transition



Cazalilla et al., Phys. Rev. Lett. 97, 076401 (2006)

Gaussian variational method

Luttinger liquid phase

$$G_{LL}^{-1} \simeq \frac{1}{2\pi} \left[\frac{u_r}{K_r} q^2 + \frac{1}{u_r K_r} \omega_n^2 \right]$$

Dissipative phase

$$G_{Dis}^{-1} \simeq \frac{1}{2\pi} \left[\frac{u_r}{K_r} q^2 + \eta \omega_n^s \right]$$

Order parameter

$$\cos 2\phi \simeq \begin{cases} \frac{1}{L^{K_r}} & \text{LL phase} \\ const - \frac{1}{L^{1-s/2}} & \text{Dissipative phase} \end{cases}$$



Statistical tilt symmetry

$$S_{Dis}[\phi + hx] = S_{Dis}[\phi]$$

Constant compressibility $\chi = u/K$



Small ω_n behaviour

$$C(\omega_n) = \sum_{q} G(q, \omega_n) \simeq \begin{cases} \omega_n^{-1} & \text{LL phase} \\ \omega_n^{-s/2} & \text{Dissipative phase} \end{cases}$$



Dissipative phase

Gapless spin density wave $S^{z}(x) \simeq \cos(2q_{F}x) \langle \cos 2\phi \rangle$



Transport properties

DC Conductivity

$$\sigma_{DC} = \lim_{\epsilon \to 0} \left(\frac{e^2}{\pi^2 \hbar} \epsilon^{1-s} \right) = \begin{cases} \infty & \text{Super Ohmic} \\ Const & \text{Ohmic} \\ 0 & \text{Sub Ohmic} \end{cases}$$

Bath induced localisation !

Majumdar et al., Phys Rev, B 107, 165113 (2023) Majumdar et al., Phys. Rev. B 108, 205138 (2023)

Commensurate case

Dissipative phase

$$G_{Dis}^{-1} \simeq \frac{1}{2\pi} \left[\frac{u}{K} q^2 + \frac{1}{uK} \omega_n^2 + \Delta^2 \right]$$

Gapped antiferromagnetic phase

Antiferromagnetic order enhanced by the bath !

Malatsetxebarria et al., Phys. Rev. A 88, 063630 (2013) Bouverot-Dupuis et al., Phys. Rev. B 109, 205148 (2024)

Phase diagram



Order induced by on site dissipation

Other examples

- Werner, Troyer, Sachdev, J. Phys. Soc. Jpn (2005)
 - Cai, Schollwöck, Pollet, PRL (2014)
 - Weber, Luitz, Assad, PRL (2022)
- Kuklov, Prokof'ev, Radzihovsky, Svistunov, PRL, PRB, PRA (2023-2024)
 - Ribeiro, McClarty, Ribeiro, Weber, PRB (2024)

Where we would like to go

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Conclusions

- Order induced or enhanced by the bath
- The bath can induce a gapless insulating phase
- Whole phase diagram as a function of the magnetic field? (commensurate/incommensurate phase transition)
- Entanglement entropy between the system and the bath?
- Can we draw some link with other open quantum systems?

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Thank you!