

Non-Normalizable Wavefunctions in Extreme Gravitational Fields: A Mathematical Perspective

(A Mathematical Perspective on Quantum Confinement and Gravity)

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Introduction - Choosing the Right Frame

By using a non-relativistic framework, where we use the Laplacian ∇^2

$$abla^2\Psi = rac{1}{\sqrt{|g|}}rac{\partial}{\partial x^\mu}\left(\sqrt{|g|}g^{\mu
u}rac{\partial\Psi}{\partial x^
u}
ight)$$

I did introduce relativistic conditions when I placed the boundary near the event horizon. Earlier, I started with a non-relativistic 4D quantum well model to establish the fundamental structure of confined states. However, once I considered the Schwarzschild metric near the horizon, I had to transition to a relativistic

framework (What Boundary?)

The Need for a 4D Schrödinger Equation

Schrödinger Equation in Flat 4D Space

In a flat 4D space (ignoring relativistic effects), the time-dependent Schrödinger equation is

$$i\hbarrac{\partial\Psi}{\partial t}=-rac{\hbar^2}{2m}\left(rac{\partial^2\Psi}{\partial x^2}+rac{\partial^2\Psi}{\partial y^2}+rac{\partial^2\Psi}{\partial z^2}+rac{\partial^2\Psi}{\partial w^2}
ight)$$

Schrödinger Equation in Non-Euclidean 4D Space

$$i\hbarrac{\partial\Psi}{\partial t}=-rac{\hbar^2}{2m}\left(rac{1}{\sqrt{|g|}}rac{\partial}{\partial x^\mu}\left(\sqrt{|g|}g^{\mu
u}rac{\partial\Psi}{\partial x^
u}
ight)
ight)\Psi$$

The Schwarzschild metric (for a spherically symmetric non-rotating mass M) in four-dimensional curved spacetime is given by:

$$ds^2 = -\left(1 - rac{2GM}{r}
ight)c^2 dt^2 + rac{dr^2}{1 - rac{2GM}{r}} + r^2 d heta^2 + r^2 \sin^2 heta \, d\phi^2$$

Here;

$$|g| = r^4 \sin^2 heta$$

Thus, the Schrödinger equation in curved Schwarzschild space is:

$$i\hbarrac{\partial\Psi}{\partial t}=-rac{\hbar^2}{2m}\left[rac{1}{r^2}rac{\partial}{\partial r}\left(r^2\left(1-rac{2GM}{r}
ight)rac{\partial\Psi}{\partial r}
ight)+rac{1}{r^2\sin heta}rac{\partial}{\partial heta}\left(\sin hetarac{\partial\Psi}{\partial heta}
ight)+rac{1}{r^2\sin^2 heta}rac{\partial^2\Psi}{\partial\phi^2}
ight]$$

This equation incorporates the effects of gravitational curvature on a quantum wavefunction.

Solving the Wavefunction Using Schwarzschild Metric

energy eigenvalues for a 4D quantum well in Flat Space

$$E_{n_x,n_y,n_z,n_w} = rac{\hbar^2 \pi^2}{2m} \left(rac{n_x^2}{L_x^2} + rac{n_y^2}{L_y^2} + rac{n_z^2}{L_z^2} + rac{n_w^2}{L_w^2}
ight)$$

Now, lets assume a reflective boundary at $r=2GM+\epsilon$, meaning:

$$\Psi(2GM+\epsilon)=0.$$

Energy Eigenvalue

$$E_n=rac{\hbar^2 k_n^2}{2m}=rac{\hbar^2 \pi^2 n^2}{2m\epsilon^2}-rac{m}{2}$$

(How??)

(what is actually ϵ ?)

Updated Energy Eigenvalue

$$E_{n_r,\ell,m,n_w} = rac{\hbar^2}{2m} \left(rac{x_{\ell n}}{R}
ight)^2 + rac{2mL^2}{\hbar^2} \pi^2 n_w^2 - rac{GMm}{r} + rac{2mr^2\ell(\ell+1)}{\hbar^2} + rac{am}{r^3} - rac{m}{2}$$

<u>Understanding Epsilon (ε)</u>

- •The probability density of finding a particle at that boundary is zero.
- •Instead of escaping, the wavefunction is reflected back into the allowed region.

Why Use a Reflective Boundary for Normalization?

In Schwarzschild spacetime, the wavefunction often diverges at the event horizon (r=2GM), making it non-normalizable. By imposing a reflective boundary at a small distance away from the horizon (r=2GM+ ϵ),

we can:

- Prevent divergence: The wavefunction is forced to remain finite and well-behaved.
- Create a quantum well-like system: The wavefunction gets trapped between the reflective boundary and another potential barrier.

Expression for €

$$\epsilon = lpha rac{\hbar}{m} + 4GMeta^2 + rac{8\pi GM\hbar}{c^3}.$$

- Quantum penetration depth (\hbar/m) .
- Gravitational blueshift effects (GM-dependent term).
- Thermal quantum effects (Hawking temperature scale).

A Blind Approach!!!

1). The 4D Quantum Schrödinger Equation in curved spacetime:

$$\left(g^{\mu
u}
abla_{\mu}
abla_{
u}-rac{V_{ ext{eff}}}{\hbar}
ight)\Psi=irac{\partial\Psi}{\partial t}$$

After some re-arrangement

$$\left(
abla^2 - 4rac{V_{
m eff}}{\hbar}
ight)\Psi = 4irac{\partial\Psi}{\partial t}$$

• EFE

$$G_{\mu
u} + \Lambda g_{\mu
u} = rac{8 \pi G}{c^4} T_{\mu
u}$$

- The structure of both equations has some resemblance (a differential term, a modifying term, and a source term).
- It can be seen as a formal analogy, but it should not be interpreted as a direct equivalence.

2) Gravitational Wave Echo Prediction

Since the quantum well imposes a reflective boundary at $r=2GM+\epsilon$, we modify the boundary condition:

[Classical GR Prediction: $A_{ref} = 0$ (full absorption)

$$h_{\mu
u}(r=2GM+\epsilon)=A_{
m in}+A_{
m ref}$$

introduce a reflection coefficient R dependent on the energy eigenvalues of the quantum well:

$$\mathcal{R}(\omega)=e^{i\delta(\omega)}$$

where $\delta(\omega)$ is the phase shift caused by quantum confinement.

$$\Delta t_{
m echo} = 2 \int_{r_p}^{2GM+\epsilon} rac{dr}{1-rac{2GM}{r}}$$

Approximating for small ϵ :

$$\Delta t_{
m echo} pprox 4GM\log\left(rac{1}{\epsilon}
ight)$$

Is this logarithmic scaling matches predictions from LQG and String Theory?

Conclusion & Seeking Guidance

- Equivalence to EFE: No direct mapping due to imaginary terms and effective potential.
- Event Horizon Boundary: Reflective conditions may not fully capture quantum effects.
- Non-Normalizability: Gravity modifies quantum states, affecting probability interpretation.
- Relativistic Framework: Transition from non-relativistic to relativistic needs better justification.
- Mathematical Consistency: Can we fully relativize the 4D quantum well equation?

further clarification and guidance needed

- Refining the transition to a relativistic framework while keeping solvability.
- Exploring quantum gravity or effective field theories for justification.
- Defining reflective boundary conditions rigorously near the horizon.
- Understanding if non-normalizability reflects real information leakage.

Note: This is a mathematical framing only; physical validity needs further refinement.

THANKYOU