



# Molecular dynamics simulations of DNA–DNA and DNA–protein interactions

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The all-atom molecular dynamics method can characterize the molecular-level interactions in DNA and DNA–protein systems with unprecedented resolution. Recent advances in computational technologies have allowed the method to reveal the unbiased behavior of such systems at the microseconds time scale, whereas enhanced sampling approaches have matured enough to characterize the interaction free energy with quantitative precision. Here, we describe recent progress toward increasing the realism of such simulations by refining the accuracy of the molecular dynamics force field, and we highlight recent application of the method to systems of outstanding biological interest.

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## Introduction

The very utility of DNA as the carrier of hereditary information is derived from its ability to undergo hybridization, whereby two DNA strands or a DNA and an RNA strand form a double-stranded (ds) DNA helix or a DNA/RNA duplex in accordance with the pattern of complementary interactions between the strands' bases, [Figure 1](#). RNA is also known to form higher-order structures, such as riboswitches in mRNA and ribozymes, driven by both canonical base pairing and higher-order non-canonical interactions. Other types of DNA–DNA interactions, such as polycation-mediated interactions

between dsDNA molecules and end-to-end base stacking interactions are central to the processes of DNA condensation [1] and the repair of double-stranded DNA breaks [2].

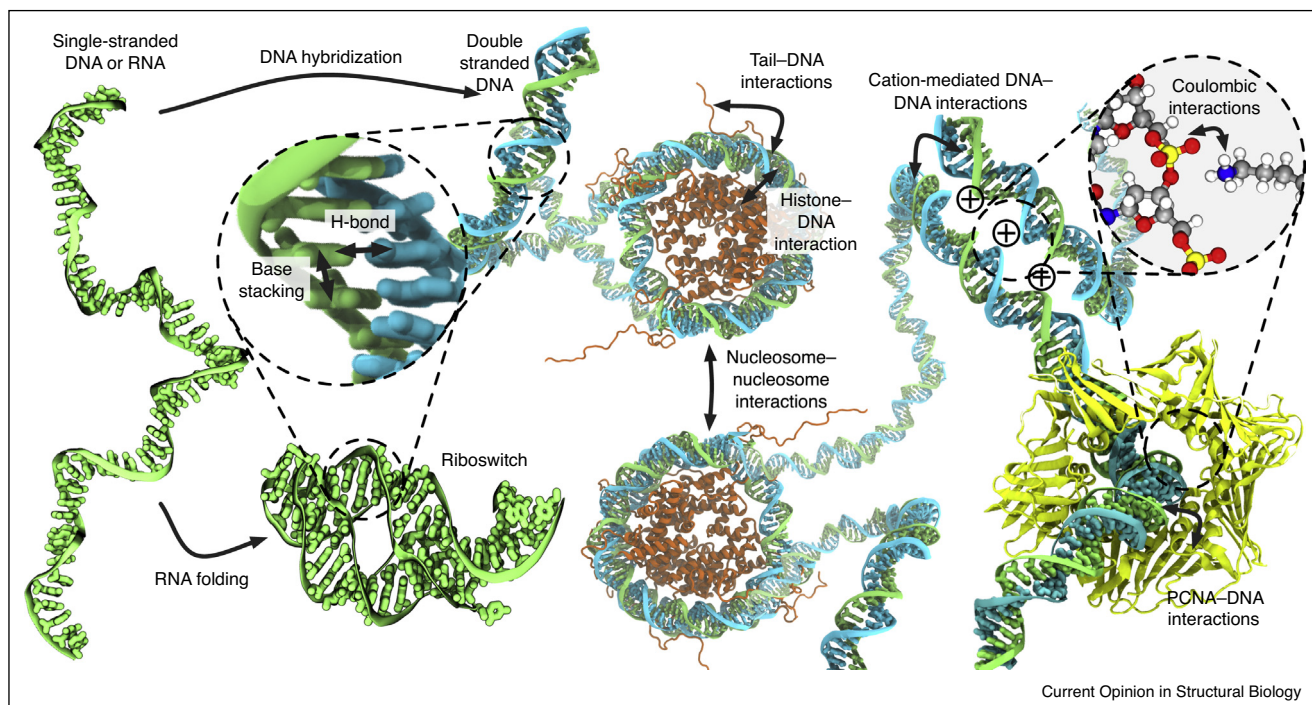
Protein–DNA interactions largely dominate the organization of DNA into higher-order structures. Thus, histone or histone-like proteins (in eukaryotes and prokaryotes, respectively) organize dsDNA at several scales, from individual protein–DNA complexes to higher-order clusters and fibers, and, ultimately, into entire chromosomes. The molecular mechanisms driving such multiscale organization are not yet fully understood, but they are known to play a central role in gene expression. Both transcription and translation require the concerted action of several protein and RNA species organized into massive supramolecular complexes [3]. Similarly, replication of genomic DNA requires the action of numerous DNA-binding proteins to synthesize DNA without losing the genome's integrity [4]. Protein–RNA interactions [5] play essential roles in emerging fields, such as gene editing by CRISPR–Cas [6] and in the formation of membraneless organelles in the nucleus and cytoplasm through phase separation [7]. The biological function of all the above systems critically depends on the strength of nucleic acid (NA)–protein interactions, as those should be strong enough to ensure stable binding of interaction partners that specifically recognize one another [8\*\*] but also weak enough to enable partner seeking by diffusion [9].

Matching their biological importance, NA–NA and protein–NA interactions have been studied extensively using the all-atom molecular dynamics (MD) approach. Paramount to the success of such studies is the precision of the molecular force field that prescribes the interaction strength between the chemical groups. However, recent long time-scale, quantitative characterization of NA–NA and NA–protein interactions revealed considerable imperfections of the existing molecular mechanics models [10,11\*\*,12,13\*]. Here, we describe recent advances in characterizing NA–NA and protein–NA interactions using all-atom MD simulation and increasing the realism of such simulations. Readers interested in a more comprehensive description of computational studies of DNA–protein systems are directed to a recent review [14\*\*].

## Molecular force fields

All-atom MD simulations of nucleic acids systems are typically performed using the AMBER or CHARMM force fields, which have been validated and improved

Figure 1



Examples of nucleic acid interactions in biological systems. From left to right: nucleic acids form secondary and tertiary structures; histone proteins organize DNA into nucleosomes and higher order structures; DNA interacts with a replication protein.

through multiple cycles of revisions. Presently, AMBER bsc0 [15], bsc1 [16], OL15 [17] and CHARMM36 parameters for RNA [18] and DNA [19] are the most up-to-date choices. A potential alternative is the OPLS force field [20] with the recently developed NA parameters [21], though it requires further validation by the community. For MD simulations of protein-NA systems, AMBER ff99SB [22] or ff14SB [23] or CHARMM36m [24] parameter sets for proteins contain the latest updates.

Historically, the improvement of NA force fields has been focused on the refinements of backbone and glycosidic torsion parameters [15,22,25]. Presently, both AMBER and CHARMM-based simulations can maintain the experimental double helical structure in tens of microseconds simulations [26], although some artifacts with the CHARMM36 simulations of longer dsDNA fragments have been reported recently [27]. It has long been recognized that MD simulations of unfolded proteins exhibit overly collapsed protein conformations [28–30,31\*,32–34], in part, because the TIP3P water model (used by both AMBER and CHARMM) energetically favors water–water interactions over water–protein interactions. Conversely, overly strong NA–NA and protein–NA interactions can be potential artifacts of standard AMBER and CHARMM simulations, which, in the case of peptide-mediated DNA–DNA interactions, leads to qualitatively incorrect

mulation outcomes [12]. Ion-specific effects are another area of concern, in particular, to simulations of dense NA systems [35].

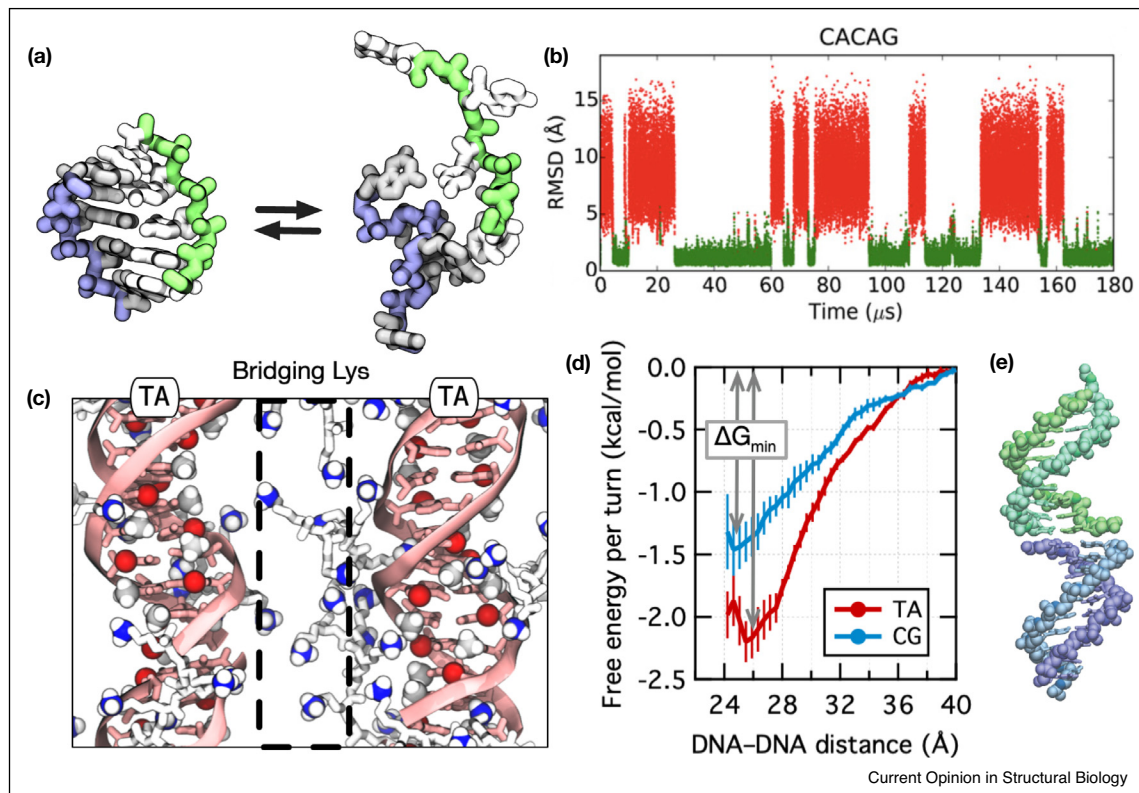
Recognizing the problem, several approaches have been developed to increase the realism of long time-scale simulations of NA systems, which we describe in the subsequent sections.

### NA–NA interactions

Hybridization is the most fundamental type of DNA and RNA self-assembly, a process in which hydrogen bonds and base stacking interactions stabilize the double helical structure in a nucleotide sequence-specific manner, Figure 2a. Because of its complexity and biological significance, the hybridization process has become a test ground for advanced simulation methods and force field refinement.

A typical problem encountered in the simulations of short oligomers hybridization or small RNA hairpin folding is the emergence of an intercalated base conformation [10]. Recognizing that the intercalation occurs because of the overestimated base–base stacking and base–backbone interactions, in comparison to base pairing, Chen and Garcia refined the Lennard-Jones (LJ) parameters of nucleobase atoms to weaken the stacking and strengthen

Figure 2



Hybridization of RNA and condensation of DNA. **(a)** Illustration of the hybridization process. **(b)** Root mean squared deviation of a CACAG RNA duplex from its folded configuration observed in a simulated tempering MD simulation. Panel (b) adapted from Ref. [11\*\*]. **(c)** Condensation of DNA duplexes mediated by poly-lysine peptides. **(d)** Condensation free energy of TA and CG repeat DNA duplexes. Panels (c) and (d) adapted from Ref. [40\*\*]. **(e)** End-to-end stacking of two DNA duplexes [47].

base pairing [10]. Following that, several revisions of the LJ parameters have been suggested by several groups [36–38].

The above-mentioned intercalation artifact is another manifestation of the imbalance between water–water and water–solute interactions. To remedy the problem, the Shaw group introduced the TIP4P-D water model and revised the AMBER ff14 RNA force field by optimizing the partial charges, the LJ and torsional parameters [11\*\*]. Note that a similar strategy was used previously by the Shaw group to refine the protein force field [31\*]. Using the revised force field, the Shaw group demonstrated dramatic improvements in the simulations of reversible hybridization of RNA duplexes, Figure 2b, as well as in the simulations of unfolded long ssRNAs, tetraloops, and riboswitches [11\*\*]. Although it has not been explicitly shown in the original study, one can expect to observe similar improvements for MD simulations of DNA hybridization.

The behavior of densely packed DNA systems, such as in a fully packed viral capsid, a folded chromosome or a

synthetic DNA nanostructure, sensitively depends on its ionic environment. When submerged in a monovalent cation solution, dsDNA molecules do not aggregate, regardless of the cation concentration. However, in the presence of tetravalent (or longer) basic peptides, for example, Lys<sub>4</sub>, dsDNA can spontaneously form a condensate (equivalently coacervate), in which dsDNA helices form a solvated, ordered phase [1,39]. This self-assembly process is driven by the Coulombic attractions between the basic groups of the peptide and the phosphate groups of DNA.

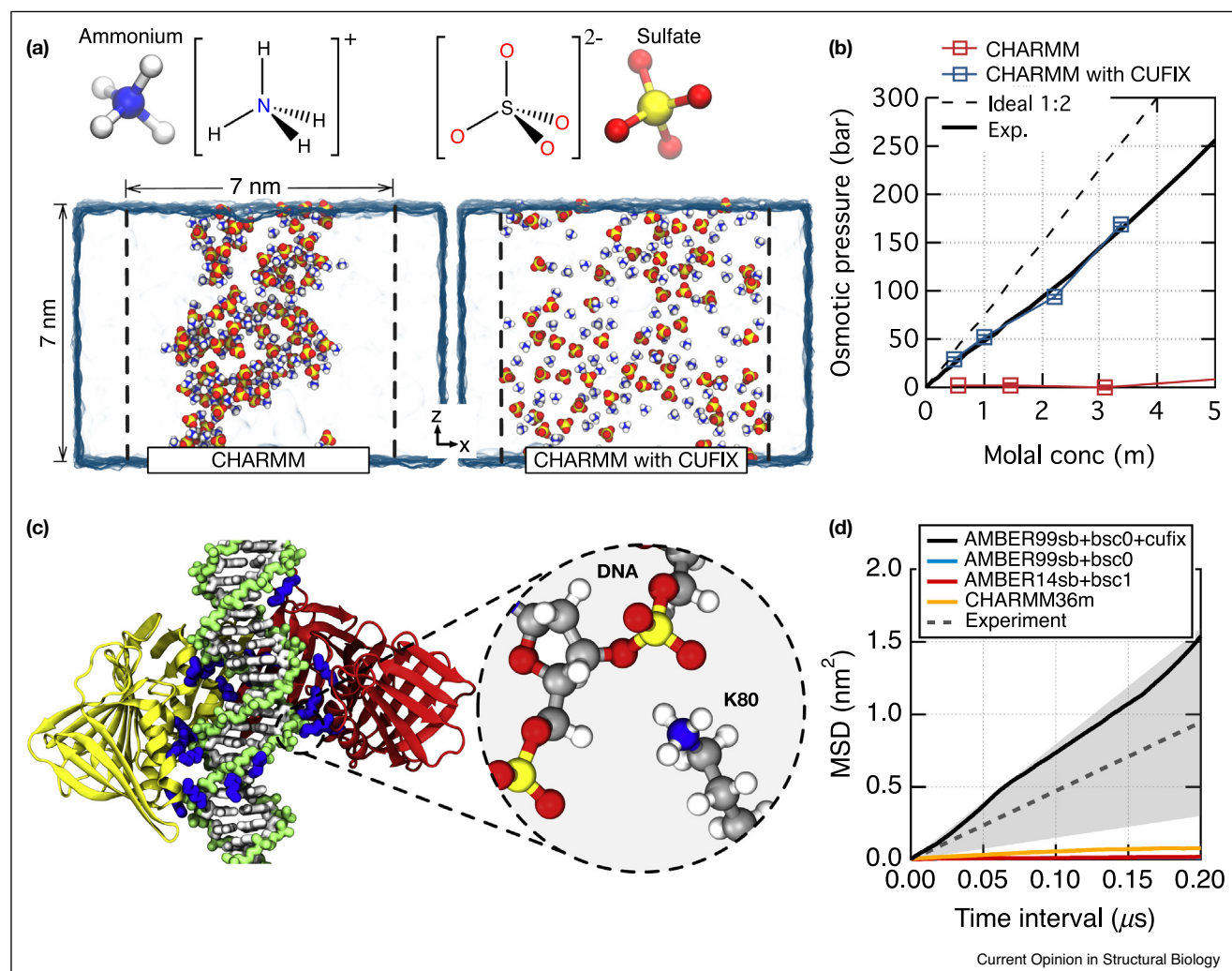
All-atom MD simulations of dsDNA molecules in aqueous solution of monovalent basic amino acids (e.g. Lys) showed pronounced aggregation of dsDNA for both standard AMBER and CHARMM parameter sets [12], a simulation artifact. The underlying cause of the aggregation was considerable overestimation of the Coulombic attractions between lysine sidechains and DNA [12]. Yoo and Aksimentiev refined the interaction strength using the experimental osmotic pressure of ammonium sulfate as a reference [12]. Briefly, ammonium sulfate solutions of various concentrations were simulated in a semi-permeable membrane setup,

Figure 3a, measuring the effect of the force field corrections on the osmotic pressure. The experimental osmotic pressure was recovered by increasing the LJ  $R_{\min}$  parameter for the amine nitrogen–sulfate oxygen pairs by about 0.16 Å, Figure 3b [12]. The simulations carried out using this non-bonded correction matched semi-quantitatively the experimental data on the magnitude of the DNA–DNA forces [12]. Readers interested in the development of the so-called CUFIX corrections are referred to a recent review [13\*].

The improved realism of MD characterization of DNA–DNA interactions enabled the prediction of previously unknown phenomena. MD simulations with the updated force field predicted that AT-rich dsDNA molecules would attract each other more strongly than GC-rich ones

when poly-lysine peptides mediate the inter-DNA interactions, Figure 2c,d [40\*\*]. This prediction immediately suggests that dsDNA molecules can undergo phase separation in a manner that depends on the DNA sequence. Both predictions were validated by single-molecule experiment [40\*\*,41] and bulk liquid–liquid phase separation assays [42]. Given that basic residues are critical functional groups of intrinsically disordered peptides that interact with DNA or RNA [43], the CUFIX corrections may improve MD description of such systems. A potential alternative approach to realizing accurate simulations of peptide-mediated DNA condensation could be the combination of the protein and nucleic acid force fields with improved water models [31\*,44\*], though this approach has not yet been validated.

Figure 3



Refinement of non-bonded interactions for MD simulation of a protein/DNA complex. **(a,b)** Calibration of charge-charge interactions against osmotic pressure data for aqueous solution of ammonium sulfate. Here, ammonium and sulfate are used as analogs of the protein's amine and DNA's phosphate groups, respectively. **(c)** All-atom model of a PCNA/DNA complex based on a crystal structure [51]. Lysine and arginine residues that form contacts with the DNA phosphates are highlighted in blue. **(d)** Mean squared displacement of PCNA on DNA simulated using four force field models and measured from experiment. Figures in panels (a)–(d) are adapted from Refs. [12] and [52], respectively.

In addition to its role in non-homologous double-strand DNA break repair, end-to-end stacking of dsDNA molecules, Figure 2e, has come to light as a possible mechanism driving the origin of life [45] and as a method of assembling complex synthetic DNA nanostructures [46]. MD simulations of the end-to-end stacking were found to overestimate the absolute magnitude of such interactions but correctly account for the effect of the termination chemistry [47]. Recent characterization of nucleotide type-specific base-stacking interactions [48] has set the stage for future refinement of the MD base-stacking models.

### Protein–NA interactions

Because of the enormous size of the genome, diffusion along DNA is essential for most DNA-binding proteins to find their targets [9,49]. An outstanding example is the DNA clamp, a ring-shaped protein complex that encircles dsDNA to ensure processive DNA replication in both eukaryotes and prokaryotes, Figure 3c. Because the diffusion coefficient of the eukaryotic clamp, proliferating cell nuclear antigen (PCNA), is greater than  $1 \text{ nm}^2/\mu\text{s}$  [50], a noticeable displacement of PCNA should be observed in a microsecond-long MD simulation. Contrary to that expectation, when a PCNA–dsDNA system is simulated using standard AMBER and CHARMM force fields, the mean-squared displacement (MSD) of the protein relative to dsDNA is orders of magnitude smaller than in experiment, Figure 3d [51,52]. Such a dramatic underestimation of MSD occurs because the lysine and arginine residues at the PCNA–DNA interface form long-lasting contact pairs with DNA phosphates, Figure 3c. In contrast, MSDs obtained from simulations carried out using the CUFIX corrections quantitatively match the experimental value, Figure 3d [52].

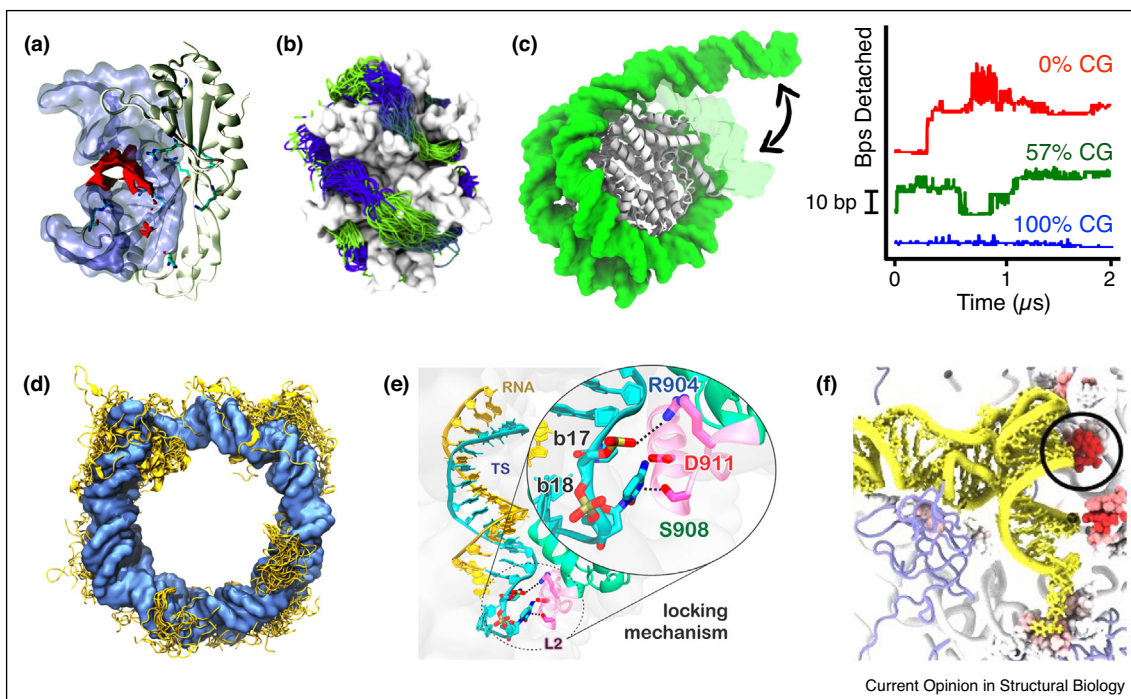
The natural outcome of a diffusive search is stable binding of the protein to a DNA fragment carrying a specific nucleotide sequence. Transcription factors (TFs) are one class of such sequence recognition proteins [9,49] that have been studied extensively through all-atom MD approaches [53,54\*,55,56\*,57\*\*,58\*]. An outstanding question in this area is the microscopic mechanism(s) enabling the recognition of the target DNA sequence. Extensive analysis of a large collection of protein–DNA complexes [8\*\*] has recently identified the sequence-specific deformation of a DNA fragment as a critical factor enabling the target recognition by a DNA binding protein, Figure 4a. The free energy of TF binding has been determined using advanced MD sampling methods such as alchemy perturbation [53] and umbrella sampling [54\*,55]. One study [52], however, implied that either standard AMBER or standard CHARMM force field systematically overestimates the binding free energy because the DNA–protein interactions typically involve direct charge–charge contacts at the protein–DNA interface [56\*], which standard MD force fields are known to overestimate. In spite

of such strong binding, TFs were observed to slide along dsDNA, albeit by a small amount, in a recent MD study [59]. Using the CUFIX corrections to charge–charge interactions is expected to reduce the strength of TF–DNA binding and to accelerate TF sliding along DNA.

In between stable binding and free diffusion lies a situation where a protein–DNA complex forms with a high affinity but remains amenable to rearrangement in response to external factors. *In vivo*, ssDNA is almost always sequestered by single-stranded DNA binding (SSB) proteins that bind ssDNA with high affinity. Despite the strong binding, SSBs can diffuse along ssDNA. The elementary steps of such a diffusion process — formation and diffusion of a small DNA bulge, Figure 4b — was recently observed in MD simulations [60]. Another example is the binding of dsDNA to histone proteins. The resulting assembly, the nucleosome, sequesters 147 bases of dsDNA, and many such nucleosomes can form higher-order structures guided by the interactions between DNA and intrinsically disordered histone tails or chromatin remodeling factors. Recent MD studies investigated DNA unwrapping from a protein core [61,62,63\*], and interactions between intrinsically disordered histone tails and the surrounding DNA [62,64,65,66,67\*]. Spontaneous and reversible nucleosome unravelling, Figure 4c, was observed in MD simulations carried out at elevated magnesium concentration [63\*]. Combined with NMR measurements, MD simulations uncovered how histone tails, and their chemical modifications, impede a zinc finger domain from binding to a nucleosome [64,67\*], Figure 4d. Other simulations characterized the electrostatic environment at the histone–DNA interface [68\*], and the effect of a centromere-specific histone variant on nucleosome elasticity [69\*]. Several recent all-atom simulations of multi-nucleosome systems investigated histone tail bridging interactions between two separate nucleosomes [65], and, combined with experiment and coarse-grained simulation, determined how a chaperone protein binds to a di-nucleosome [70]. All-atom MD simulations have also examined how relatively small histone-like proteins can strongly bend [71\*] or bridge [72] regions of bacterial DNA.

RNA–protein interactions feature prominently in both CRISPR–Cas9 and ribosomes. The CRISPR–Cas9 system has received significant attention because of its ability to edit genetic information in live cells [73,74]. Several recent all-atom MD studies have investigated the source of CRISPR–Cas9’s unintended interaction with off-target DNA sequences [75\*\*,76,77], one of which [75\*\*] identified a conformational ‘locking mechanism’, Figure 4e, which could be enhanced further through mutagenesis. Lastly, we highlight a simulation study of ssRNA’s interaction with the protein and RNA components of a ribosome, Figure 4f [78\*], which found diffusion of aminoacyl-tRNA to be significantly impeded by relatively few direct contacts with the ribosome.

Figure 4



All-atom simulations of protein–nucleic acid systems. **(a)** Binding of a protein (white) bends a segment of dsDNA (blue). Red shows a  $-7.4 \text{ kcal mol}^{-1}$  interaction potential isosurface. Adapted from Ref. [8\*\*]. **(b)** Ensemble of conformations explored by ssDNA bound to a single-stranded binding protein (white) within  $10 \mu\text{s}$ . The DNA is depicted as a tube that is colored from green to blue every 10 nt. Adapted from Ref. [60]. **(c)** Left, partial unwrapping of 601L DNA (green) from a histone core (white). Right, spontaneous unwrapping for three DNA sequences. Adapted from Ref. [63\*]. **(d)** Non-specific binding of histone tails (gold) to nucleosomal DNA (blue). Adapted from Ref. [67\*]. **(e)** MD simulation of a CRISPR-Cas9 system. The inset illustrates the interactions between the target strand of an RNA:DNA hybrid (TS, cyan) with loop 2 (L2, pink) of the protein catalytic domain. Adapted from Ref. [75\*\*]. **(f)** An ‘elbow’ of aminoacyl-tRNA (yellow) interacts with the RNA (white tubes) and protein (blue tubes) components of a ribosome. The black circle highlights a prominent contact between rRNA and tRNA. Adapted from Ref. [78\*].

## Conclusions

As the scope of all-atom MD simulations evolves from individual proteins or NAs to systems containing hundreds of such biomolecules, so too will the need to refine the underlying computational models to accurately describe non-bonded interactions between those molecules. An outstanding challenge for the field lies in describing intrinsically disordered regions (IDRs) of the DNA-binding proteins, which are omnipresent and play important roles in the recruitment and formation of transient biomolecular complexes. One particularly exciting and challenging class of systems are NA-IDR condensates [40\*\*,79], the biological roles of which we are just beginning to grasp. Here, in addition to further refinement of the all-atom molecular force field, methods that permit adequate sampling of the conformational space in an entangled polymer melt environment will be needed [80,81]. In general, the field would greatly benefit from a closer integration of coarse-grained models, such as Martini [82], oxDNA [83], 3SPN.2 [84] and ABSINTH [85], with all-atom approaches, allowing for

mixed-resolution, accuracy-when-needed types of description of very large protein-NA systems.

## Conflict of interest statement

Nothing declared.

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