

# Lecture 1: General introduction

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3rd February, 2025

# Open quantum systems

**Open quantum systems:** a system that interacts with an external quantum system (the environment or a bath)



# Quantum measurement: indirect measurements



# Photon counting and Homodyne detection



Figure : Two types of detection : Photon counting and Homodyne detection.

# **Control features**

• **Open-loop control:** mostly with Hamiltonian control, implemented by  $H = H_0 + \sum_{j=1}^{n} u_j(t)H_j$ 

Major goal: Controllability and optimal control

 Closed-loop control: measurement-based feedback, coherent feedback

Major goal: Compensating decoherence induced by the environment

#### Measurement-based feedback



# Example: LKB photon box



Here 
$$\rho \rightarrow \frac{V_{\pm}\rho V_{\pm}^{\dagger}}{\operatorname{tr}(V_{\pm}\rho V_{\pm}^{\dagger})}$$
 with probability  $\operatorname{tr}(V_{\pm}\rho V_{\pm}^{\dagger})$ 

# Feedback stabilizing photon number state



C. Sayrin et al., Nature, 2011

Control input *u*: application of a unitary  $U_u = e^{-iuH}$ :  $\rho \rightarrow U_u \left( \frac{V_{\pm \rho} V_{\pm}}{\operatorname{tr}(V_{\pm \rho} V_{\pm}^{\dagger})} \right) U_u^{\dagger}$ 

# Other examples of measurement-based feedback



Campagne-Ibarcq et al., Nature, 2020

Vijay et al., Nature, 2012

# Coherent feedback (without measurement)



# Examples of coherent feedback





#### Photonic QEC network schematic Kerckhoff et al., PRL, 2010

Coherent feedback with a dynamic compensator Mabuchi, PRA, 2008

# Some references on quantum trajectories

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- Strunz, Diósi, Gisin, Yu, Quantum trajectories for Brownian motion, 1999

# Some references on feedback control

- Belavkin, Theory of the control of observable quantum systems, 1983
- Wiseman, Quantum theory of continuous feedback, 1994
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- Doherty, Habib, Jacobs, Mabuchi, Tan, Quantum feedback and classical control theory, 2000
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- van Handel, Stockton, Mabuchi, Feedback control of quantum state reduction
- Mirrahimi, van Handel, Stabilizing feedback controls for quantum systems, 2007
- James, Nurdin, Petersen, H-Infinity Control of Linear Quantum Stochastic Systems, 2008
- Gough, James, The series product and its application to quantum feedforward and feedback networks, 2009

# Heuristic derivation of the Lindblad equation <sup>1</sup>

 $(\mathcal{H}_{\mathsf{Total}}, \rho_{\mathsf{Total}}, \mathcal{H}_{\mathsf{Total}})$ 



$$ho_{ ext{Total},t} = U_t(
ho_0 \otimes 
ho_{ ext{Bath}})U_t^{\dagger} \quad 
ho_{ ext{Bath}} = \sum_{
u} \lambda_{
u} \ket{
u} ra{
u}$$

<sup>1</sup>Lindblad-Gorini-Kossakawski-Sudarshan, 1976.

Different approaches to derive quantum trajectories

# **Repeated interaction**



Consider a discrete-time system with interaction time  $\tau$  at time  $n\tau$  it comes in interaction with a fresh ancillary :

 $\mathcal{V}_{\mathcal{S}} = \mathsf{Tr}_{a} (V(\rho \otimes \sigma) V^{\dagger}), \quad V \text{ is unitary acting on } \mathcal{H}_{\mathcal{S}} \otimes \mathcal{H}_{B}$ 

$$\mathrm{d}\hat{\rho}_t = \mathcal{L}(\hat{\rho}_t)\mathrm{d}t + \left(L\hat{\rho}_t + \hat{\rho}_t L^{\dagger} - \mathrm{Tr}\left((L + L^{\dagger})\hat{\rho}_t\right)\hat{\rho}_t\right)\mathrm{d}W_t$$

# Quantum filtering



The evolution is described by the following unitary operator: <sup>2</sup>

$$\mathrm{d}U_t = \left\{ L\mathrm{d}A_t^{\dagger} - L^{\dagger}\mathrm{d}A_t - \frac{1}{2}L^{\dagger}L\mathrm{d}t - \mathrm{i}H\mathrm{d}t \right\} U_t, \quad U_0 = I$$

Take  $X_t = j_t(X)$  and the observation process  $(Y_t)_{t \ge 0}$ 

$$dj_t(X) = j_t(\mathcal{L}^{\dagger}(X)) dt + j_t([\mathcal{L}^{\dagger}, X]) dA_t + j_t([X, L]) dA_t^{\dagger}$$
$$dY_t = j_t(\mathcal{L} + \mathcal{L}^{\dagger}) dt + dA_t + dA_t^{\dagger}$$

<sup>&</sup>lt;sup>2</sup>Hudson and Parthasarathy, 1984.

# Quantum filtering

 $[Y_t, Y_r] = 0, t_0 \le r \le t,$  (self-non-demolition property)  $[X_t, Y_s] = 0, t_0 \le s \le t,$  (non-demolition property)

Quantum filter<sup>3</sup>:  $\hat{X}(t) \equiv \pi_t(X) = \mathbb{E}(j_t(X)|\mathcal{Y}_t)$ 

 $\mathrm{d}\pi_t(X) = \pi_t(\mathcal{L}^{\dagger}(X))\,\mathrm{d}t + \left(\pi_t(XL + L^{\dagger}X) - \pi_t(L + L^{\dagger})\pi_t(X)\right)\,\mathrm{d}W_t,$ 

where W is a Wiener process,  $dW_t = dY_t - \pi_t(L + L^{\dagger}) dt$ .

In the Schrödinger picture,  $\pi_t(X) = tr(\rho_t X)$ ,

$$\mathrm{d}\rho_t = \mathcal{L}(\rho_t)\,\mathrm{d}t + \left(L\rho_t + \rho_t L^{\dagger} - \mathrm{Tr}\left((L+L^{\dagger})\rho_t\right)\rho_t\right)\mathrm{d}W_t$$

<sup>&</sup>lt;sup>3</sup>Belavkin, 1980.

# Barchielli's approach

Probability space<sup>4</sup>: 
$$\left(\Omega, \mathcal{F}, \left(\mathcal{F}_{t}\right)_{t \geq 0}, \mathbb{Q}, \left(W_{t}\right)_{t \geq 0}\right)$$

#### Linear Stochastic Schrödinger Equation:

$$\mathrm{d}|\psi_t\rangle = \left(-\mathrm{i}H - \frac{1}{2}L^{\dagger}L\right)|\psi_t\rangle\mathrm{d}t + L|\psi_t\rangle\mathrm{d}W_t$$

Equation of the Propagator:

$$\mathrm{d}\boldsymbol{S}_t = \left(-\mathrm{i}\boldsymbol{H} - \frac{1}{2}\boldsymbol{L}^{\dagger}\boldsymbol{L}\right)\boldsymbol{S}_t\mathrm{d}t + \boldsymbol{L}\boldsymbol{S}_t\mathrm{d}\boldsymbol{W}_t$$

<sup>4</sup>Barchielli, 1983.

# Barchielli's approach

Probability space: 
$$\left(\Omega, \mathcal{F}, \left(\mathcal{F}_{t}\right)_{t \geq 0}, \mathbb{Q}, \left(\boldsymbol{W}_{t}\right)_{t \geq 0}\right)$$

Linear Stochastic Master Equation: denote  $\sigma_t = S_t \rho_0 S_t^{\dagger}$ 

$$\mathrm{d}\sigma_t = \mathcal{L}\left(\sigma_t\right)\mathrm{d}t + \left(L\sigma_t + \sigma_t L^{\dagger}\right)\mathrm{d}W_t$$

where

$$\mathcal{L}(\sigma_t) = -\mathrm{i}[H, \sigma_t] + L\sigma_t L^{\dagger} - \frac{1}{2} \left\{ L^{\dagger} L, \sigma_t \right\}$$

# Barchielli's approach

New probability space:

$$\left(\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \ge 0}, \mathbb{Q}, (W_t)_{t \ge 0}\right) \longrightarrow \left(\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \ge 0}, \tilde{\mathbb{Q}}, (B_t)_{t \ge 0}\right)$$

**Stochastic Master Equation:** Denote  $\rho_t = \sigma_t / \text{Tr}(\sigma_t)$ 

$$\mathrm{d}\rho_{t} = \mathcal{L}\left(\rho_{t}\right)\mathrm{d}t + \left(L\rho_{t} + \rho_{t}L^{\dagger} - \mathrm{Tr}\left((L+L^{\dagger})\rho_{t}\right)\rho_{t}\right)\mathrm{d}B_{t}$$

where

$$\mathcal{L}(\sigma_t) = -\mathrm{i}[H, \sigma_t] + L\sigma_t L^{\dagger} - \frac{1}{2} \left\{ L^{\dagger} L, \sigma_t \right\}$$

# Non-Markovian quantum trajectories

- Up to here, environmental correlation times are assumed negligibly short compared to the system's characteristic time scale (Markovian approximation)
- Strunz, Diósi, and Gisin, Open system dynamics with non-Markovian quantum trajectories, PRL, 1999
- Gardiner & Zoller, Quantum Noise: A Handbook of Markovian and Non-Markovian Quantum Stochastic Methods, 2004
- Breuer & Petruccione, The Theory of Open Quantum Systems, 2002

# Thank you!