

# Kinetic Mixing in $SO(10)$ , Gravitational Waves and PTA data

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Hearing BSM with Cosmic Sources of Gravitational Waves  
@ ICTS - TIFR, Bengaluru

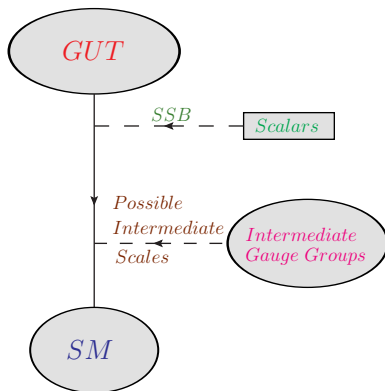
- 1 Introduction
- 2  $SO(10)$  Symmetry Breaking, Monopoles and Strings
- 3 GWB from Quasi-stable Strings and PTA Data
- 4 Unification Solutions
- 5 Summary

# 1 *Introduction*

# Grand Unification Beyond the SM

- The basic idea in a Grand Unified Theory (GUT) is that the SM,  $SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$ , is embedded in a larger simple group,  $\mathcal{G}$ .

## Schematic view



- $SU(5)$  (rank = 4):  $\bar{5} + 10 \Rightarrow$  SM fermions.

Georgi, Glashow, PRL **32**, 438 (1974)

- $SO(10)$  (rank = 5):  $16 \Rightarrow$  SM fermions  $\oplus \nu_L^C$ .

Fritzsch, Minkowski, Ann. Phys. **93**, 93-266 (1975)

- $E(6)$  (rank = 6):  $27 \Rightarrow$  SM fermions  $\oplus \nu_L^C \oplus$

$$\underbrace{\left(2, \pm \frac{1}{2}, 1\right) + \left(1, -\frac{1}{3}, 3\right) + \left(1, \frac{1}{3}, \bar{3}\right) + (1, 0, 1)}_{\text{Exotic fermions}}$$

Exotic fermions

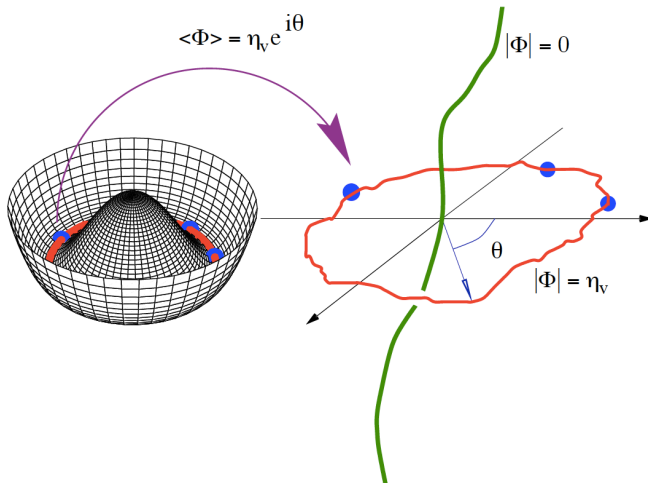
Gursey, Ramond, Sikivie, PLB **60** (1976) 177

Shafi, PLB **79** (1978) 301

# Prediction of topological defects

- Topological defects may appear during the SSB of a group  $\mathcal{G}$  down to its subgroup  $\mathcal{H}$ .
- Non-trivial homotopy group  $\Pi_k(\mathcal{M})$  of the vacuum manifold ( $\mathcal{M} = \mathcal{G}/\mathcal{H}$ ) implies formation of topological defects.
- Various types of topological defects which can be formed are : domain walls ( $k = 0$ ), cosmic strings ( $k = 1$ ), monopoles ( $k = 2$ ) etc.  
 $SU(N)_X \rightarrow U(1)_X$ : Monopoles,  $U(1) \rightarrow Z_N$ : Strings

# Cosmic string



Vachaspati et. al. arXiv:1506.04039

# Cosmic string network

- String tension  $\mu \simeq \pi v^2$ ,  $v$  is the VEV that form the string.
- Strings inter-commute, form loops, radiate GWs and the evolution of the network enters a 'scaling' regime.
- Scaling energy density  $\rho_s \sim \mu/t^2$ . Critical density:  $\rho_c \sim 1/Gt^2$  in RD and MD.

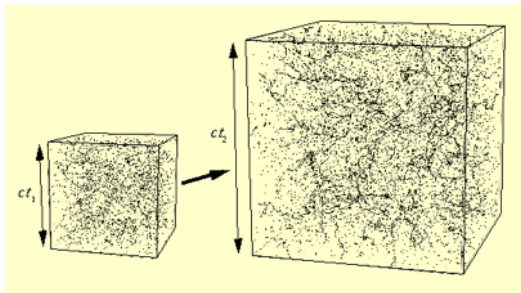


Image source: [ctc.cam.ac.uk](http://ctc.cam.ac.uk)

Kibble, NPB 252 (1985) 227; Vachaspati, Vilenkin PRD 31 (1985) 3052; Bennett, Bouchet, PRL 60 (1988) 257 ...



# Strings and gravitational waves

- Loops of **initial length**  $l_i = \alpha t_i$  ( $\alpha \simeq 0.1$ ) decay via emission of gravity waves. Blanco-Pillado, Olum, Shlaer, Phys. Rev. D **89** (2014) 023512

$$\frac{dE_{\text{GW}}^{(k)}}{dt} = \Gamma_k G\mu^2; \quad \Gamma_k \propto k^{-n} \quad \text{with } n = \begin{cases} 4/3 & \text{cusps} \\ 5/3 & \text{kinks} \\ 2 & \text{kink-kink collisions.} \end{cases}$$

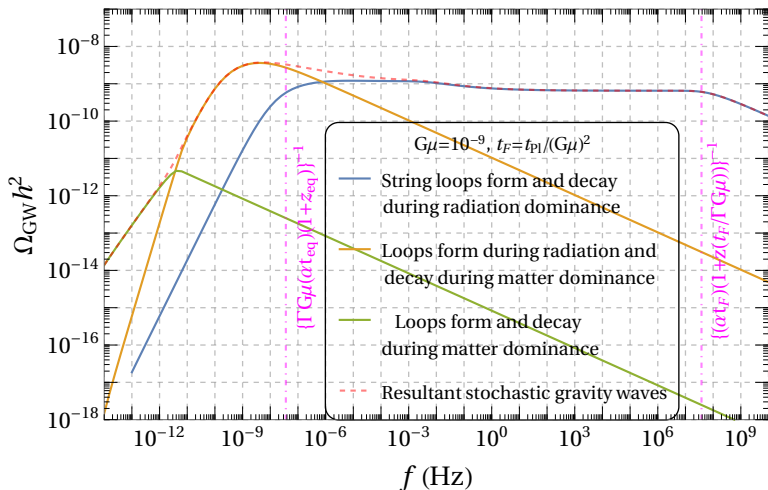
- The redshifted frequency of a normal mode  $k$ , emitted at time  $\tilde{t}$ , as observed today, is given by Vilenkin, Shellard, 1994, CUP

$$f = \frac{a(\tilde{t})}{a(t_0)} \frac{2k}{\alpha t_i - \Gamma G\mu(\tilde{t} - t_i)}, \quad \text{with } \Gamma = \sum \Gamma_k \sim 50$$

Redshift

$$\frac{dl}{dt} = -\Gamma G\mu \Rightarrow \text{Loop size at time } \tilde{t}$$

# Stochastic gravitational wave background



# Observational constraints from defects

- Stable domain walls contradict standard cosmology.  
Zeldovich, Kobzarev, Okun, Zh. Eksp. Teor. Fiz. **67**, 3-11 (1974)
- Upper bound on comoving monopole number density:  
 $Y_M = n_M/s \gtrsim 10^{-27}$ . MACRO: EPJC 25 511, IceCube: PRL 128 (2022) 051101, ANTARES: JHEAp 34 (2022) 1, ...
- CMB constraint on stable strings:  $G\mu \lesssim 10^{-7}$ .  
PhysRevD.93.123503, ...
- LIGO-VIRGO O3 data constraint on “undiluted” strings:  
 $G\mu \lesssim 10^{-7}$  around **decaHz** frequencies. PhysRevLett.126.241102
- PTA experiments put a constraint on stable cosmic strings :  
 $G\mu \lesssim 10^{-10}$  around the **nanoHertz** frequencies.  
arXiv:2306.16219,...

## 2 *SO(10) Symmetry Breaking, Monopoles and Strings*

# SO(10) symmetry breaking, monopoles and strings

$$\begin{aligned}SO(10) &\xrightarrow[M_U]{\langle 210 \rangle} SU(4)_c \times SU(2)_L \times SU(2)_R \\ &\xrightarrow[M_U]{\langle (15,1,3) \in 210 \rangle} SU(3)_c \times U(1)_{B-L} \times SU(2)_L \times U(1)_R \\ &\xrightarrow[M_R]{\langle (1,1,1, -\frac{1}{2}) \in 16 \rangle} SU(3)_c \times SU(2)_L \times U(1)_Y.\end{aligned}$$

- Symmetry breaking  $SU(4)_C \rightarrow SU(3)_C \times U(1)_{B-L}$  produces ‘Red’ monopoles with magnetic fluxes

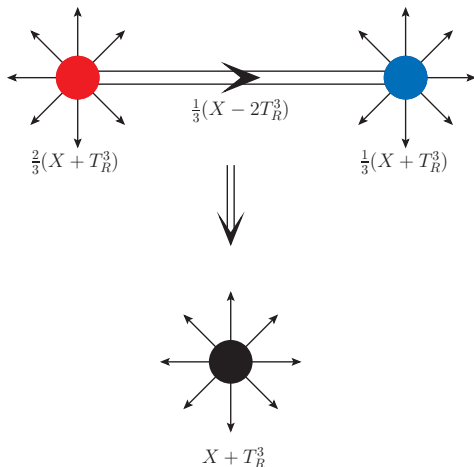
$$X \equiv B - L + 2T_c^8/3 = \text{diag}(1, 1, -1, -1).$$

- $SU(2)_R \rightarrow U(1)_R$  generates ‘Blue’ monopoles with fluxes

$$T_R^3 = \text{diag}(1, -1).$$

# Strings connecting monopoles

- $U(1)_R \times U(1)_{B-L} \rightarrow U(1)_Y$  generates topologically unstable strings.
- These strings connects
  - ① a blue monopole to a red monopole.
  - ② a monopole to its anti-monopole.
  - ③ ends on itself forming a loop.
- Red and blue monopoles combined to form stable Schwinger monopoles.



Lazarides, Shafi, JHEP 10 (2019) 193

Lazarides, RM, Shafi, JCAP 05 (2024)

### 3 *GWB from Quasi-stable Strings and PTA Data*

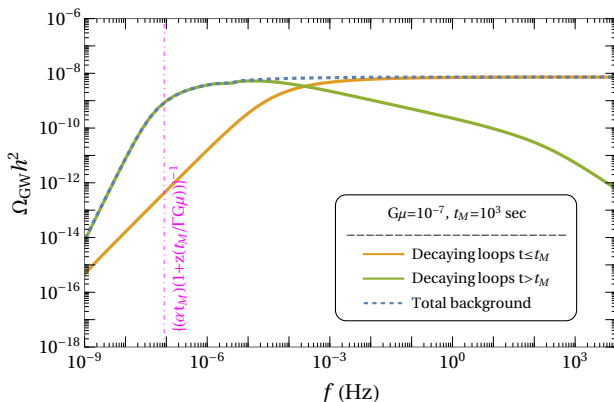
# Formation of quasi-stable strings

- Magnetic monopoles, created prior to the cosmic strings, experience “**partial**” inflation.
- The strings make random walks with step of the order of the horizon, and form a network of stable strings before the horizon reentry of the monopoles.
- The strings inter-commute and form loops which decay into gravitational waves.  $f = \frac{1}{1+z(t)} \frac{2k}{l(t)}$ ,  $k \in Z^+$ .
- As monopoles reenter the horizon at a cosmic time  $t_M$ , we obtain monopoles connected by string segments which eventually decay.

Lazarides, **RM**, Shafi, JCAP 08 (2022) 042

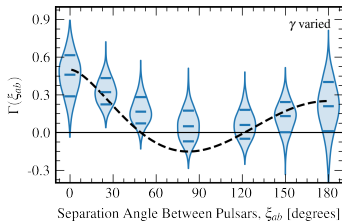


# Gravitational waves from quasi-stable strings



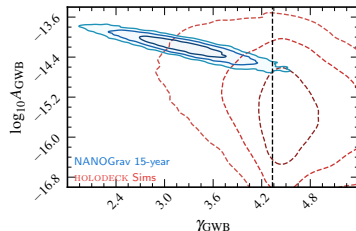
- Large string loops and segments ( $> 2t_M$ ) are absent.
- Gravitational wave spectrum in the low frequency region  $f \lesssim 1/t_M(1 + z(t_M/\Gamma G\mu))$  is suppressed.

# Evidence of GWB in PTA

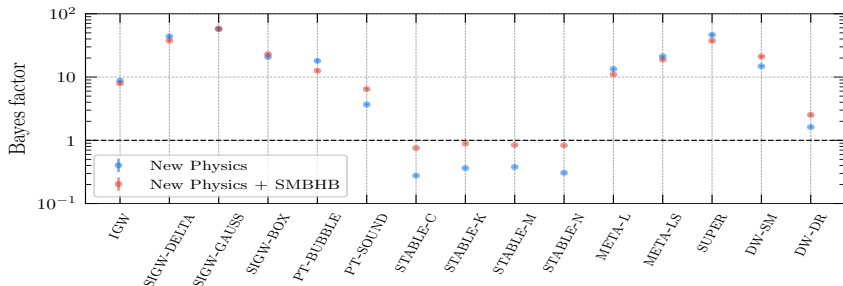


NG15:2306.16219,...

# Mergers of SMBHBs?

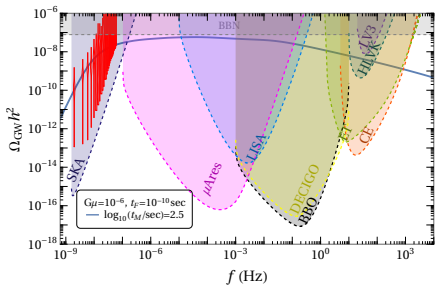


# Cosmological Origin? New Physics?

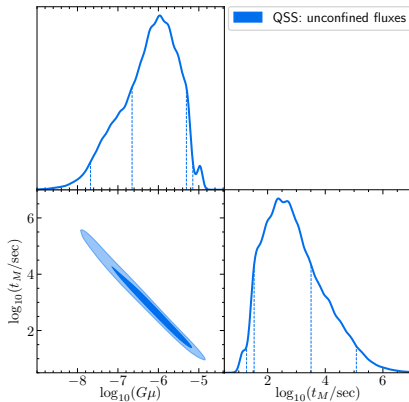


# Gravitational waves and NANOGrav 15 year data

Parameters	Bayesian Credible Intervals	
	68%	95%
$\log_{10}(G\mu)$	$[-6.65, -5.31]$	$[-7.68, -5.15]$
$\log_{10}(t_M/\text{sec})$	$[1.53, 3.51]$	$[1.27, 5.08]$



RM, Shafi, JHEP 10 (2024) 157



## 4 *Unification Solutions*

# Dim-5 operators and kinetic mixing of $U(1)_{B-L} \times U(1)_R$

- $F_{\mu\nu}^p = \partial_\mu A_\nu^p - \partial_\nu A_\mu^p$  ( $p = X, R$ ;  $X \equiv B - L$ ) mixes as:

$$-\frac{1}{4}F_{\mu\nu}^X F^{X\mu\nu} - \frac{1}{4}F_{\mu\nu}^R F^{R\mu\nu} - \frac{\epsilon}{2}F_{\mu\nu}^X F^{R\mu\nu},$$

$\epsilon$  is the mixing parameter.

Holdom, PLB 166 (1986)

- Dim-5 operator suppressed by the cut-off scale  $\Lambda$ :

$$\frac{\mathcal{C}}{\Lambda} G^{a\mu\nu} \Phi_{ai} G_{\mu\nu}^i, \quad \Phi_{ai} \equiv (15, 1, 3) \in 210_H$$

induces  $\epsilon \sim \mathcal{O}(M_U/\Lambda)$  at the GUT scale as  $\Phi_{ai}$  gets a VEV  $\sim M_U$ .

RM, Shafi, JHEP 10 (2024) 157

- The canonical form is obtained through change of basis  $A^p \rightarrow B^p$ :

$$\begin{pmatrix} A_\mu^X \\ A_\mu^R \end{pmatrix} = \begin{pmatrix} 1 & \frac{-\epsilon}{\sqrt{1-\epsilon^2}} \\ 0 & \frac{1}{\sqrt{1-\epsilon^2}} \end{pmatrix} \begin{pmatrix} B_\mu^X \\ B_\mu^R \end{pmatrix}.$$

# Abelian mixing and matching of gauge couplings

- In the covariant derivative part, the diagonal gauge coupling  $\text{diag}[g_X, g_R] \rightarrow G$ :

$$G = \begin{pmatrix} g_{XX} & g_{XR} \\ 0 & g_{RR} \end{pmatrix},$$

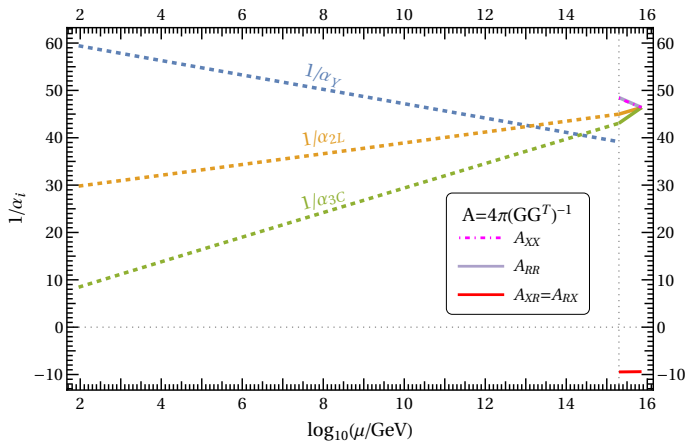
$$\text{with } g_{XX} = g_X, \quad g_{XR} = -\frac{\epsilon g_X}{\sqrt{1 - \epsilon^2}}, \quad g_{RR} = \frac{g_R}{\sqrt{1 - \epsilon^2}}.$$

- At the GUT scale  $M_U$ :  $g_X = g_R \equiv g_U$ .
- We have the matching condition at the breaking scale  $M_R$  given by

$$\frac{1}{\alpha_Y(M_R)} = 4\pi P(GG^T)^{-1}P^T \quad \text{with } P = \left(\sqrt{\frac{2}{5}}, \sqrt{\frac{3}{5}}\right)$$

$$\Rightarrow \frac{1}{g_Y^2} = \frac{3g_{XX}^2 - 2\sqrt{6}g_{XX}g_{XR} + 2(g_{XR}^2 + g_{RR}^2)}{5g_{XX}^2g_{RR}^2}.$$

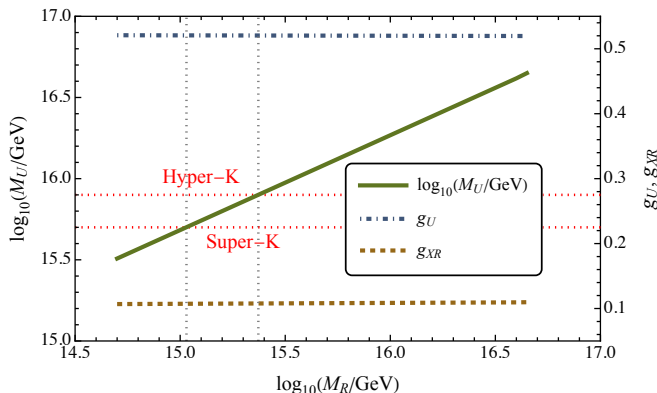
# Unification solution



- The unification occurs for  $\epsilon \sim M_U/\Lambda \sim 0.1$

RM, Shafi, JHEP 10 (2024) 157

# Unification solution



- The unification occurs for  $\epsilon \sim M_U/\Lambda \sim 0.1$ .
- String tension ( $\mu \sim \pi M_R^2$ ):  $\log_{10}(G\mu) \in [-7.7, -4.7]$
- Seesaw scale  $\sim \epsilon M_R^2/M_U \sim 10^{14}$  GeV.



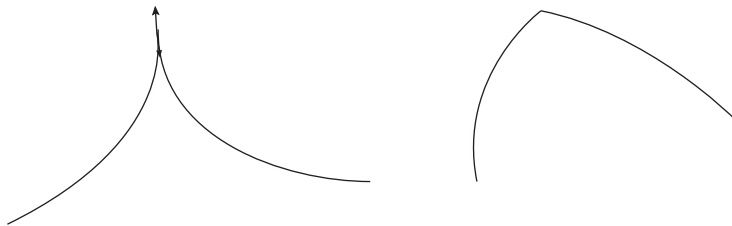
## 5 *Summary*

- Realistic grand unified theory of  $SO(10)$  predicts formation of superheavy strings which connects monopoles.
- Quasistable strings can be formed if the monopoles experience partial inflation.
- The GWs from superheavy QSS with  $G\mu \sim 10^{-6}$  can explain the evidence of GWs in recent PTA data with a monopole horizon reentry time  $t_M \sim 10^2$  sec.
- Realistic grand unified theory of  $SO(10)$  predicts the formation of such a quasistable string network.

*Thank You*

*Back up slides*

# Cusp and Kink



# Pulsar Timing Arrays

- Pulsars are rapidly spinning neutron stars with a strong magnetic field  $\Rightarrow$  Radiate beam of radio waves.
- Repeating pulses are observed as the radio beam intersects the observers periodically.
- Millisecond pulsar (MSP) produces exceedingly stable and regular pulse profile  $\Rightarrow$  “Perfect Clock”.

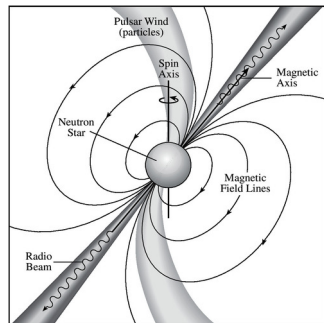
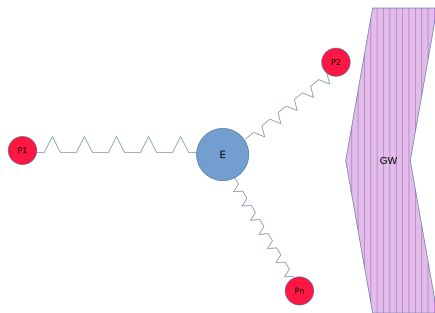


Image source: K.R. Lang, NASA's Cosmos

# Pulsar Timing Arrays



- Measurement of the time of arrival (ToA) of pulses can reveal tiny distortion of spacetime fabric due to gravity waves (GWs)  $\Rightarrow$  Pulsar timing!

Image source: [Wikipedia](#)

- Difference between observed ToA and the expected ToA from timing model gives time residual.
- Time residual contains information about other signals like GWs.

# Pulsar Timing Arrays

- Impossible to distinguish between GWs signal and other source of signal in the timing residual of a single pulsar.
- Need correlations between the timing residuals of different pulsars  $\Rightarrow$  Pulsar Timing Array (PTA).
- Gravity waves generate unique quadrupolar correlations between timing residuals of pulsar pairs.
- Correlations depend on the angular separations between the pulsar pairs and follow the Hellings and Downs correlation curve. **APJ. 265, L39 (1983)**

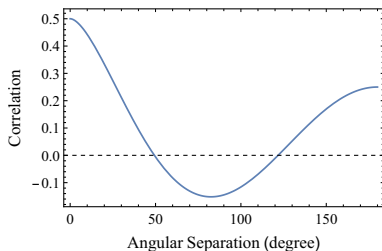
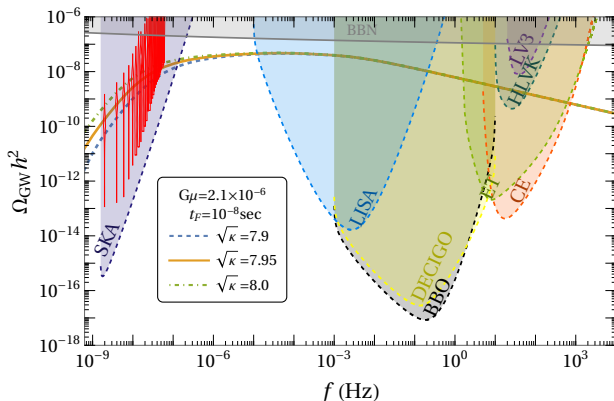


Figure: Hellings and Downs curve.



# Metastable Strings: PTA and LIGO-VIRGO



- The GWs from MSS explain PTA data at nanoHertz frequency, but violate the bound from LIGO-VIRGO third observing run!
- An early matter domination or partial inflation can reduce the spectra at high  $f$ .

Lazarides, **RM**, Moursy, Shafi, JCAP 03 (2024) 006

**RM**, Park, JCAP 01 (2024) 015

# Formation of Metastable Strings (MSS)

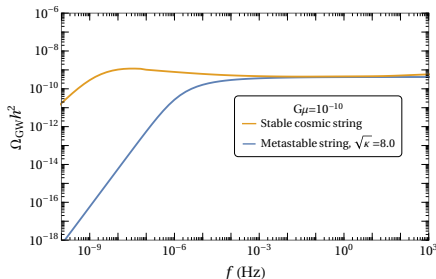
- Magnetic monopoles, created prior to the cosmic strings, experience inflation.
- The lifetime of decay of the strings via quantum mechanical tunneling is much smaller than the age of Universe.
- The strings form a network of stable strings before the time  $t_s = 1/\sqrt{\Gamma_d}$ .
- The strings network disappear at a time  $t_e \sim 1/\sqrt{\Gamma_d \Gamma G \mu}$ .

Leblond, Shlaer, Siemens, PRD **79** (2009) 123519

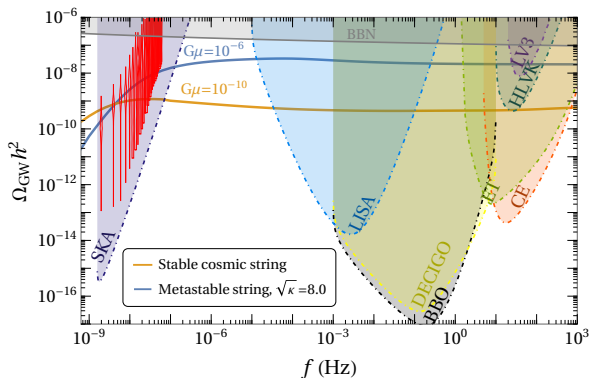
Buchmuller, Domcke, Schmitz, JCAP **12** (2021) 006

# Metastable Strings and GWs

- The strings inter-commute and form loops which decay into gravitational waves.
- String loops larger than  $\alpha t_s$  are absent.
- Gravitational wave spectrum in the low frequency region,  $f \lesssim 1/\Gamma G\mu t_e(1+z_e)$ , becomes suppressed.



# PTA and Observational Prospects of Strings



- Pulsar Timing Arrays (PTAs)

- ① found evidence of a stochastic background which can be explained by superheavy “metastable” strings  $G\mu \sim 10^{-6} - 10^{-5}$ .
- ② put a constraint “undiluted stable” cosmic strings :  $G\mu \lesssim 10^{-10}$ .

in the **nanoHertz** frequencies. [arXiv:2306.16219,...](https://arxiv.org/abs/2306.16219)