

Periodically driven Langevin systems with vanishingly small viscous drives

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Shakul Awasthi and SD (arXiv:2402.16512)

- Periodically driven thermodynamic systems can exist in time-dependent nonequilibrium states (Oscillating states)
- Prototypical example:
Driven under-damped Brownian particle
 - Viscous drive $\gamma(t)$; Noise drive $D(t)$
 - Potential drive $U(x; \lambda)$ with $\lambda \rightarrow \lambda(t)$
- These OS are characterized by $T_b = D/\gamma$ and γ .
 - What are the properties of OS for given T_b as $\gamma \rightarrow 0$?
- Vanishing viscous limit is singular.
 - We have a singular perturbation theory
- We find two classes of OS behaviour in low viscous limit

Driven under-damped Brownian motion

- Fokker-Planck equation:

$$\frac{\partial p}{\partial t} = \left[-\frac{\partial}{\partial x} v - \frac{\partial}{\partial v} (-\gamma v + f) + D \frac{\partial^2}{\partial v^2} \right] p$$

- Harmonic $f_h = -kx$ and Anharmonic: $f_{ah} = \sum \lambda_n x^n$
- When undriven, the large-time distribution is equilibrium
- When driven
 - viscous drive $\gamma(t)$; noise drive $D(t)$,
 - mechanical drive: $k(t)$ and $\{\lambda_n(t)\}$

then the large-time state is **Oscillating state** (OS):

$$\lim_{N \rightarrow \infty} p(x, v, NT + t) = p_{os}(x, v, t)$$

Low-viscous drives

- Perturbative parameter α
 $\gamma_\alpha = \alpha\gamma(t)$ and $T_b = D_\alpha/\gamma_\alpha$
- Inertial term is important; effects of drive are drastic.
- The limit is singular:

$$\lim_{\alpha \rightarrow 0} \text{Fokker-Planck} = \text{Liouville}$$

- large t and small α do not commute:

$$\lim_{\alpha \rightarrow 0} \lim_{N \rightarrow \infty} \neq \lim_{N \rightarrow \infty} \lim_{\alpha \rightarrow 0}$$

Singular perturbation theory

- Alternative way of identifying OS is demand periodicity:

$$P(t + T; \alpha; P_{in}) = P(t; \alpha; P_{in})$$

- Follow the standard perturbative steps
 - $L_{FP} = L_0 + \alpha L_1$ and $P(t) = \sum \alpha^n P^{(n)}(t)$
 - Solve the hierarchy of equations for arbitrary $P_{in}^{(n)}$
where $P_{in} = \sum \alpha^n P_{in}^{(n)}$.
- Impose the non-standard periodicity demand (Floquet theory)
 - Periodicity at $O(\alpha^n)$ uniquely determines OS to $(n - 1)$ -th order.
- The algorithm can be employed to moments dynamics, to determine OS moments and in turn P_{OS} .

Two distinct cases

- Consider $U(x) = k_0 x^2$ with drives $\gamma(t)$ and $T_b(t)$ of period $T = 2\pi/\omega$.
- Case 1: When $k_0 \neq n_0^2 \omega^2 / 4$
 - $P_{os} \rightarrow P_{eq}$ as $\alpha \rightarrow 0$ with T_{eff} .
 - Steady temperature T_s within in spite of varying $T_b(t)$ without.
 - For small α , $T_s = T_s(t)$
- Case 2: When $k_0 = n_0^2 \omega^2 / 4$
 - As $\alpha \rightarrow 0$, P_{os} resonates with time-period T/n_0
 - For small α , OS become T -periodic
- The results hold beyond harmonic potentials.