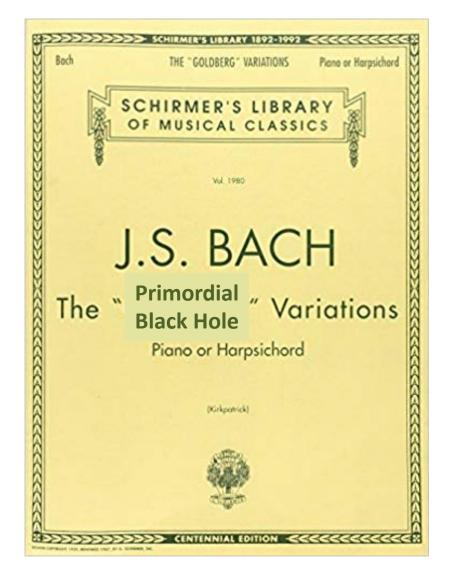
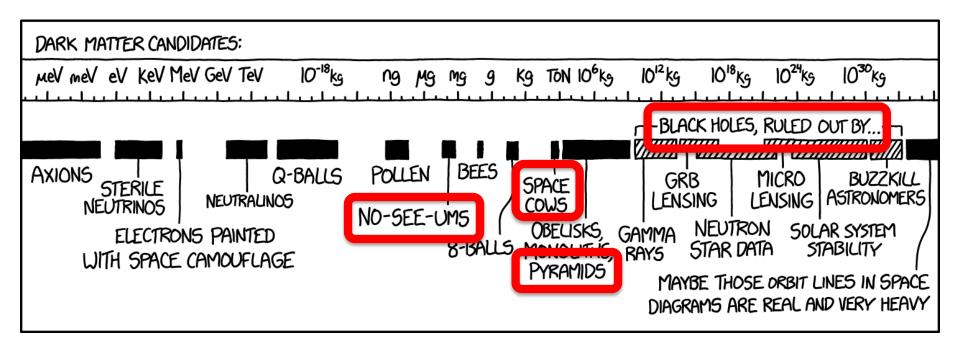


Stefano Profumo

University of California, Santa Cruz





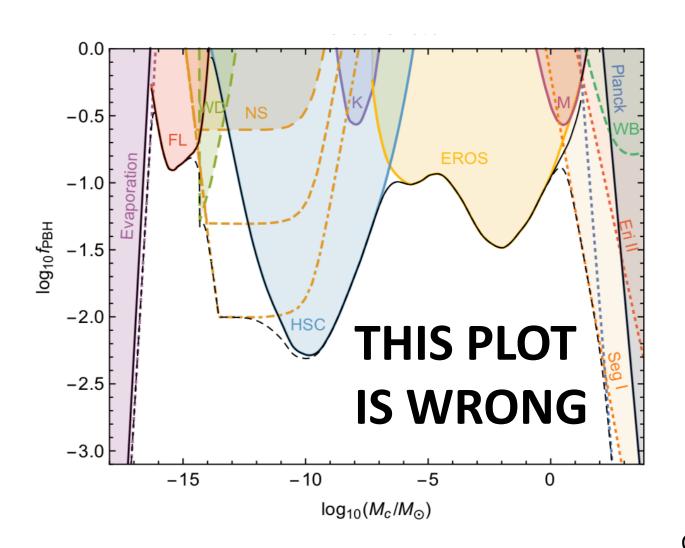


- Why it is interesting to consider PBH as Dark Matter
- > Where it is interesting to look for PBH as Dark Matter
- ➤ ...some "NO-SEE-UMS", "SPACE COWS", "PYRAMIDS"

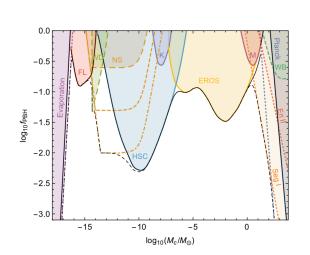
source: xkcd

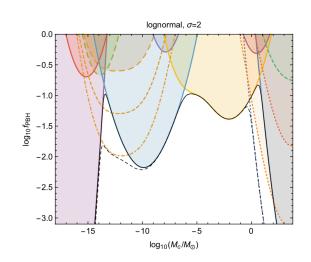
Can there be enough PBH around to be the DM?

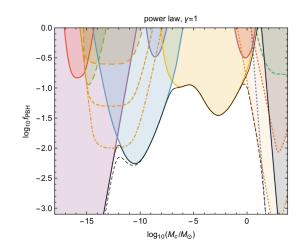
What is the maximal fraction of dark matter in PBH?



The fraction of PBH that could be the dark matter depends on the mass function!







...what is the mathematical function that maximizes the mass fraction of primordial black holes compatibly with constraints?

The Maximal-Density Mass Function for Primordial Black Hole Dark Matter



Benjamin V. Lehmann, Stefano Profumo and Jackson Yant

Department of Physics, University of California Santa Cruz, 1156 High St., Santa Cruz, CA 95064, USA Santa Cruz Institute for Particle Physics, 1156 High St., Santa Cruz, CA 95064, USA

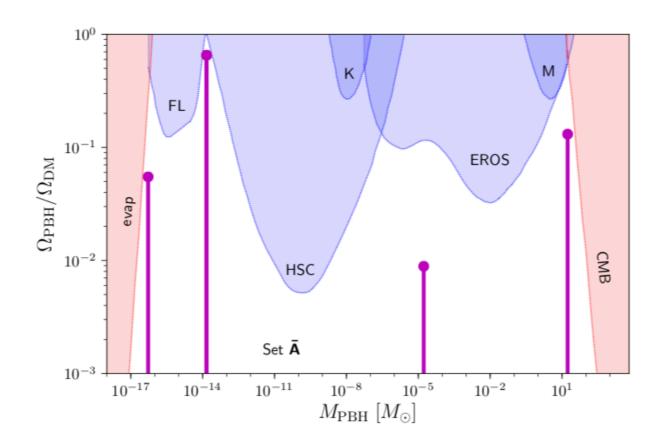
E-mail: blehmann@ucsc.edu, profumo@ucsc.edu, jyant@ucsc.edu

Abstract. The advent of gravitational wave astronomy has rekindled interest in primordial black holes (PBH) as a dark matter candidate. As there are many different observational probes of the PBH density across different masses, constraints on PBH models are dependent on the functional form of the PBH mass function. This complicates granted statements about

Answer: with N independent constraints, the optimal function is a linear combination of N delta functions with calculable relative weights

 $\min \{ \|\mathbf{x}\| \mid \mathbf{x} \in \operatorname{conv} \{ \mathbf{g}(M) \mid M \in U \} \}$

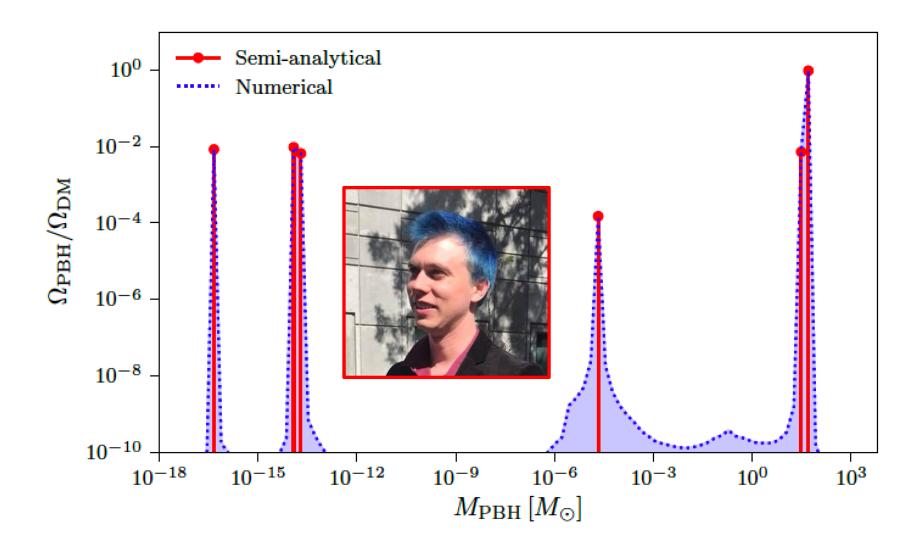
* Lehmann, Profumo and Yant, JCAP 2018



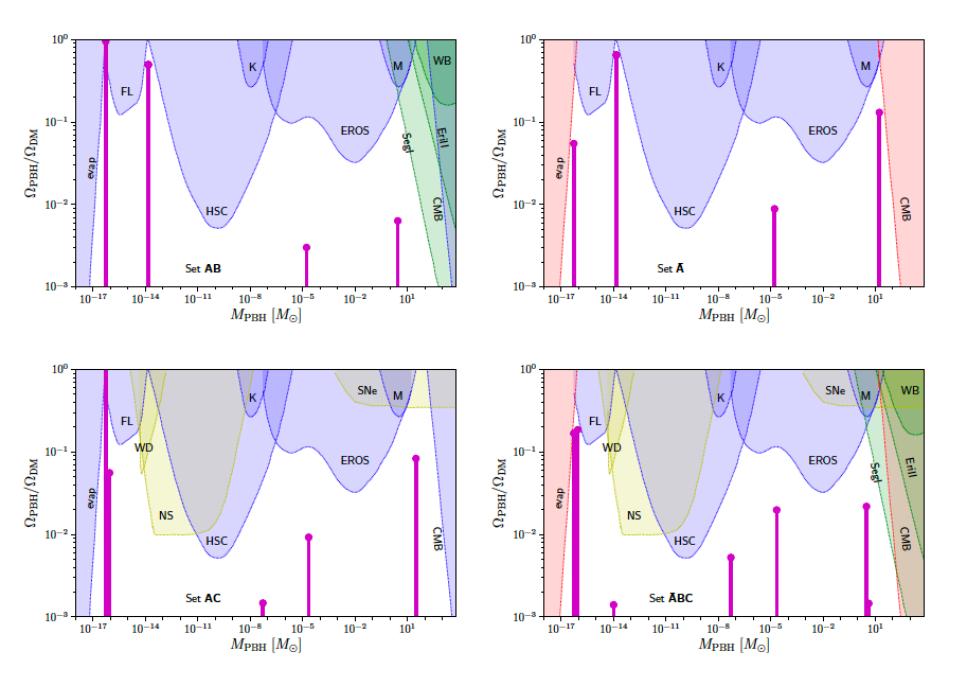
Answer: with N independent constraints, the optimal function is a linear combination of N delta functions with calculable relative weights

^{*} Lehmann, Profumo and Yant, JCAP 2018

Numerical validation



^{*} Lehmann, Profumo and Yant, JCAP 2018



* Lehmann, Profumo and Yant, JCAP 2018

	$f_{ m mono}$	$f_{ m max,all}$	$f_{ m max,GW}$	$\sigma[\psi]/M_{\odot}$	$\langle M/M_{\odot} \rangle$
Α	27.17	27.25	2.580	2.259	31.09
AB	1.372	1.965	5.139	0.162	0.009
AC	1.371	1.443	0.566	7.294	1.807
ABC	1.371	1.402	2.936	0.220	0.015
Ā	0.991	1.502	2.171	4.827	1.492
ĀΒ	0.991	1.437	11.07	0.221	0.017
ĀC	0.330	0.484	0.364	7.963	5.430
ĀBC	0.330	0.405	0.982	0.741	0.182

So YES, depending on the constraints choice, PBH can be 100% of the dark matter!

^{*} Lehmann, Profumo and Yant, JCAP 2018

Is there a goldilocks signature of PBH?

Yes! BH merger with a sub-Chandrasekhar mass (1.4 M_{sun})

LIGO search results are out*

Is there a goldilocks signature of PBH?

Yes! BH merger with a sub-Chandrasekhar mass (1.4 M_{sun})

Given a mass function, one can calculate:

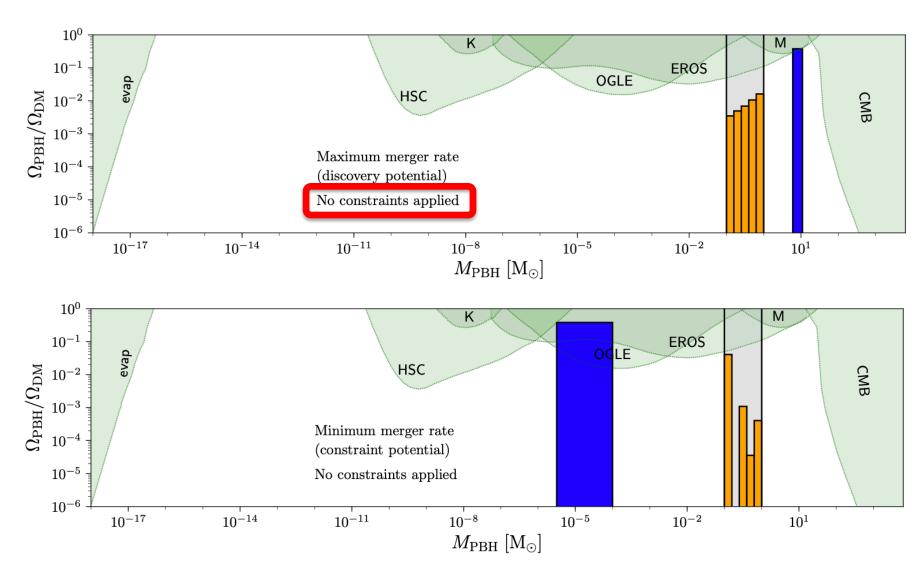
1. Rate of "goldilocks events"

$$R_{\rm DP}(\psi) = \int_{\rm DP^2} dm_1 dm_2 \, \mathcal{R}(m_1, m_2) V_{\rm eff}(m_1, m_2),$$

2. Mass fraction of light+detectable BHs

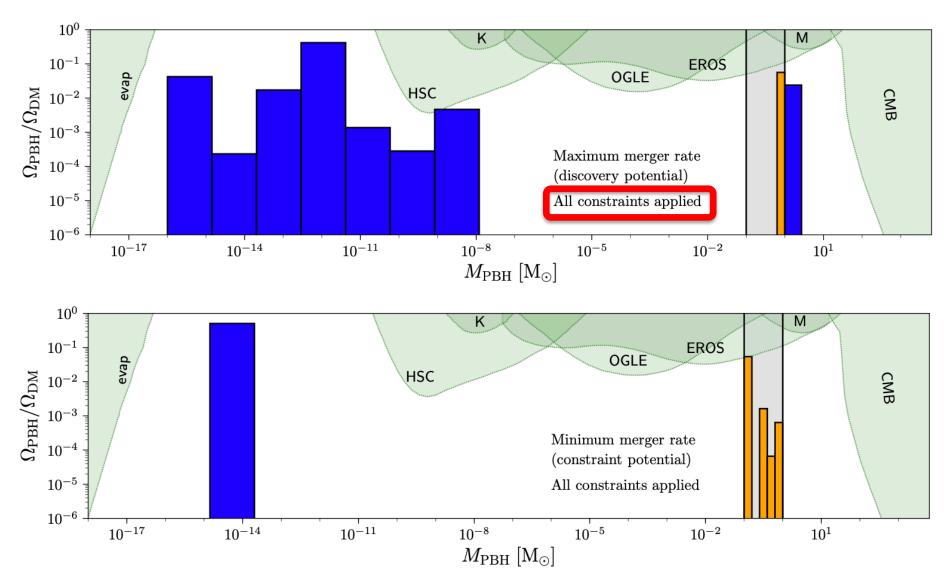
$$r_{\rm DP} = \frac{1}{f_{\rm PBH}} \int_{m_{\rm DP}^{\rm min}}^{m_{\rm DP}^{\rm max}} {\rm d}m \, \psi(m).$$

We can numerically compute the maximal and minimal possible "goldilocks event rate"



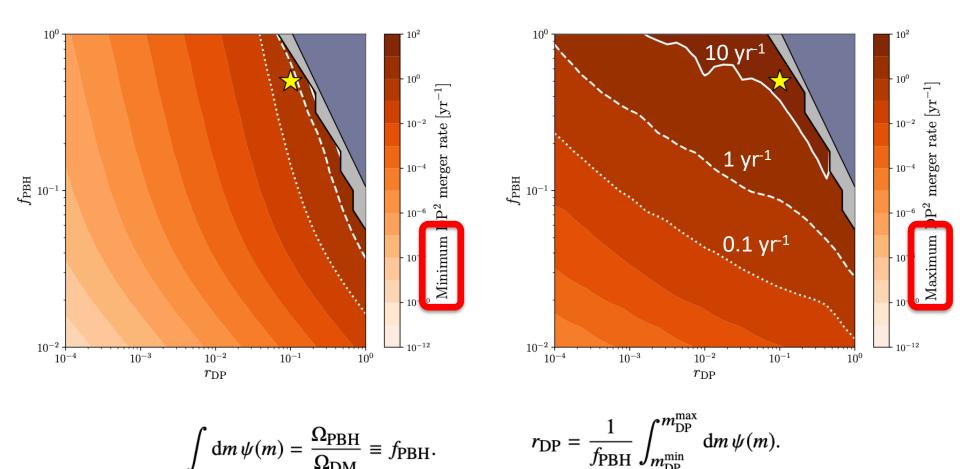
^{*} Lehmann, Profumo and Yant, MNRAS

We can numerically compute the maximal and minimal possible "goldilocks event rate"



^{*} Lehmann, Profumo and Yant, MNRAS

We can numerically compute the maximal and minimal possible "goldilocks event rate"



Besides the mass, LIGO informs us about the spin of BHs...

Besides the mass, LIGO informs us about the spin of BHs...

LIGO/Virgo Collaboration arXiv:1811.12940

Event	$m_1/{ m M}_{\odot}$	$m_2/{ m M}_{\odot}$	$\mathcal{M}/\mathrm{M}_{\odot}$	$\chi_{ m eff}$	$M_{\rm f}/{ m M}_{\odot}$	$a_{ m f}$	$E_{\rm rad}/({\rm M}_{\odot}c^2)$	$\ell_{\rm peak}/({\rm erg~s^{-1}})$	$d_L/{ m Mpc}$	z	$\Delta\Omega/deg^2$
GW150914	$35.6^{+4.8}_{-3.0}$	$30.6^{+3.0}_{-4.4}$	$28.6^{+1.6}_{-1.5}$	$-0.01^{+0.12}_{-0.13}$	$63.1^{+3.3}_{-3.0}$	$0.69^{+0.05}_{-0.04}$	$3.1^{+0.4}_{-0.4}$	$3.6^{+0.4}_{-0.4} \times 10^{56}$	430+150	$0.09^{+0.03}_{-0.03}$	180
GW151012	$23.3^{+14.0}_{-5.5}$	$13.6^{+4.1}_{-4.8}$	$15.2^{+2.0}_{-1.1}$	$0.04^{+0.28}_{-0.19}$	$35.7^{+9.9}_{-3.8}$	$0.67^{+0.13}_{-0.11}$	$1.5^{+0.5}_{-0.5}$	$3.2^{+0.8}_{-1.7} \times 10^{56}$	1060^{+540}_{-480}	$0.21^{+0.09}_{-0.09}$	1555
GW151226	$13.7^{+8.8}_{-3.2}$	$7.7^{+2.2}_{-2.6}$	$8.9^{+0.3}_{-0.3}$	$0.18^{+0.20}_{-0.12}$	$20.5^{+6.4}_{-1.5}$	$0.74^{+0.07}_{-0.05}$	$1.0^{+0.1}_{-0.2}$	$3.4^{+0.7}_{-1.7} \times 10^{56}$	440^{+180}_{-190}	$0.09^{+0.04}_{-0.04}$	1033
GW170104	$31.0^{+7.2}_{-5.6}$	$20.1^{+4.9}_{-4.5}$	$21.5^{+2.1}_{-1.7}$	$-0.04^{+0.17}_{-0.20}$	$49.1^{+5.2}_{-3.9}$	$0.66^{+0.08}_{-0.10}$	$2.2^{+0.5}_{-0.5}$	$3.3^{+0.6}_{-0.9} \times 10^{56}$	960^{+430}_{-410}	$0.19^{+0.07}_{-0.08}$	924
GW170608	$10.9^{+5.3}_{-1.7}$	$7.6^{+1.3}_{-2.1}$	$7.9^{+0.2}_{-0.2}$	$0.03^{+0.19}_{-0.07}$	$17.8^{+3.2}_{-0.7}$	$0.69^{+0.04}_{-0.04}$	$0.9^{+0.05}_{-0.1}$	$3.5^{+0.4}_{-1.3} \times 10^{56}$	320^{+120}_{-110}	$0.07^{+0.02}_{-0.02}$	396
GW170729	$50.6^{+16.6}_{-10.2}$	$34.3^{+9.1}_{-10.1}$	$35.7^{+6.5}_{-4.7}$	$0.36^{+0.21}_{-0.25}$	$80.3^{+14.6}_{-10.2}$	$0.81^{+0.07}_{-0.13}$	$4.8^{+1.7}_{-1.7}$	$4.2^{+0.9}_{-1.5}\times10^{56}$	2750^{+1350}_{-1320}	$0.48^{+0.19}_{-0.20}$	1033
GW170809	$35.2^{+8.3}_{-6.0}$	$23.8^{+5.2}_{-5.1}$	$25.0^{+2.1}_{-1.6}$	$0.07^{+0.16}_{-0.16}$	$56.4^{+5.2}_{-3.7}$	$0.70^{+0.08}_{-0.09}$	$2.7^{+0.6}_{-0.6}$	$3.5^{+0.6}_{-0.9}\times10^{56}$	990^{+320}_{-380}	$0.20^{+0.05}_{-0.07}$	340
GW170814	$30.7^{+5.7}_{-3.0}$	$25.3_{-4.1}^{+2.9}$	$24.2^{+1.4}_{-1.1}$	$0.07^{+0.12}_{-0.11}$	$53.4^{+3.2}_{-2.4}$	$0.72^{+0.07}_{-0.05}$	$2.7^{+0.4}_{-0.3}$	$3.7^{+0.4}_{-0.5} \times 10^{56}$	580^{+160}_{-210}	$0.12^{+0.03}_{-0.04}$	87
GW170817	$1.46^{+0.12}_{-0.10}$	$1.27^{+0.09}_{-0.09}$	$.186^{+0.00}_{-0.00}$	$0.00^{+0.02}_{-0.01}$	≤ 2.8	≤ 0.89	≥ 0.04	$\geq 0.1 \times 10^{56}$	40^{+10}_{-10}	$0.01^{+0.00}_{-0.00}$	16
GW170818	$35.5^{+7.5}_{-4.7}$	$26.8^{+4.3}_{-5.2}$	$26.7^{+2.1}_{-1.7}$	$-0.09^{+0.18}_{-0.21}$	$59.8^{+4.8}_{-3.8}$	$0.67^{+0.07}_{-0.08}$	$2.7^{+0.5}_{-0.5}$	$3.4^{+0.5}_{-0.7} \times 10^{56}$	1020^{+430}_{-360}	$0.20^{+0.07}_{-0.07}$	39
GW170823	39.6 ^{+10.0} _{-6.6}	$29.4^{+6.3}_{-7.1}$	$29.3^{+4.2}_{-3.2}$	$0.08^{+0.20}_{-0.22}$	$65.6^{+9.4}_{-6.6}$	$0.71^{+0.08}_{-0.10}$	$3.3^{+0.9}_{-0.8}$	$3.6^{+0.6}_{-0.9} \times 10^{56}$	1850^{+840}_{-840}	$0.34^{+0.13}_{-0.14}$	1651

Masses Spin



Slide credit: Nico Fernandez (UCSC → UIUC)

Effective Spin

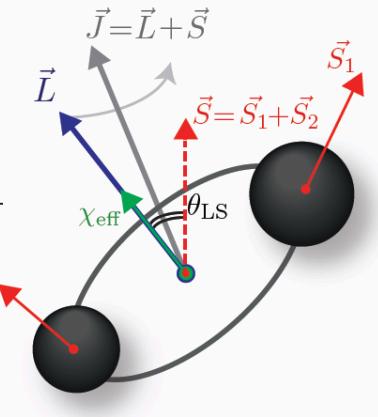
$$\chi = \frac{|\vec{S}|}{Gm^2}$$

Dimensionless spin parameter

$$\chi_{\text{eff}} = \frac{m_1 \chi_1 \cos \theta_1 + m_2 \chi_2 \cos \theta_2}{m_1 + m_2}$$

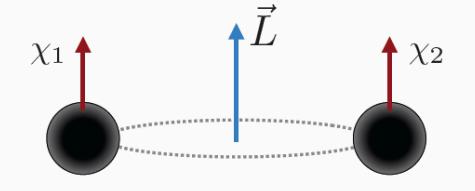


- Direction. +++
- Spin magnitude. ++
- masses. +



Effective Spin = 1

$$\chi_{\text{eff}} = \frac{m_1 \chi_1 \cos \theta_1 + m_2 \chi_2 \cos \theta_2}{m_1 + m_2}$$
$$\cos \theta = \hat{\chi} \cdot \hat{L}$$





 $\chi_{\rm eff} \approx 1$

Most black holes from stellar binaries probably start off with their spins aligned

Effective Spin = 0

$$\chi_{\mathrm{eff}} = m_1 + m_2$$

$$\cos \theta = \hat{\chi} \cdot \hat{L}$$

$$\chi_{\mathrm{eff}} \approx 0$$

 $m_1 \chi_1 \cos \theta_1 + m_2 \chi_2 \cos \theta_2$

Spins are essentially isotropic in the dynamical formation scenario. Binary was probably formed in a cluster

Effective Spin = 0

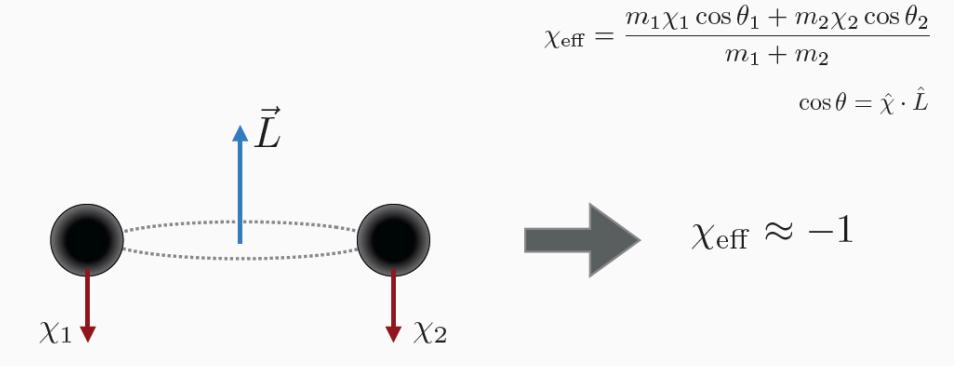
$$\chi_{\text{eff}} = \frac{m_1 \chi_1 \cos \theta_1 + m_2 \chi_2 \cos \theta_2}{m_1 + m_2}$$

$$\cos \theta = \hat{\chi} \cdot \hat{L}$$

$$\chi_{\text{eff}} \approx 0$$

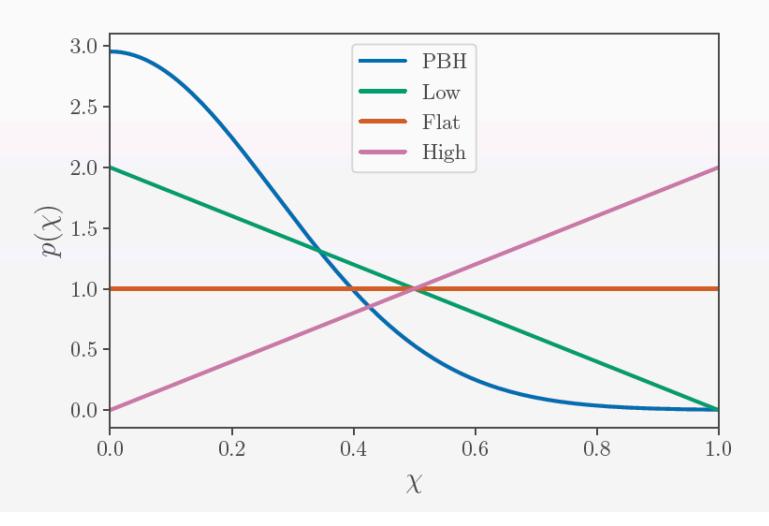
Spin magnitudes are close to zero (expected from PBHs).

Effective Spin = -1



Both spins are anti-aligned with its orbit (rare)

Magnitude Spin Priors



Fernandez and Profumo, 1905.13109 (JCAP); Slide credit: Nico Fernandez (UCSC → UIUC)

Model Selection

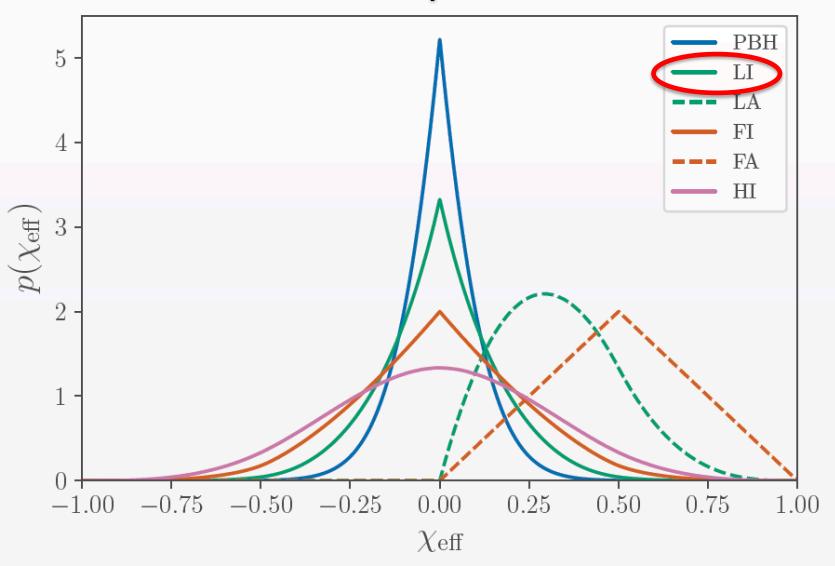
- Spin magnitude: Low (L), Flat (F),
 High (H) and PBH
- Spin orientations: Isotropic (I) and Aligned (A)

Example:

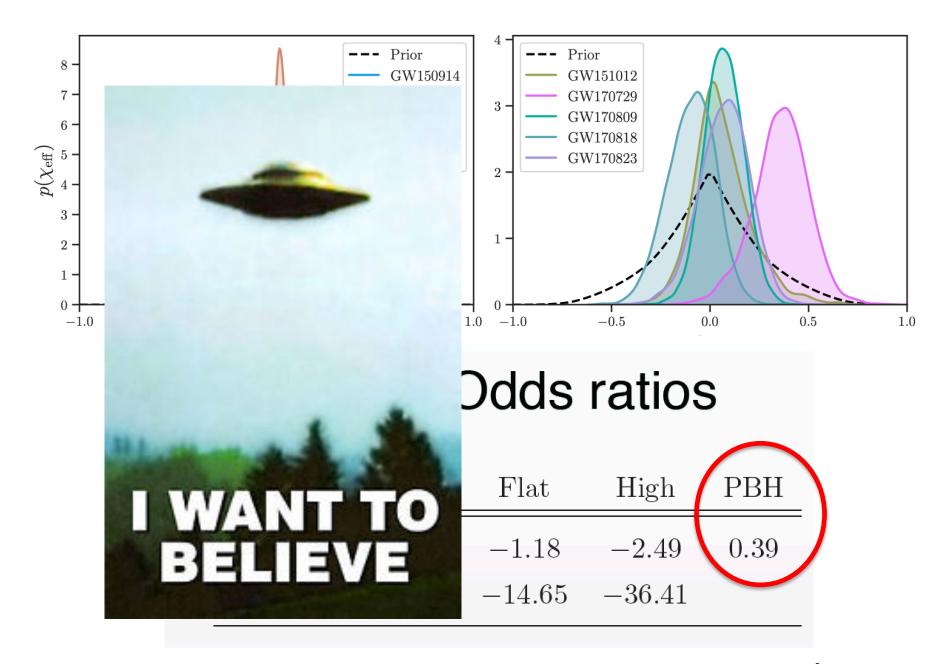
FI = Flat spin magnitude and isotropic spins (LIGO)

FA = Flat spin magnitude and align spins

Effective Spin Priors

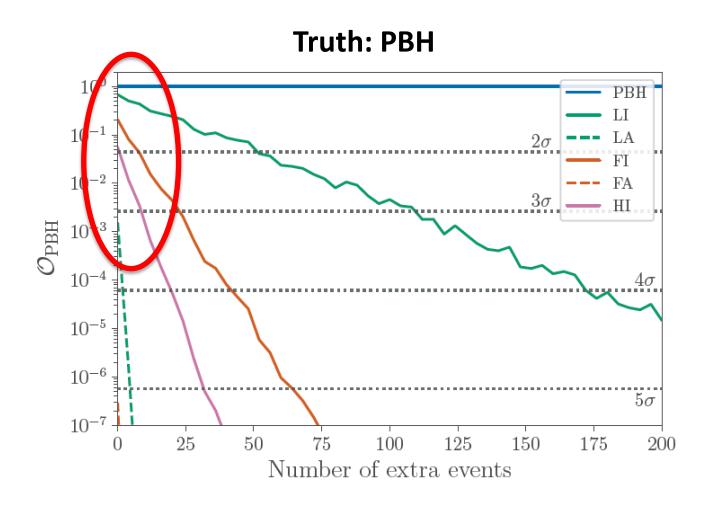


Fernandez and Profumo, 1905.13109 (JCAP); Slide credit: Nico Fernandez (UCSC → UIUC)

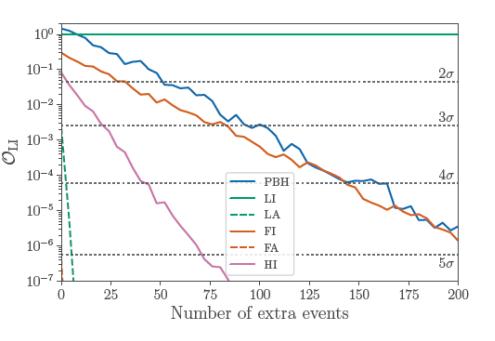


Fernandez and Profumo, 1905.13109 (JCAP); Slide credit: Nico Fernandez (UCSC → UIUC)

Evolution of the Odds ratios



Evolution of the Odds ratios

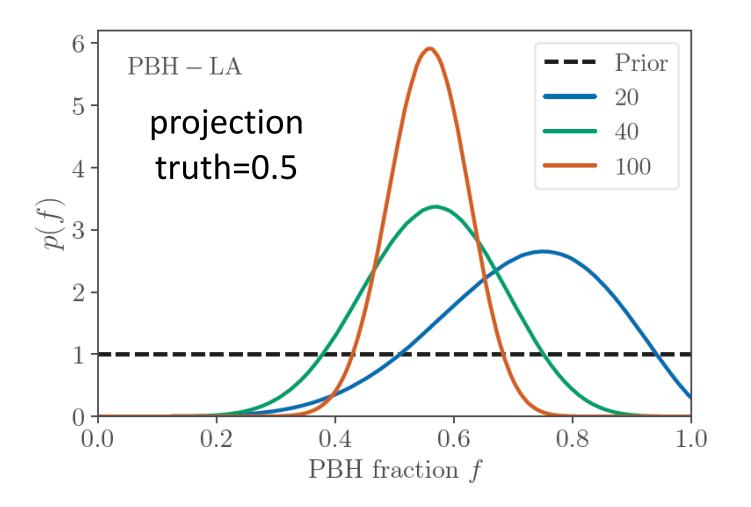


Truth: Low-isotropic

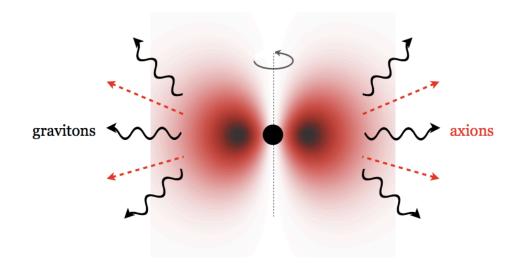
Fernandez and Profumo, 1905.13109 (JCAP); Slide credit: Nico Fernandez (UCSC → UIUC)

What about mixed models?

What about mixed models?

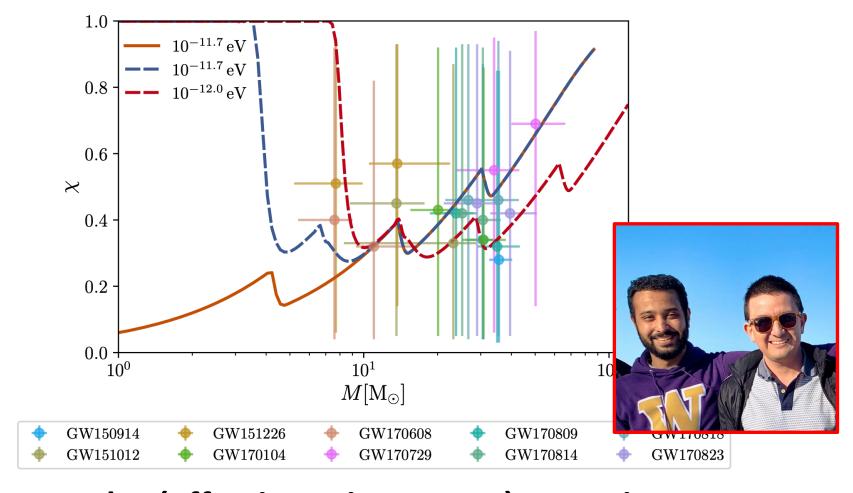


Fernandez and Profumo, 1905.13109 (JCAP); Slide credit: Nico Fernandez (UCSC → UIUC)



Assuming an initial spin and alignment distribution, one can compute the "best-fit" axion mass

Similarly, spin measurements can put constraints on axion-like particles



Regge plot (effective spin vs mass) assuming Flat priors for both mass and spin*

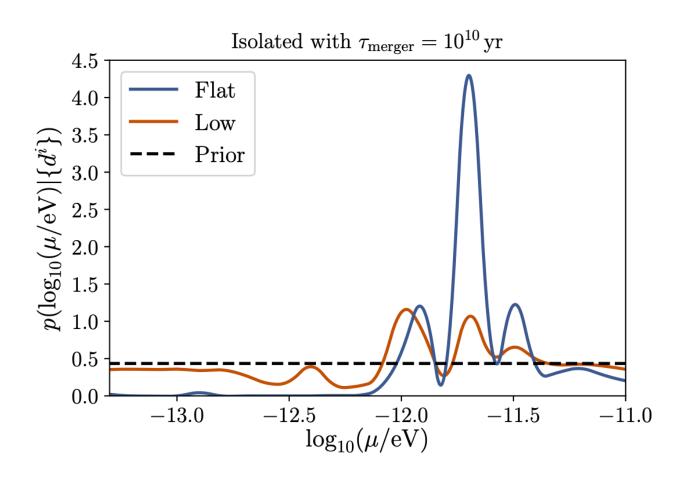
*Fernandez, Ghalsasy, Profumo, 1911.07862



binaries tend to have large spins...

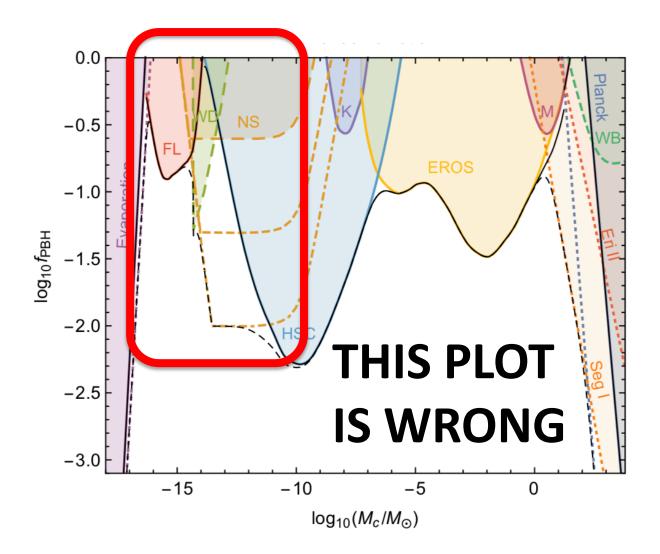
... but these are massive, so high-I is non-super-radiant!

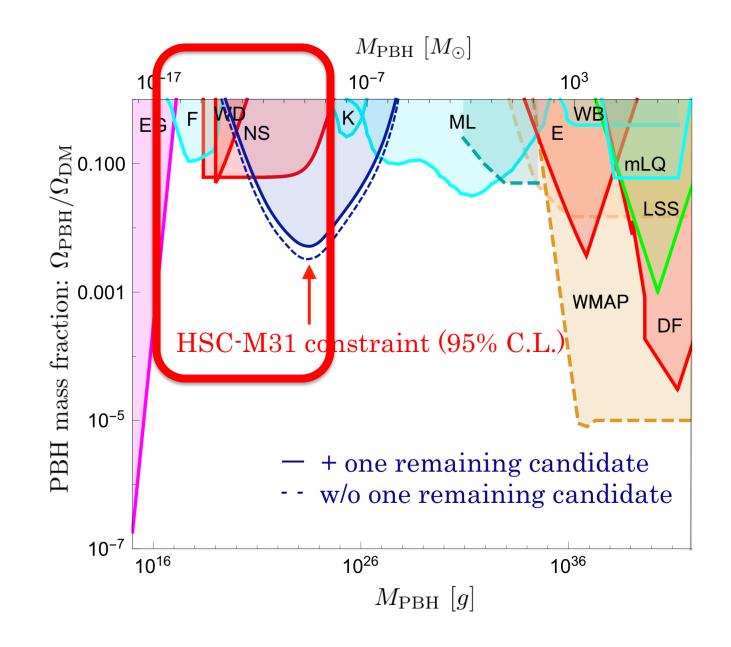
^{*}Fernandez, Ghalsasy, Profumo, 1911.07862



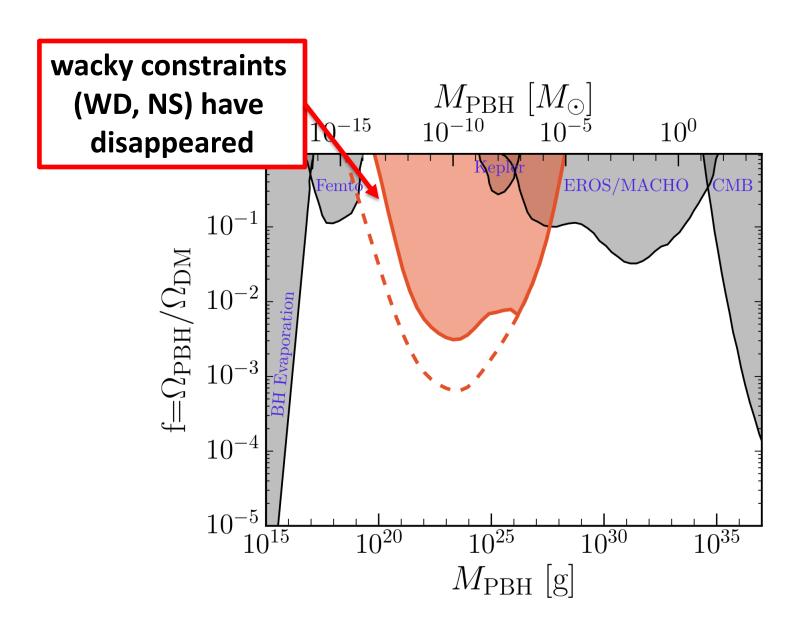
Posterior Probability for ALP mass

*Fernandez, Ghalsasy, Profumo, 1911.07862

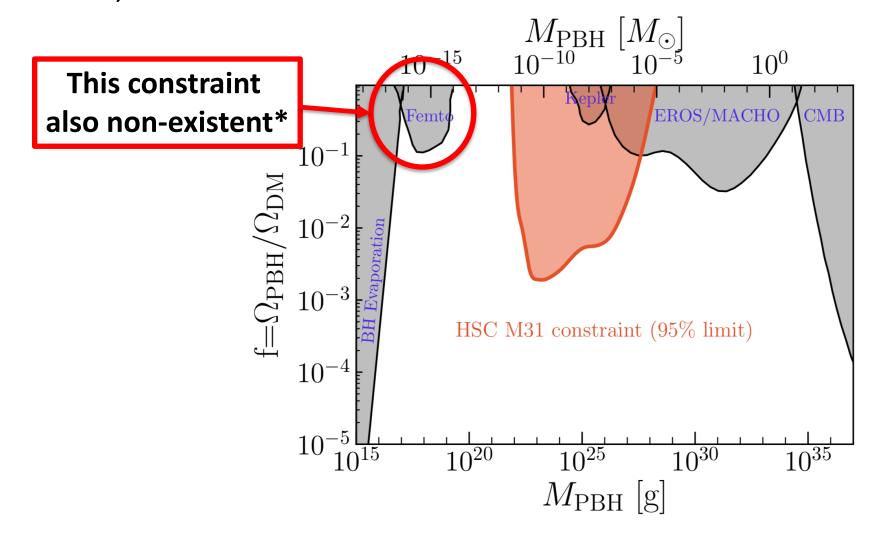




SUBARU HSC microlensing, 1701.02151 VERSION 1



SUBARU HSC microlensing, 1701.02151 VERSION 2: wave effects



SUBARU HSC microlensing, VERSION 3: finite source AND wave effects

...but assuming all stars have $R = R_{sun}$!

...but are these bounds robust?

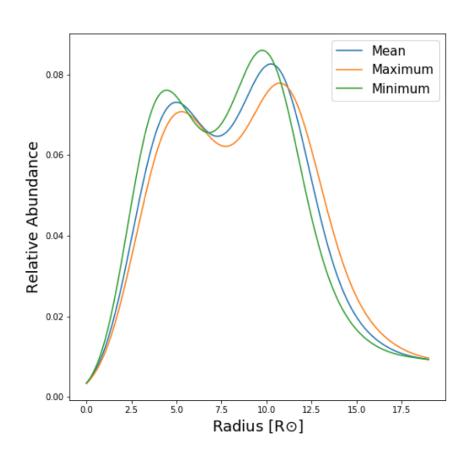
A few (worrisome) assumptions:

- > All stars are at the same distance
- > All stars have the same size (1 R_{sun})
- > DM is completely smooth



* Smyth, Profumo et al, 1910.01285, PRD

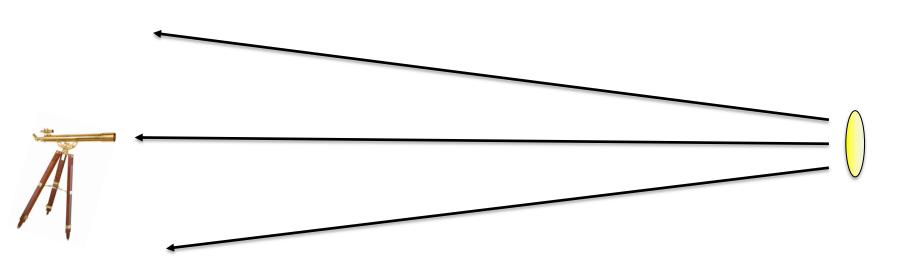
Sun-like stars are however too dim for HSC!

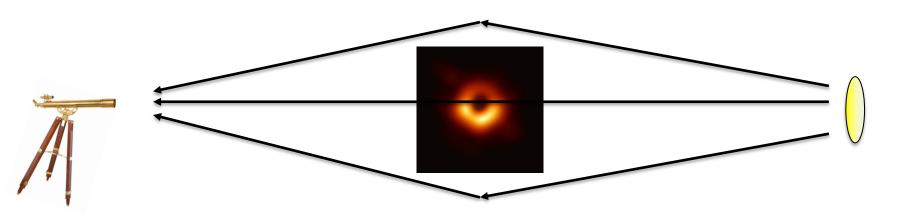


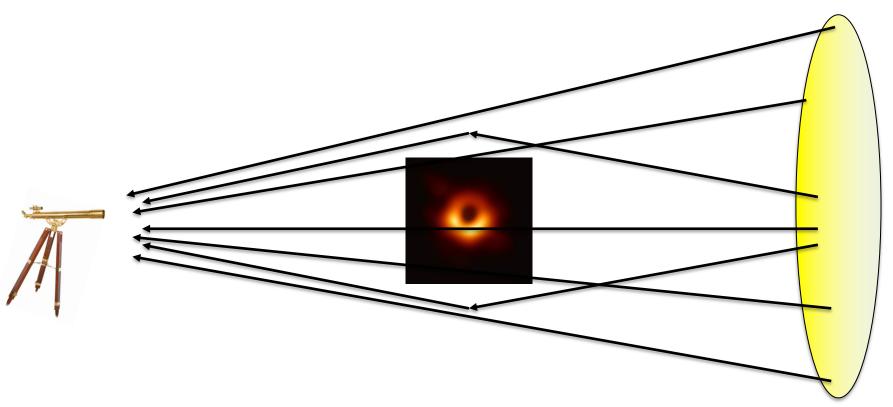


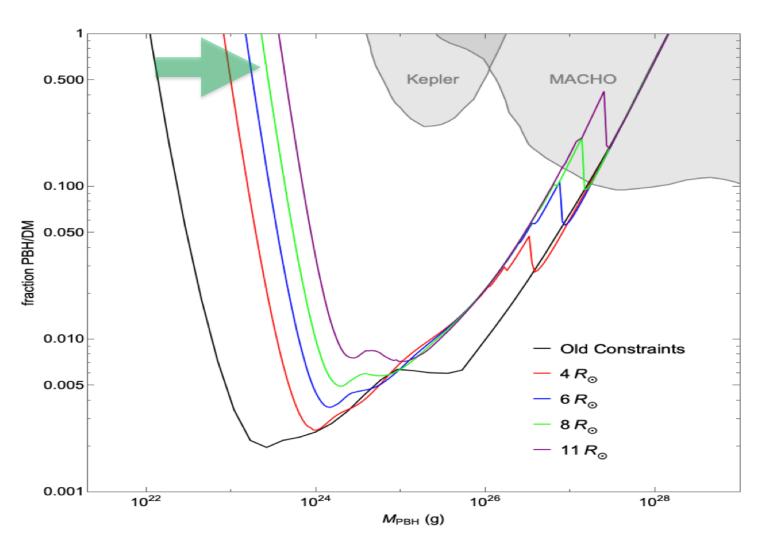
Stars that contribute to the microlensing constraints are ~ 100x larger in the sky than the Sun!

^{*} Profumo, Smyth+ PRD 2020

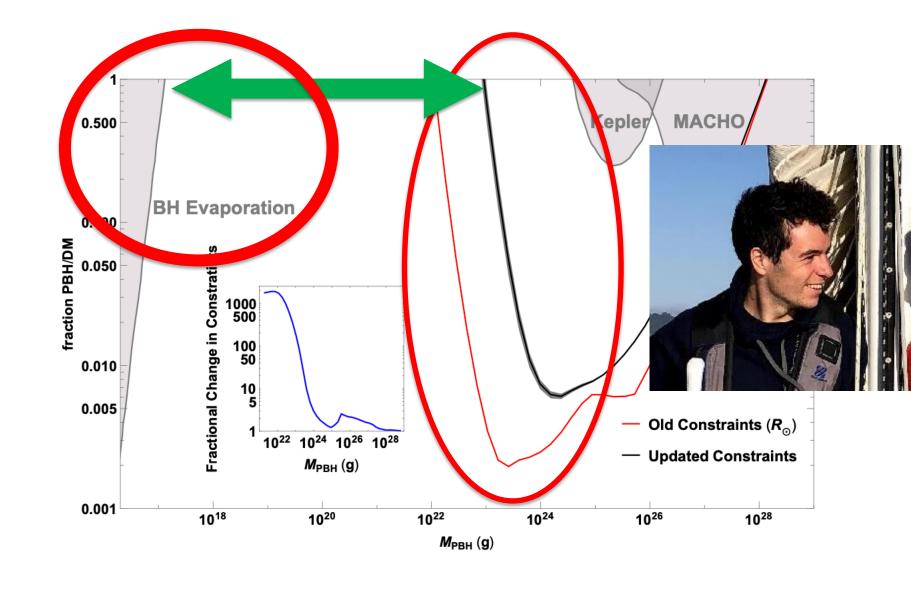








^{*} Profumo, Smyth+ PRD 2020



How do we go after them? Capture and perturbation around PSR?

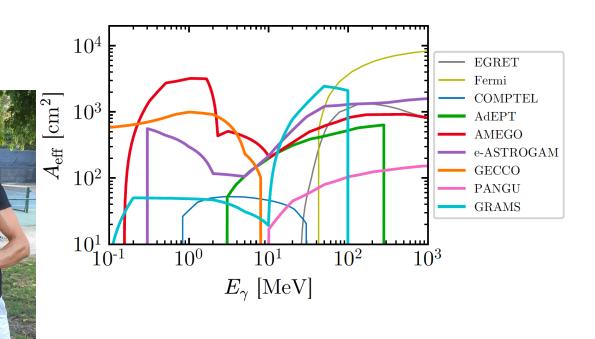
* Profumo, Smyth+ PRD 2020

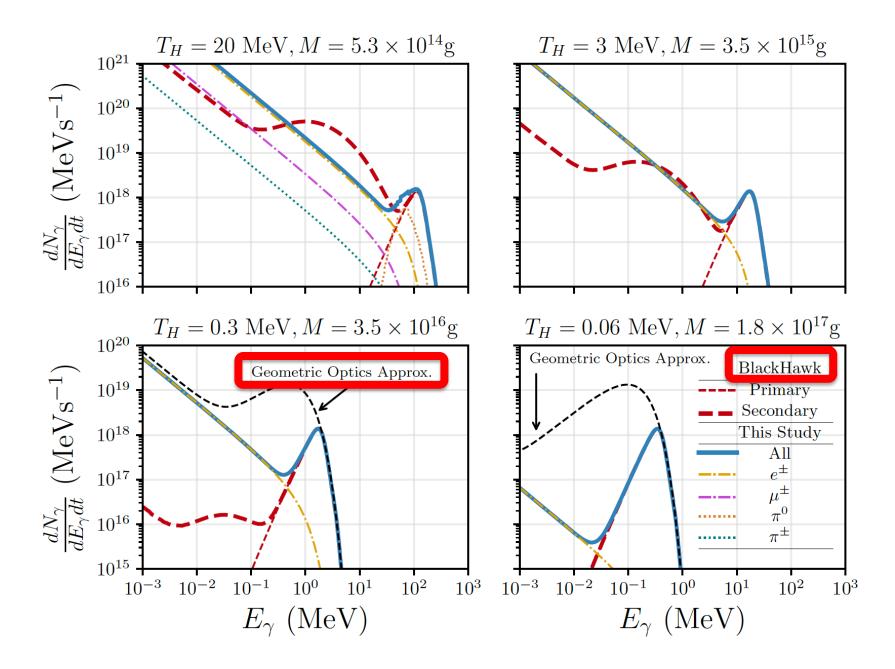
Lightest PBH that can be dark matter...

$$\tau(M) \simeq 200 \ \tau_U \left(\frac{M}{10^{15} \ \mathrm{g}}\right)^3 \simeq 200 \ \tau_U \left(\frac{10 \ \mathrm{MeV}}{T_H}\right)^3$$

are ~ asteroid/comet/PYRAMID mass

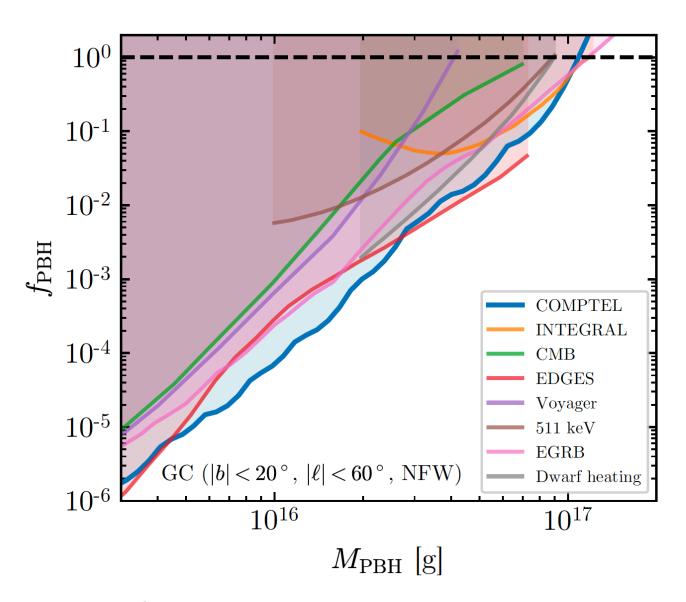
can't be much hotter than 10 MeV





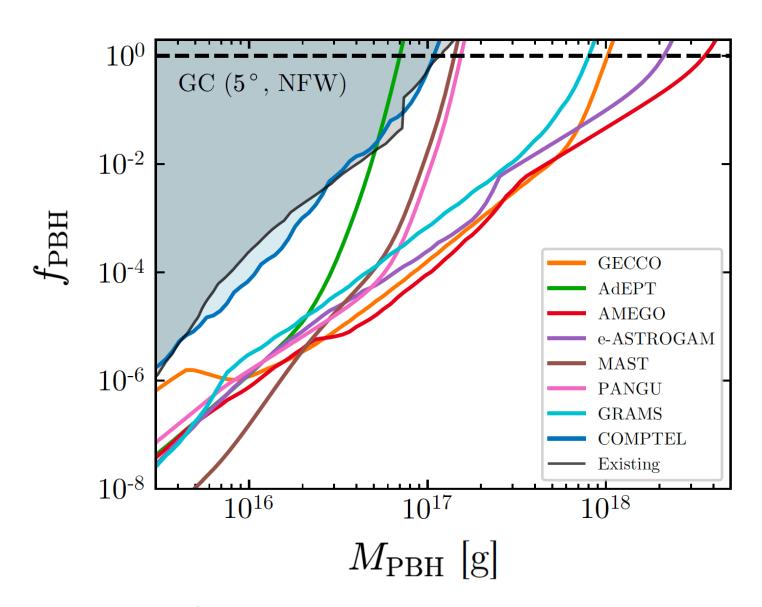
Coogan, Morrison & Profumo, 2010.04797

Our new COMPTEL constraints are among strongest/robust



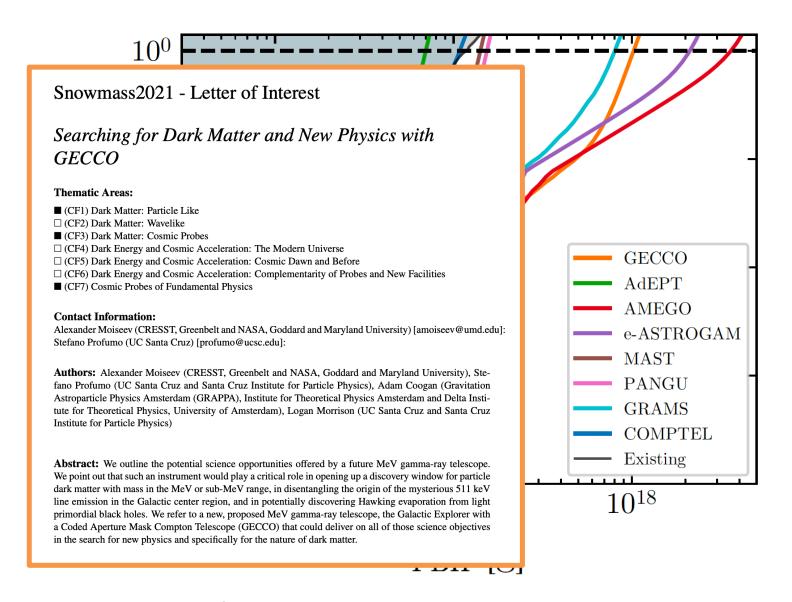
Coogan, Morrison & Profumo, 2010.04797

New MeV Telescopes could discover Hawking evaporation!



Coogan, Morrison & Profumo, 2010.04797

New MeV Telescopes could discover Hawking evaporation!



Coogan, Morrison & Profumo, 2010.04797

...even if PBH are NOT the dark matter, they can PRODUCE the dark matter via Hawking evaporation!





John Tamanas

Country	WCA ID	Gender	Competitions
United States	2007TAMA02	Male	41

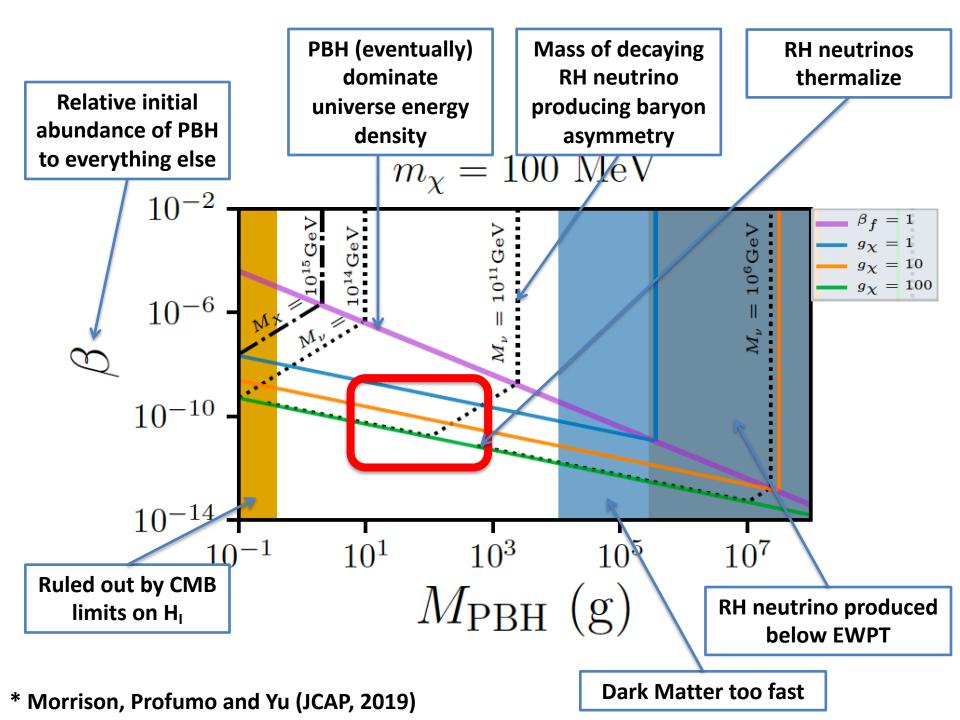
Current Personal Records

Event	NR	CR	WR	Single	Average
3x3x3 Cube	330	424	1485	8.16	10.13
2x2x2 Cube	195	265	901	1.55	3.49
4x4x4 Cube	1115	1644	7465	51.91	58.40
5x5x5 Cube	1654	2403	9997	2:28.52	2:43.81
■ 3x3x3 Blindfolded	666	900	4609	5:47.28	

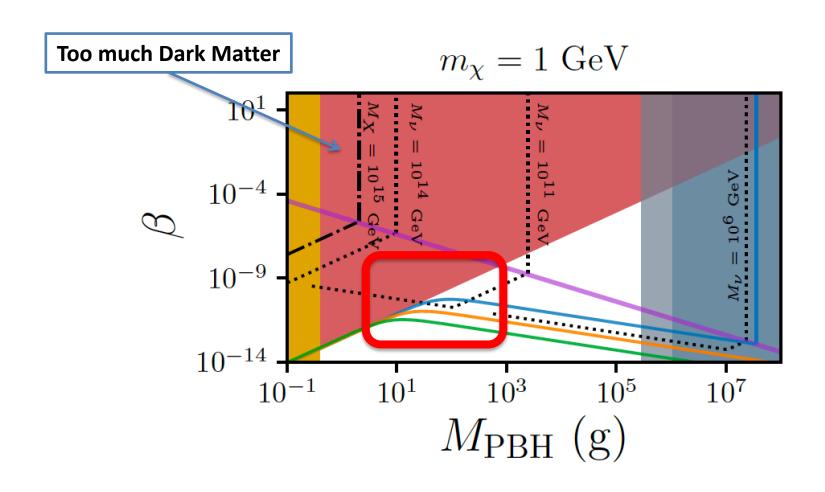
...even if PBH are NOT the dark matter, they can PRODUCE the dark matter via Hawking evaporation!

Mass (g)	$T_H ext{ (GeV)}$	τ (s)	$T_{\rm evap} = T(\tau) \; ({\rm GeV})$
$5M_P \simeq 10^{-4}$	1.7×10^{17}	10^{-41}	2×10^{17}
1	1.7×10^{13}	4×10^{-29}	2×10^{11}
10^{3}	1.7×10^{10}	4×10^{-20}	6×10^{6}
10^{6}	1.7×10^{7}	4×10^{-11}	200
10^{9}	1.7×10^{4}	0.04	0.006
10^{12}	17	$4 \times 10^7 \sim 1 \text{ yr}$	$\sim 1 \; \mathrm{keV}$

^{*} Morrison, Profumo and Yu (JCAP, 2019)



Dark Matter can be a mix of Planck-scale relics from PBH evaporation, and stuff the PBH evaporated into!



^{*} Morrison, Profumo and Yu (JCAP, 2019)

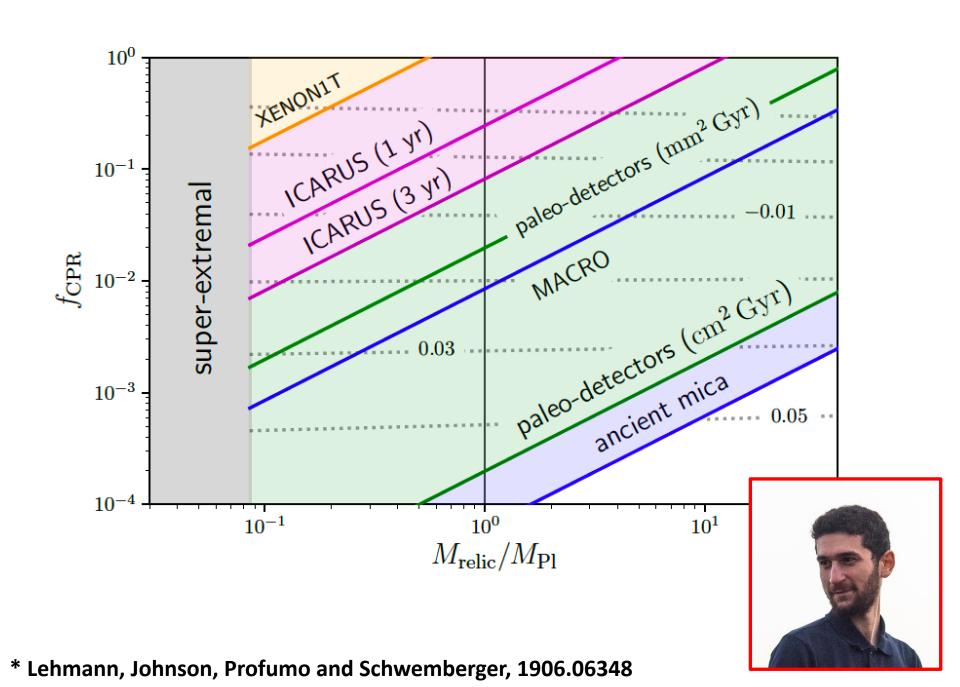
As BH approach the Planck scale, they can acquire a significant relic electric charge

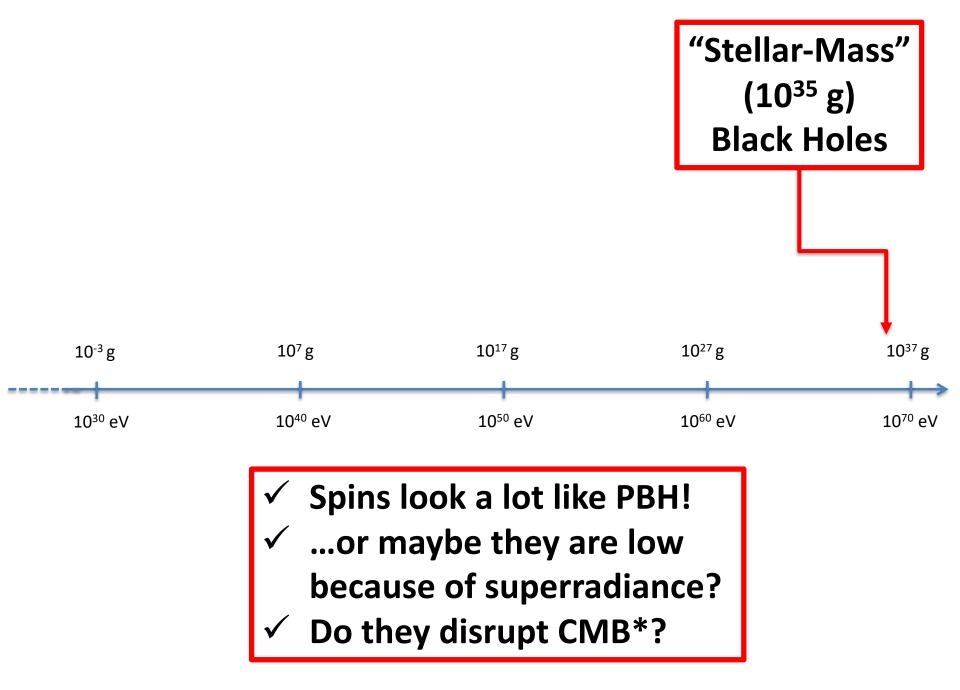
(under simple assumptions) $P(Q) \sim \exp\left(-4\pi\alpha(Q/e)^2\right)$ the relic charge is approximately Gaussian* $(8\pi\alpha)^{-1/2} \approx 2.34$

If evaporation stops around the Planck scale (because of extremality, or because of quantum gravity) we are left with a population of charged, Planck-scale relics!

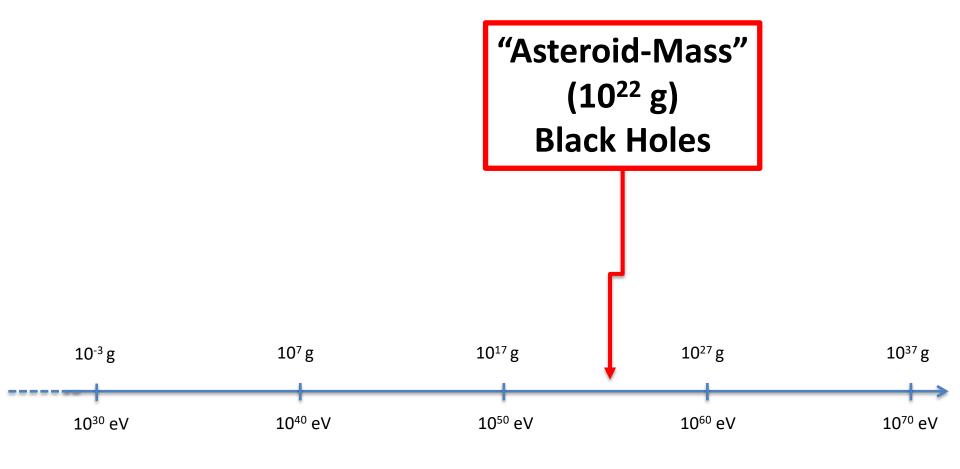
^{*} Page, 1977

^{**} Lehmann, Johnson, Profumo and Schwemberger, 1906.06348

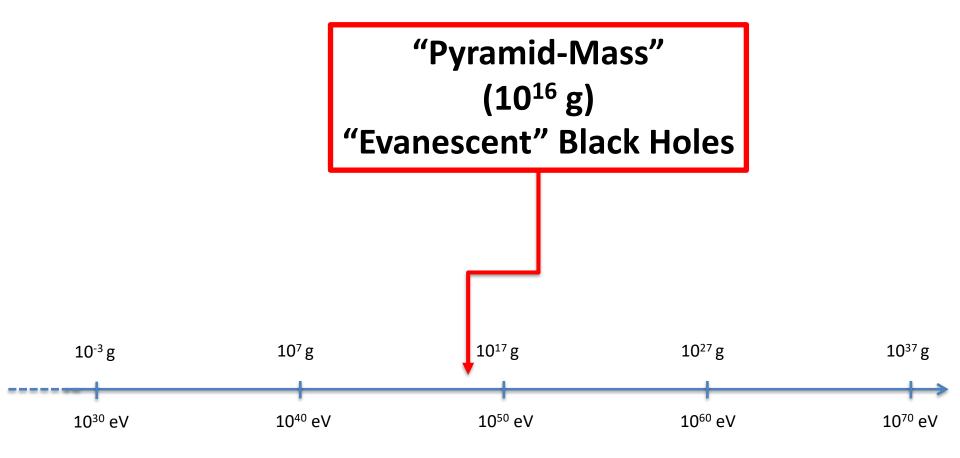




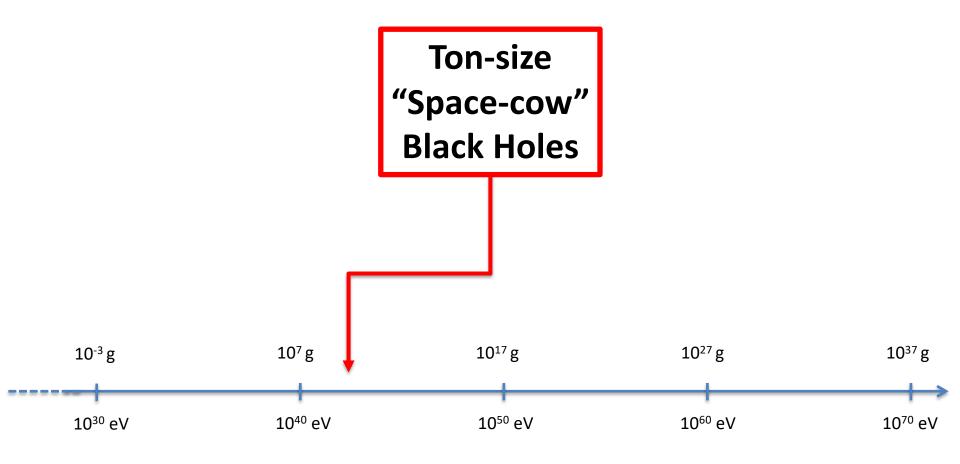
^{*} Gaspari, Lehmann, Profumo, in preparation



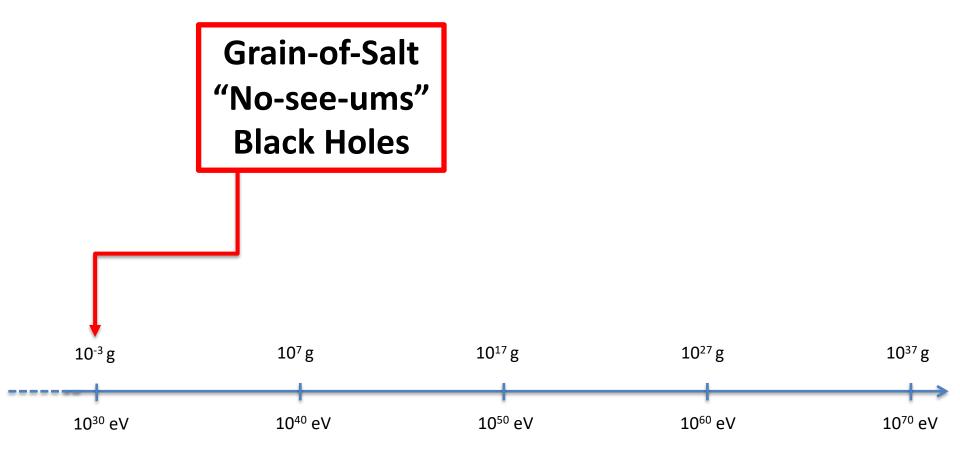
- ✓ Microlensing a lot trickier than previously thought!
- **✓** Detection strategies? PTA?



- ✓ Best constraints: COMPTEL
- ✓ Future MeV telescopes



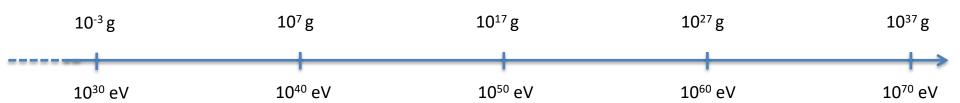
✓ Decays can produce DM, BAU, Planck relics





✓ Detectable!

In the era of gravitational wave astronomy,
the physics of macroscopic DM candidates
offers many opportunities for the ingenuity
of theorists and the craft of observers



Merger rate calculation (Cheng+Huang, 2018; Raidal +, 2017)

$$\tilde{\tau}(m_1, m_2, m_3) = \frac{348}{85} \frac{\alpha^4 \beta^7 a_{\text{eq}}^4 m_3^7 \tilde{x}(m_1, m_2)^4}{G^3 m_1 m_2 (m_1 + m_2)^8}.$$

$$\mathcal{G}(\psi; m_1, m_2, m_3) = \Gamma\left(\frac{58}{37}, \frac{\tilde{N}(\psi; m_1, m_2) t^{3/16}}{\tilde{\tau}(m_1, m_2, m_3)^{3/16}}\right) - \Gamma\left(\frac{58}{37}, \frac{\tilde{N}(\psi; m_1, m_2) t^{-1/7}}{\tilde{\tau}(m_1, m_2, m_3)^{-1/7}}\right),$$

$$\mathcal{R}(m_1, m_2) = \frac{9\bar{m}(\psi)^3 \tilde{N}(\psi; m_1, m_2)^{\frac{53}{37}}}{296\pi\delta_{\text{dc}}\tilde{x}(m_1, m_2)^3 t^{34/37}} \times \frac{\psi(m_1)\psi(m_2)}{m_1m_2} \int dm_3 \frac{\mathcal{G}(\psi; m_1, m_2, m_3)}{\tilde{\tau}(m_1, m_2, m_3)^{3/37}} \frac{\psi(m_3)}{m_3}.$$

